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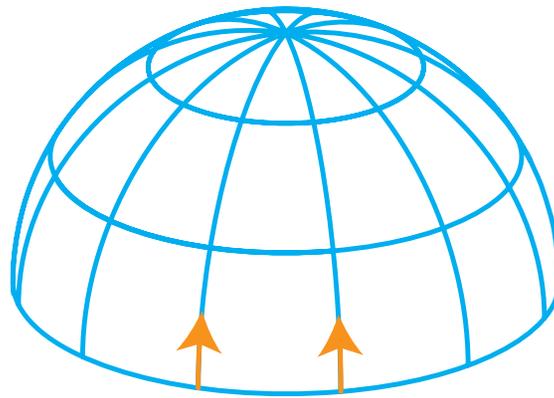
Cosmic Geometry

Newton vs Einstein

- Even though locally Newtonian gravity is an excellent approximation to General Relativity, in cosmology we deal with spatial and temporal scales across which the global picture benefits from a basic understanding of General Relativity.
- An example is: as we have seen in the previous set of notes, it is much more convenient to think of the space between galaxies expanding rather than galaxies receding through space
- While latter is a good description locally, its preferred coordinates place us at center and does not allow us to talk about distances beyond which galaxies are receding faster than light - though these distances as we shall see are also not directly observable
- To get a global picture of the expansion of the universe we need to think geometrically, like Einstein not Newton
- Best of both: think globally (Einstein), act locally (Newton)

Gravity as Geometry

- Einstein says Gravity as a force is really the geometry of spacetime
- Force between objects is a fiction of geometry - imagine the curved space of the 2-sphere - e.g. the surface of the earth
- Two people walk from equator to pole on lines of constant longitude
- Intersect at poles as if an attractive force exists between them
- Both walk on geodesics or straight lines of the shortest distance



Gravity as Geometry

- General relativity has two aspects
 - A **metric** theory: geometry tells matter how to **move**
 - **Field equations**: matter tells geometry how to **curve**
- Metric defines distances or separations in the spacetime and freely falling matter moves on a path that extremizes the distance
- Expansion of the universe carries two corresponding pieces
 - Friedmann-Robertson-Walker geometry or metric tells matter, including light, how to move – allows us to chart out the expansion with light
 - Matter content of the universe tells it how to expand – allows us to infer the components of the universe
- Useful to **separate** out these two pieces both conceptually and for understanding **alternate cosmologies**

FRW Geometry

- FRW geometry = homogeneous and isotropic on large scales
- Universe observed to be nearly isotropic (e.g. CMB, radio point sources, galaxy surveys)
- Copernican principle: we're not special, must be isotropic to all observers (all locations)

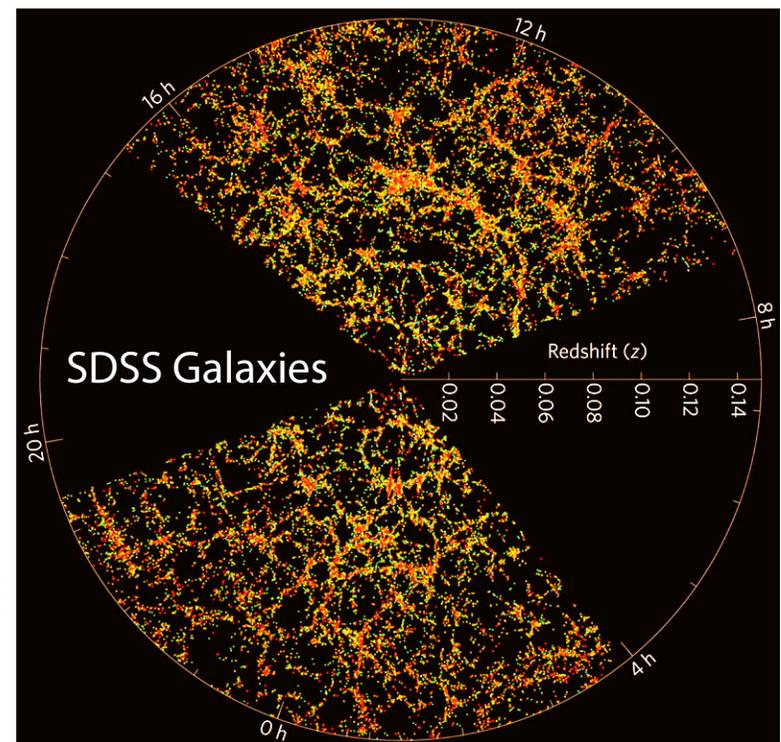
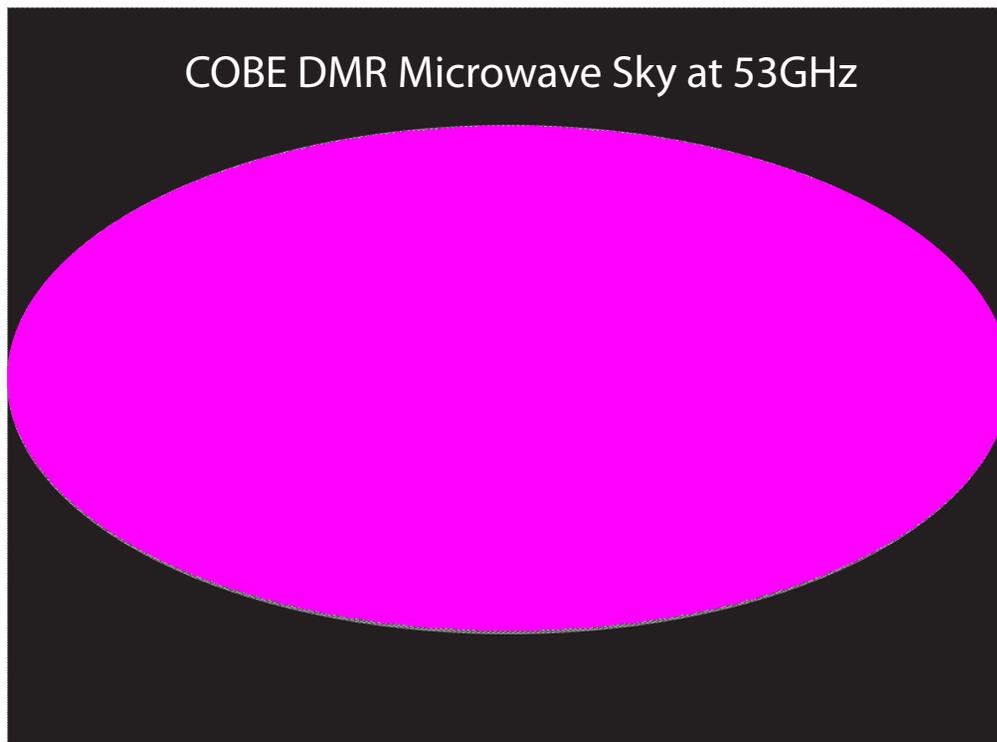
Implies homogeneity

