

**Supplement**  
Inflationary Perturbations

# Field perturbations

- Let's define the perturbed scalar field as  $\phi = \phi_0 + \phi_1$  where  $\phi_0$  is the unperturbed field (i.e.  $\phi_1 = \delta\phi$ )
- Field fluctuations obey a damped harmonic oscillator equation (with dots referring to conformal time derivatives)

$$\ddot{\phi}_1 + 2\frac{\dot{a}}{a}\dot{\phi}_1 + k^2\phi_1 \approx 0$$

- We want a simple harmonic oscillator that we can then quantize so define  $u \equiv a\phi_1$

$$\dot{u} = \dot{a}\phi_1 + a\dot{\phi}_1$$

$$\ddot{u} = \ddot{a}\phi_1 + 2\dot{a}\dot{\phi}_1 + a\ddot{\phi}_1$$

$$\ddot{u} + \left[k^2 - \frac{\ddot{a}}{a}\right]u = 0$$

- Note Friedmann equations say if  $p = -\rho$ ,  $\ddot{a}/a = 2(\dot{a}/a)^2$

# Harmonic Oscillator

- Now let's look at the oscillator equation

$$\ddot{u} + \left[ k^2 - 2 \left( \frac{\dot{a}}{a} \right)^2 \right] u = 0$$

or for conformal time measured from the end of inflation

$$\tilde{\eta} = \eta - \eta_{\text{end}}$$

$$\tilde{\eta} = \int_{a_{\text{end}}}^a \frac{da}{Ha^2} \approx -\frac{1}{aH}$$

- So we can rewrite this as

$$\ddot{u} + \left[ k^2 - \frac{2}{\tilde{\eta}^2} \right] u = 0$$

# Quantum Fluctuations

- Simple harmonic oscillator  $\ll$  Hubble length

$$\ddot{u} + k^2 u = 0$$

- Quantize the simple harmonic oscillator

$$\hat{u} = u(k, \tilde{\eta})\hat{a} + u^*(k, \tilde{\eta})\hat{a}^\dagger$$

where  $u(k, \tilde{\eta})$  satisfies classical equation of motion and the creation and annihilation operators satisfy

$$[a, a^\dagger] = 1, \quad a|0\rangle = 0$$

- Normalize wavefunction  $[\hat{u}, d\hat{u}/d\tilde{\eta}] = i$

$$u(k, \eta) = \frac{1}{\sqrt{2k}} e^{-ik\tilde{\eta}}$$

# Quantum Fluctuations

- Zero point fluctuations of ground state

$$\begin{aligned}\langle u^2 \rangle &= \langle 0 | u^\dagger u | 0 \rangle \\ &= \langle 0 | (u^* \hat{a}^\dagger + u \hat{a}) (u \hat{a} + u^* \hat{a}^\dagger) | 0 \rangle \\ &= \langle 0 | \hat{a} \hat{a}^\dagger | 0 \rangle |u(k, \tilde{\eta})|^2 \\ &= \langle 0 | [\hat{a}, \hat{a}^\dagger] + \hat{a}^\dagger \hat{a} | 0 \rangle |u(k, \tilde{\eta})|^2 \\ &= |u(k, \tilde{\eta})|^2 = \frac{1}{2k}\end{aligned}$$

- Classical equation of motion take this quantum fluctuation outside horizon where it freezes in.

# Slow Roll Limit

- Classical equation of motion then has the exact solution

$$u = \frac{1}{\sqrt{2k}} \left( 1 - \frac{i}{k\tilde{\eta}} \right) e^{-ik\tilde{\eta}}$$

- For  $|k\tilde{\eta}| \ll 1$  (late times,  $\gg$  Hubble length) fluctuation freezes in

$$\lim_{|k\tilde{\eta}| \rightarrow 0} u = -\frac{1}{\sqrt{2k}} \frac{i}{k\tilde{\eta}} \approx \frac{iHa}{\sqrt{2k^3}}$$

$$\phi_1 = \frac{iH}{\sqrt{2k^3}}$$

- Power spectrum of field fluctuations

$$\Delta_{\phi_1}^2 = \frac{k^3 |\phi_1|^2}{2\pi^2} = \frac{H^2}{(2\pi)^2}$$