Verified through galaxy redshift surveys

- FRW cosmology (homogeneity, isotropy & field equations) generically implies the expansion of the universe, except for special unstable cases

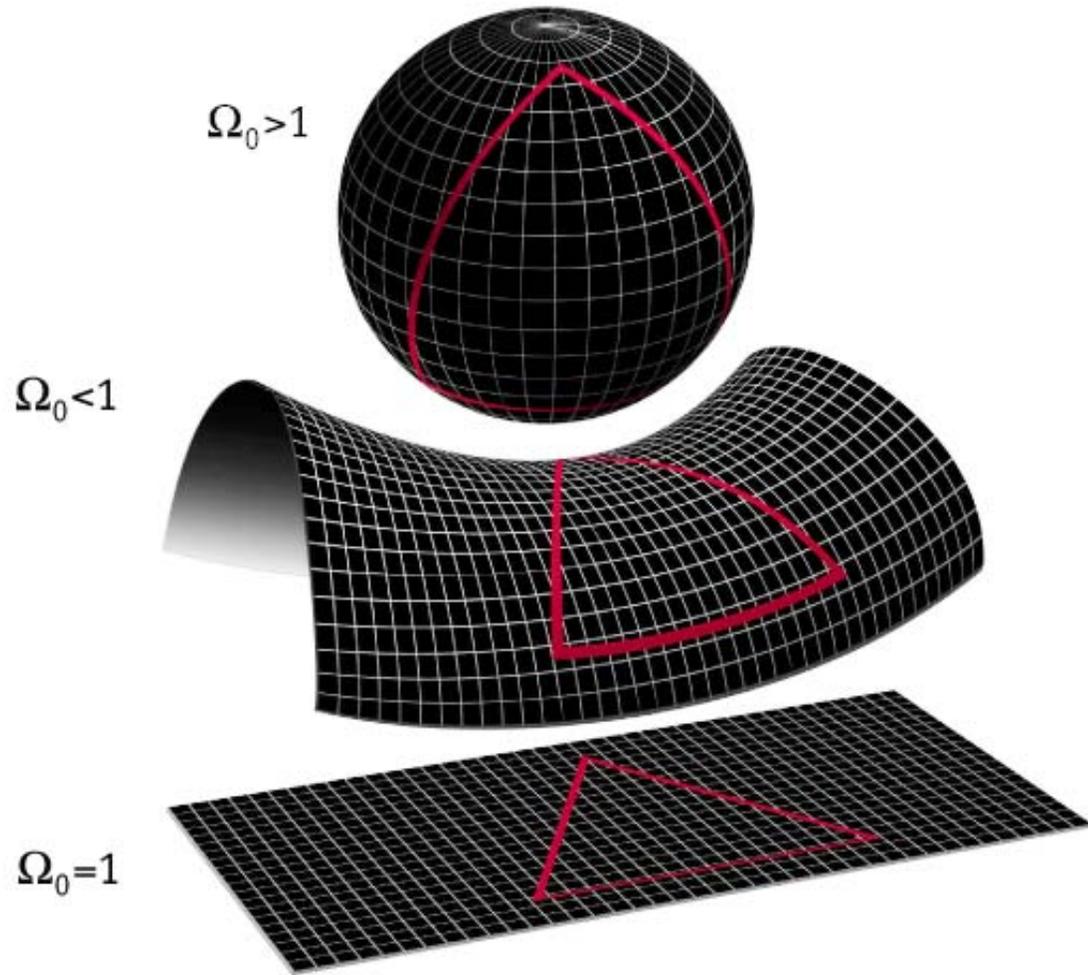
Isotropy & Homogeneity

- Isotropy: CMB isotropic to 10^{-3} , 10^{-5} if dipole subtracted
- Redshift surveys show return to homogeneity on the $>100\text{Mpc}$ scale



FRW Geometry

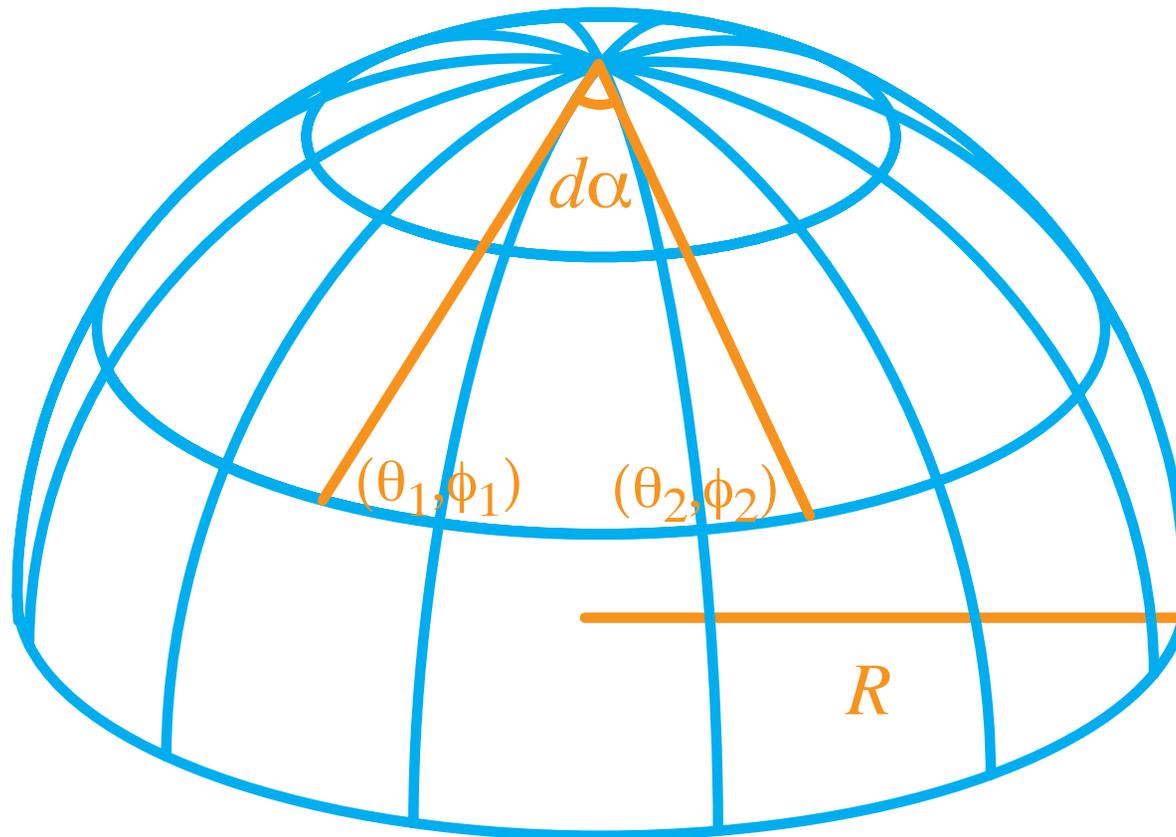
- Spatial geometry is that of constant curvature
 - Positive: sphere
 - Negative: saddle
 - Flat: plane
- Metric tells us how to measure distances on this surface



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FRW Geometry

- Closed: sphere of radius R and (real) curvature $K = 1/R^2$
- Suppress 1 dimension α represents total angular separation between two points on the sky (θ_1, ϕ_1) and (θ_2, ϕ_2)



FRW Geometry

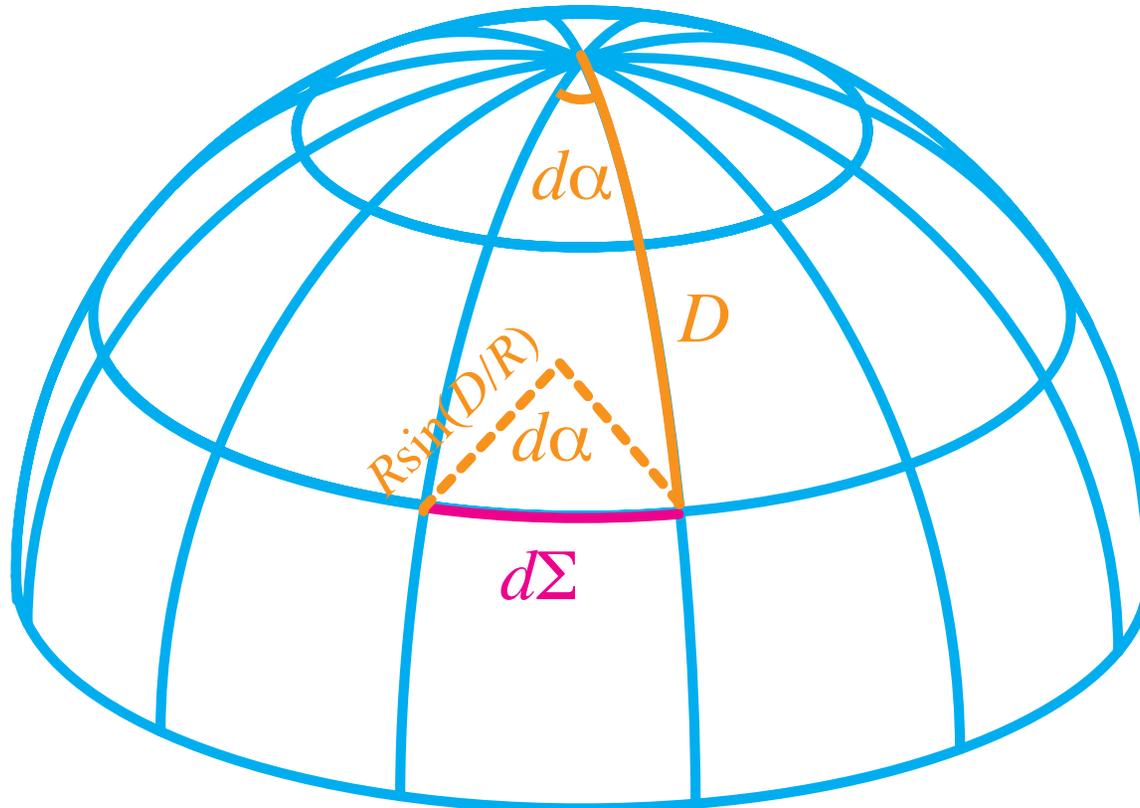
- Geometry tells matter how to move: take (null) geodesic motion for light along this generalized sense of longitude or radial distance D
- This arc distance is the distance our photon traveler sees
- We receive light from two different trajectories as observer at pole
- Compared with our Euclidean expectation that the angle between the rays should be related to the separation at emission Σ as $d\alpha \approx \Sigma/D$ the angular size appears larger because of the “lensing” magnification of the background
- This leads to the so called angular diameter distance - the most relevant sense of distance for the observer
- In General Relativity, we are free to use any distance coordinate we like but the two have distinct uses

FRW Geometry

- To define the angular diameter distance, look for a D_A such that

$$d\Sigma = D_A d\alpha$$

Draw a circle at the distance D , its radius is $D_A = R \sin(D/R)$

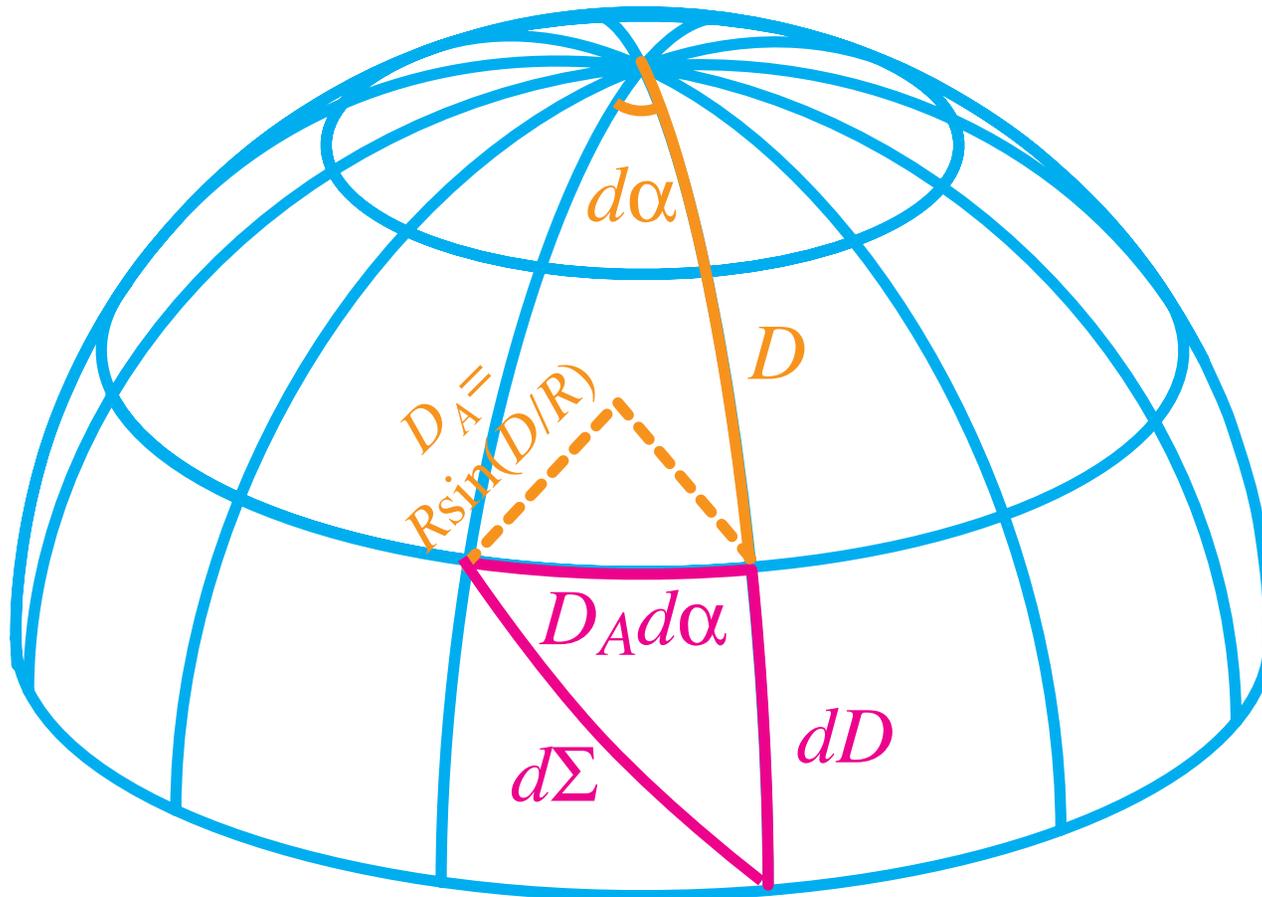


FRW Geometry

- Angular diameter distance
- Positively curved geometry $D_A < D$ and objects are further than they appear
- Negatively curved universe R is imaginary and
$$R \sin(D/R) = i|R| \sin(D/i|R|) = |R| \sinh(D/|R|)$$
and $D_A > D$ objects are closer than they appear
- Flat universe, $R \rightarrow \infty$ and $D_A = D$

FRW Geometry

- Now add that point 2 may have a different radial distance
- What is the distance $d\Sigma$ between points 1 (θ_1, ϕ_1, D_1) and point 2 (θ_2, ϕ_2, D_2) , separated by $d\alpha$ in angle and dD in distance?



Angular Diameter Distance

- For small angular and radial separations, space is nearly flat so that the Pythagorean theorem holds for differentials

$$d\Sigma^2 = dD^2 + D_A^2 d\alpha^2$$

- Now restore the fact that the angular separation can involve two angles on the sky - the curved sky is just a copy of the spherical geometry with unit radius that we were suppressing before

$$\begin{aligned} d\Sigma^2 &= dD^2 + D_A^2 d\alpha^2 \\ &= dD^2 + D_A^2 (d\theta^2 + \sin^2 \theta d\phi^2) \end{aligned}$$

- D_A useful for **describing observables** (flux, angular positions)
- D useful for **theoretical constructs** (causality, relationship to temporal evolution)

Alternate Notation

- Aside: line element is often also written using D_A as the coordinate distance

$$dD_A^2 = \left(\frac{dD_A}{dD} \right)^2 dD^2$$

$$\left(\frac{dD_A}{dD} \right)^2 = \cos^2(D/R) = 1 - \sin^2(D/R) = 1 - (D_A/R)^2$$

$$dD^2 = \frac{1}{1 - (D_A^2/R)^2} dD_A^2$$

and defining the curvature $K = 1/R^2$ the line element becomes

$$d\Sigma^2 = \frac{1}{1 - D_A^2 K} dD_A^2 + D_A^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

where $K < 0$ for a negatively curved space

Line Element or Metric Uses

- Metric also defines the volume element

$$\begin{aligned}dV &= (dD)(D_A d\theta)(D_A \sin \theta d\phi) \\ &= D_A^2 dD d\Omega\end{aligned}$$

where $d\Omega = \sin \theta d\theta d\phi$ is solid angle

- Most of classical cosmology boils down to these three quantities, (comoving) radial distance, (comoving) angular diameter distance, and volume element
- For example, distance to a high redshift supernova, angular size of the horizon at last scattering and BAO feature, number density of clusters...

Comoving Coordinates

- Remaining degree of freedom (preserving homogeneity and isotropy) is the **temporal evolution** of overall **scale factor**
- Relates the **geometry** (fixed by the radius of curvature R) to **physical coordinates** – a function of time only

$$d\sigma^2 = a^2(t)d\Sigma^2$$

our conventions are that the scale factor today $a(t_0) \equiv 1$

- Similarly **physical distances** are given by $d(t) = a(t)D$,
 $d_A(t) = a(t)D_A$.
- Distances in **upper case** are **comoving**; lower, physical
Comoving coordinates do not change with time and
Simplest coordinates to work out geometrical effects

Time and Conformal Time

- Spacetime separation (with $c = 1$)

$$\begin{aligned} ds^2 &= -dt^2 + d\sigma^2 \\ &= -dt^2 + a^2(t)d\Sigma^2 \end{aligned}$$

- Taking out the scale factor in the time coordinate

$$ds^2 \equiv a^2(t) (-d\eta^2 + d\Sigma^2)$$

$d\eta = dt/a$ defines **conformal time** – useful in that photons travelling radially from observer on null geodesics $ds^2 = 0$

$$\Delta D = \Delta\eta = \int \frac{dt}{a}$$

so that **time** and **distance** may be interchanged

FRW Metric

- Aside for advanced students: Relationship between coordinate differentials and space-time separation defines the **metric** $g_{\mu\nu}$

$$ds^2 \equiv g_{\mu\nu} dx^\mu dx^\nu = a^2(\eta)(-d\eta^2 + d\Sigma^2)$$

Einstein summation - **repeated** lower-upper pairs **summed**

- Usually we will use **comoving coordinates** and **conformal time** as the x^μ unless otherwise specified – metric for other choices are related by $a(t)$
- Aside: scale factor plays the role of a conformal rescaling (which preserves spacetime “angles”, i.e. light cone and causal structure - hence conformal time)

Horizon

- Distance travelled by a photon in the whole lifetime of the universe defines the **horizon**
- Since $ds = 0$, the horizon is simply the **elapsed conformal time**

$$D_{\text{horizon}}(t) = \int_0^t \frac{dt'}{a} = \eta(t)$$

- Horizon always **grows with time**
- Always a point in time before which two observers separated by a distance D could **not** have been in **causal contact**
- **Horizon problem**: why is the universe homogeneous and **isotropic** on large scales especially for objects seen at early times, e.g. CMB, when horizon small

Special vs. General Relativity

- From our class perspective, the big advantage of comoving coordinates and conformal time is that we have largely reduced general relativity to special relativity
- In these coordinates, aside from the difference between D and D_A , we can think of photons propagating in flat spacetime
- Now let's relate this discussion to observables
- Rule of thumb to avoid dealing with the expansion directly:
 - Convert from physical quantities to conformal-comoving quantities at emission
 - In conformal-comoving coordinates, light propagates as usual
 - At reception $a = 1$, conformal-comoving coordinates are physical, so interpret as usual

Hubble Parameter

- Useful to define the expansion rate or Hubble parameter

$$H(t) \equiv \frac{1}{a} \frac{da}{dt} = \frac{d \ln a}{dt}$$

fractional change in the scale factor per unit time - $\ln a = N$ is also known as the **e-folds** of the expansion

- Cosmic time becomes

$$t = \int dt = \int \frac{d \ln a}{H(a)}$$

- Conformal time becomes

$$\eta = \int \frac{dt}{a} = \int \frac{d \ln a}{aH(a)}$$

- Advantageous since conservation laws give matter evolution with a ; $a = (1 + z)^{-1}$ is a direct observable...

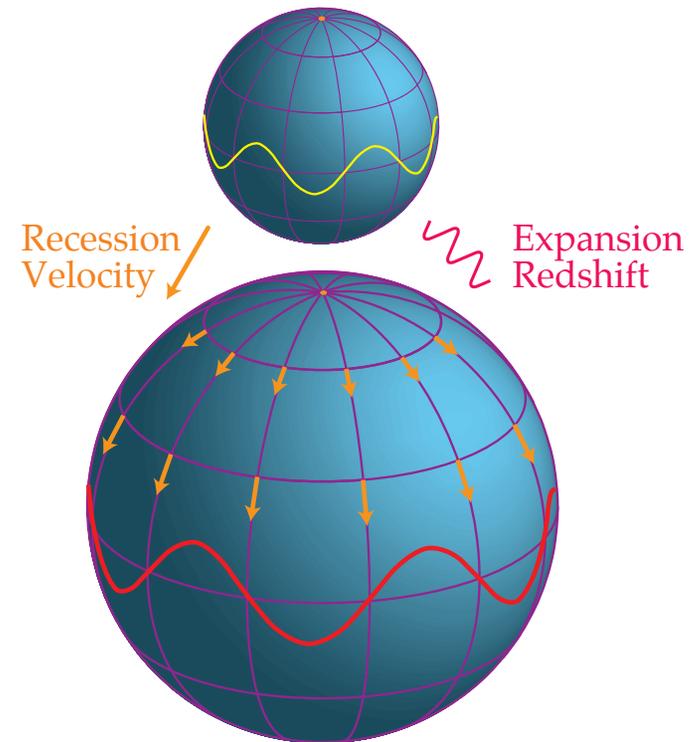
Redshift

- Wavelength of light “stretches” with the scale factor
- The physical wavelength λ_{emit} associated with an observed wavelength today λ_{obs} (or comoving=physical units today) is

$$\frac{\lambda_{\text{emit}}}{\lambda_{\text{obs}}} = \frac{a(t_{\text{emit}})}{a(t_{\text{obs}})} = a(t_{\text{emit}})$$

so that the redshift of spectral lines measures the scale factor of the universe at t , $1 + z = 1/a$.

- Interpreting the redshift as a Doppler shift, objects recede in an expanding universe



Distance-Redshift Relation

- Given **atomically known** rest wavelength λ_{emit} , redshift can be precisely measured from spectra
- Combined with a measure of **distance**, distance-redshift $D(z) \equiv D(z(a))$ can be inferred - given that photons travel $D = \Delta\eta$ this tells us how the scale factor of the universe evolves with time.
- Related to the **expansion history** as

$$D(a) = \int dD = \int_a^1 \frac{d \ln a'}{a' H(a')}$$
$$[d \ln a' = -d \ln(1 + z) = -a' dz]$$
$$D(z) = - \int_z^0 \frac{dz'}{H(z')} = \int_0^z \frac{dz'}{H(z')}$$

Hubble Law

- Note limiting case is the Hubble law

$$\lim_{z \rightarrow 0} D(z) = z/H(z=0) \equiv z/H_0$$

independently of the geometry and expansion dynamics

- Hubble constant usually quoted as dimensionless h

$$H_0 = 100h \text{ km s}^{-1} \text{ Mpc}^{-1}$$

- Observationally $h \sim 0.7$ (see below)
- With $c = 1$, $H_0^{-1} = 9.7778 (h^{-1} \text{ Gyr})$ defines the time scale (Hubble time, \sim age of the universe)
- As well as $H_0^{-1} = 2997.9 (h^{-1} \text{ Mpc})$ a length scale (Hubble scale \sim Horizon scale)

Standard Ruler

- **Standard Ruler**: object of known physical size λ
- Let's apply our rule of thumb: at emission the comoving size is Λ :

$$\lambda = a(t)\Lambda$$

Now everything about light is normal: the object of comoving size Λ subtends an **observed angle** α on the sky α

$$\alpha = \frac{\Lambda}{D_A(z)}$$

- This is the easiest way of thinking about it. But we could also define an effective **physical** distance $d_A(z)$ which corresponds to what we would infer in a non expanding spacetime

$$\alpha \equiv \frac{\lambda}{d_A(z)} = \frac{a\Lambda}{d_A(z)} = \frac{\Lambda}{D_A(z)} \rightarrow d_A(z) = aD_A(z) = \frac{D_A(z)}{1+z}$$

Standard Ruler

- Since $D_A \rightarrow D_A(D_{\text{horizon}})$ whereas $(1+z)$ unbounded, **angular size** of a fixed physical scale at high redshift actually **increases** with **radial distance**
- Paradox: the further away something is, the bigger it appears
 - Easily resolved by thinking about comoving coordinates - a fixed physical scale λ as the universe shrinks and $a \rightarrow 0$ will eventually encompass the whole observable universe out to the horizon in comoving coordinates so of course it subtends a big angle on the sky!
 - But there are no such bound objects in the early universe - there is no causal way such bigger-than-the-horizon objects could form
- Knowing λ or Λ and measuring α and z allows us to infer the comoving angular diameter distance $D_A(z)$

Standard Candle

- **Standard Candle:** objects of same luminosity L , measured flux F
- Apply rules again: at emission in conformal-comoving coordinates
 - L is the energy per unit time at emission
 - Since $E \propto \lambda^{-1}$ and comoving wavelength $\Lambda \propto \lambda/a$ so comoving energy $\mathcal{E} \propto \Lambda^{-1} \propto aE$
 - Per unit time at emission $\Delta t = a\Delta\eta$ in conformal time
 - So observed luminosity today is $\mathcal{L} = \mathcal{E}/\Delta\eta = a^2L$
 - All photons must pass through the sphere at a given distance, so the comoving surface area is $4\pi D_A^2$
- Put this together to the observed flux at $a = 1$

$$F = \frac{\mathcal{L}}{4\pi D_A^2} = \frac{L}{4\pi D_A^2} \frac{1}{(1+z)^2}$$

Notice the flux is diminished by two powers of $(1+z)$

Luminosity Distance

- We can again define a physical “luminosity” distance that corresponds to our non-expanding spacetime intuition

$$F \equiv \frac{L}{4\pi d_L^2}$$

- So luminosity distance

$$d_L = (1 + z)D_A = (1 + z)^2 d_A$$

- As $z \rightarrow 0$, $d_L = d_A = D_A$
- But as $z \rightarrow \infty$, $d_L \gg d_A$ - key to understanding Olber's paradox

Olber's Paradox Redux

- Surface brightness - object of physical size λ

$$S = \frac{F}{\Delta\Omega} = \frac{L}{4\pi d_L^2} \frac{d_A^2}{\lambda^2}$$

- In a non-expanding geometry (regardless of curvature), surface brightness is conserved $d_A = d_L$

$$S = \text{const.}$$

– each sight line in universe full of stars will eventually end on surface of star, night sky should be as bright as sun (not infinite)

- In an expanding universe

$$S \propto (1 + z)^{-4}$$

Olber's Paradox Redux

- Second piece: **age finite** so even if stars exist in the early universe, not all sight lines end on stars
- But even as **age** goes to infinity and the number of sight lines goes to 100%, **surface brightness** of distant objects (of fixed physical size) goes to **zero**
 - Angular size increases
 - Redshift of “luminosity” i.e. energy and arrival time dilation

Measuring $D(z)$

- Astro units side: since flux ratios are very large in cosmology, its more useful to take the log

$$m_1 - m_2 = -2.5 \log_{10}(F_1/F_2)$$

related to d_L **by definition** by inverse square law

$$m_1 - m_2 = 5 \log_{10}[d_L(z_1)/d_L(z_2)]$$

- To quote in terms of a single object, introduce **absolute magnitude** as the magnitude that would be measured for the object at 10 pc

$$m - M = 5 \log_{10}[d_L(z)/10\text{pc}]$$

Knowing absolute magnitude is the same as knowing the absolute distance, otherwise distances are relative

Measuring $D(z)$

- If absolute magnitude unknown, then both standard candles and standard rulers measure relative sizes and fluxes – ironically this means that measuring the change in H is easier than measuring H_0 – acceleration easier than rate!

For **standard candle**, e.g. **compare magnitudes** low z_0 to a high z object - using the Hubble law $d_L(z_0) = z_0/H_0$ we have

$$\Delta m = m_z - m_{z_0} = 5 \log_{10} \frac{d_L(z)}{d_L(z_0)} = 5 \log_{10} \frac{H_0 d_L(z)}{z_0}$$

Likewise for a **standard ruler** comparison at the two redshifts

$$\frac{d_A(z)}{d_A(z_0)} = \frac{H_0 d_A(z)}{z_0}$$

- Distances are measured in units of h^{-1} Mpc.

Measuring $D(z)$

- Since z is a direct observable, in both cases $H_0 D_A(z)$ is the measured quantity
- We can relate that back to $H_0 D(z)$ recalling that

$$H_0 D_A = H_0 R \sin(H_0 D / H_0 R)$$

and use h^{-1} Mpc as the unit for all lengths – furthermore, local observations are at distances much smaller than R so

$H_0 D_A = H_0 D$ is a good approximation

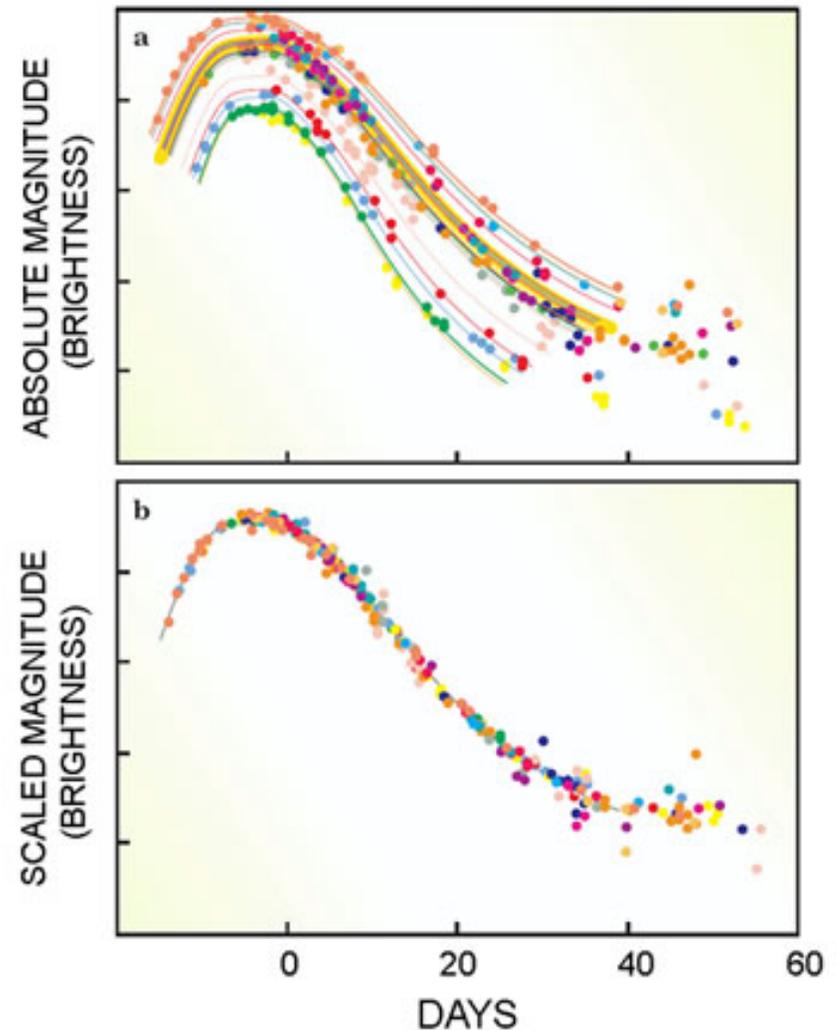
- Then since $D(z) = \int dz / H(z)$ we have

$$H_0 D(z) = \int dz \frac{H_0}{H(z)}$$

- Fundamentally our low to high z comparison tells us the **change** in expansion rate $H(z)/H_0$

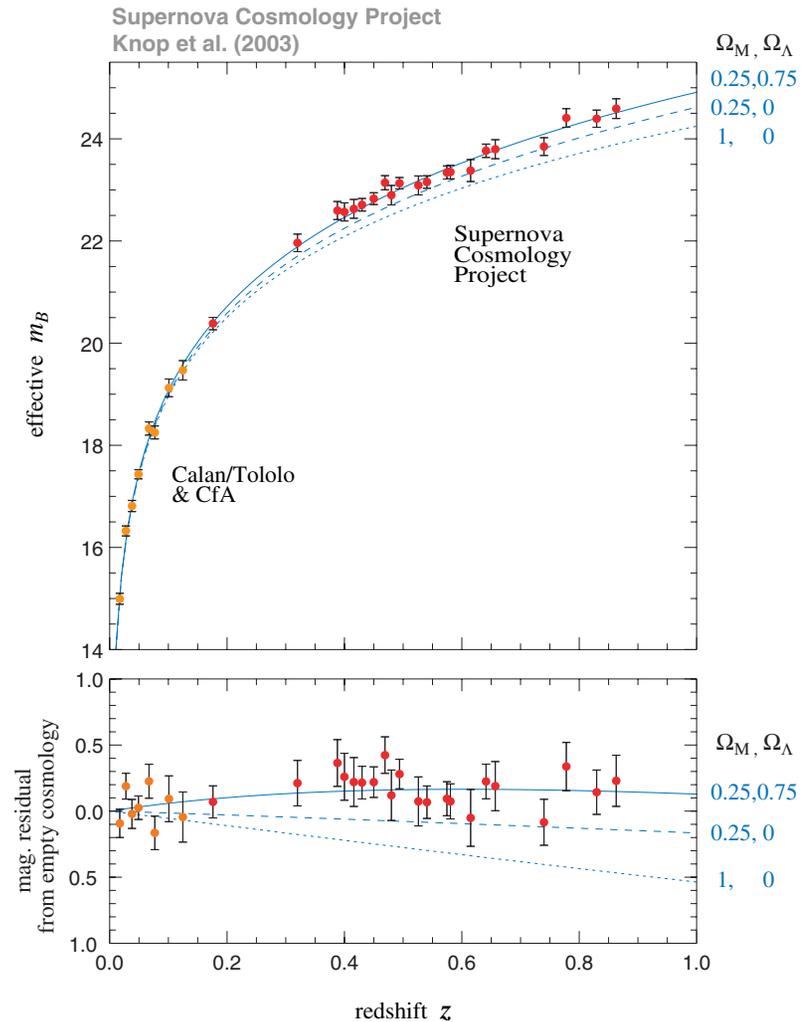
Supernovae as Standard Candles

- Type 1A supernovae are **white dwarfs** that reach **Chandrasekar mass** where electron degeneracy pressure can no longer support the star, hence a **very regular explosion**
- Moreover, the scatter in absolute magnitude is correlated with the **shape** of the light curve - the rate of decline from peak light, empirical “**Phillips relation**”
- Higher ^{56}Ni , **brighter** SN, higher opacity, **longer** light curve duration



Beyond Hubble's Law

- Type 1A are therefore “standardizable” candles leading to a very low scatter $\delta m \sim 0.15$ and visible out to high redshift $z \sim 1$
- Two groups in 1999 found that SN more distant at a given redshift than expected
- Cosmic acceleration



Acceleration of the Expansion

- Using SN as a **relative indicator** (independent of absolute magnitude), comparison of low and high z gives

$$H_0 D(z) = \int dz \frac{H_0}{H}$$

more distant implies that $H(z)$ not increasing at expected rate, i.e. is more constant

- Take the limiting case where $H(z)$ is a **constant** (a.k.a. **de Sitter expansion**)

$$H = \frac{1}{a} \frac{da}{dt} = \text{const}$$

$$\frac{dH}{dt} = \frac{1}{a} \frac{d^2 a}{dt^2} - H^2 = 0$$

$$\frac{1}{a} \frac{d^2 a}{dt^2} = H^2 > 0$$

Acceleration of the Expansion

- Indicates that the **expansion** of the universe is **accelerating**
- Intuition tells us (FRW dynamics shows) **ordinary matter** decelerates expansion since gravity is **attractive**
- **Ordinary expectation** is that

$$H(z > 0) > H_0$$

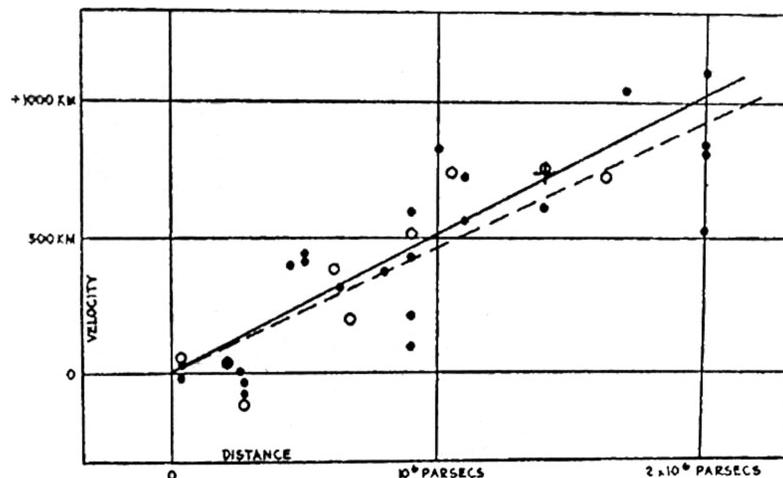
so that the Hubble parameter is higher at high redshift

Or equivalently that **expansion rate decreases** as it **expands**

- Notice that this a purely geometric inference and does not yet say anything about what causes acceleration – topic of next set of lectures on cosmic dynamics

Hubble Constant

- Getting H_0 itself is harder since we need to know the absolute distance d_L to the objects: $H_0 = z_0/d_L$
- Hubble actually inferred **too large** a Hubble constant of $H_0 \sim 500\text{km/s/Mpc}$
- Miscalibration of the Cepheid distance scale - **absolute measurement hard, checkered history**
- Took 70 years to settle on this value with a factor of **2 discrepancy** persisting until late 1990's - which is after the projects which discovered acceleration were conceived!
- H_0 now measured as $73.48 \pm 1.66\text{km/s/Mpc}$ by SHOES calibrating off AGN **water maser**

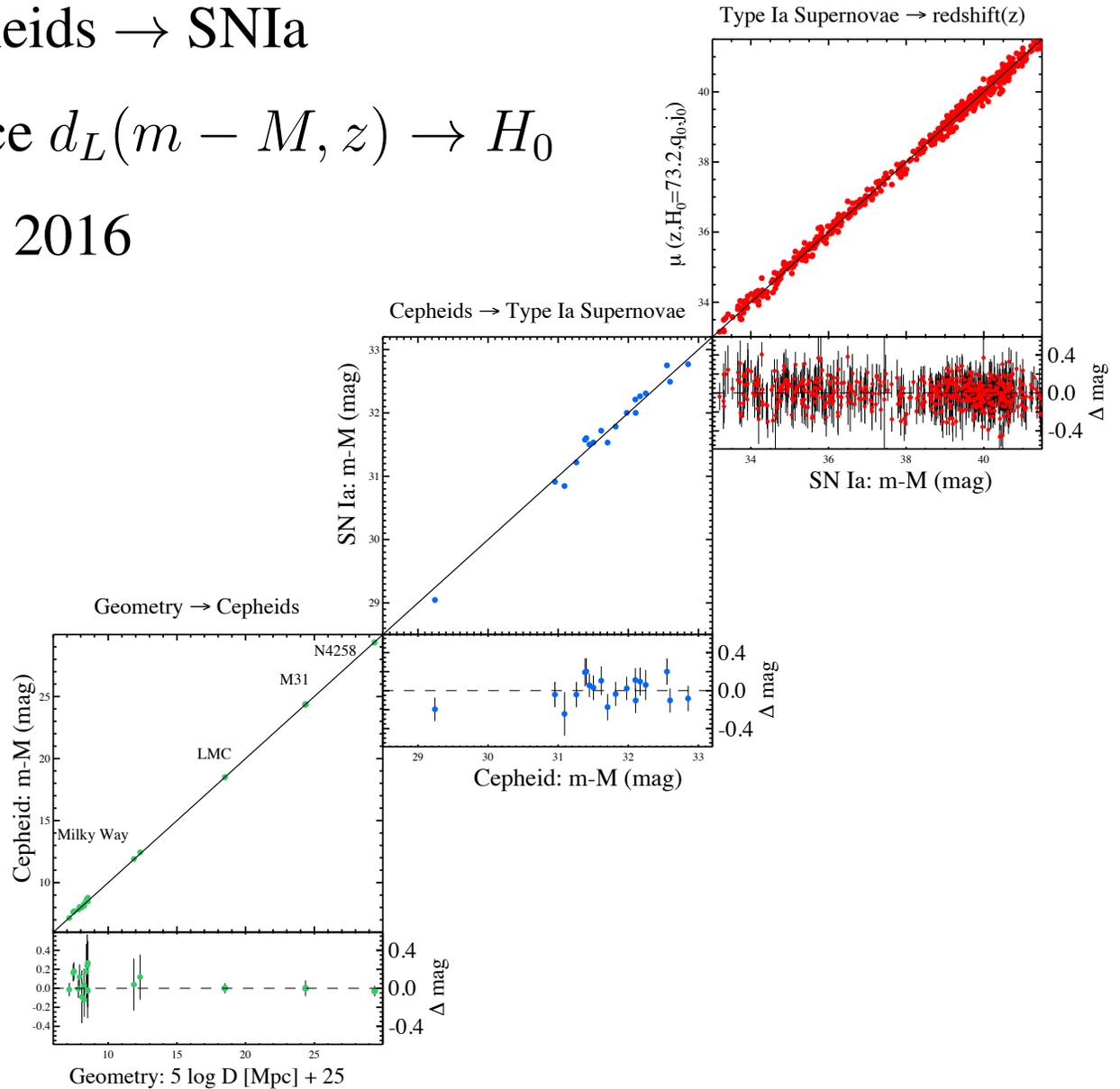


Hubble Constant History

- Difficult measurement since local galaxies have peculiar motions and so their velocity is not entirely due to the “Hubble flow”
- A “distance ladder” of cross calibrated measurements
- Primary distance indicators cepheids, novae planetary nebula, tip of red giant branch, or globular cluster luminosity function, AGN water maser
- Use more luminous secondary distance indications to go out in distance to Hubble flow
 - Tully-Fisher, fundamental plane, surface brightness fluctuations, Type 1A supernova

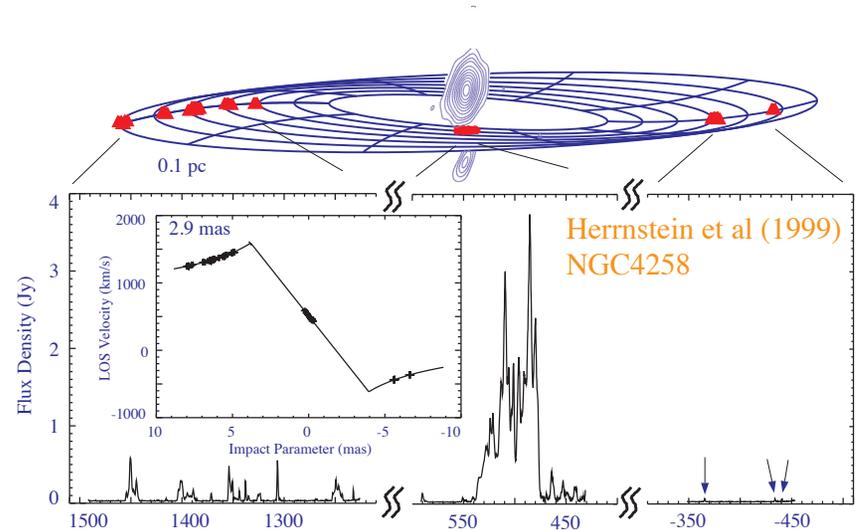
Modern Distance Ladder

- Geometry \rightarrow Cepheids \rightarrow SNIa
- Luminosity distance $d_L(m - M, z) \rightarrow H_0$
- SH0ES, Riess et al 2016



Maser-Cepheid-SN Distance Ladder

- Water maser around AGN, gas in Keplerian orbit
- Measure proper motion, radial velocity, acceleration of orbit
- Method 1: radial velocity plus orbit infer tangential velocity = distance \times angular proper motion



$$v_t = d_A(d\alpha/dt)$$

- Method 2: centripetal acceleration and radial velocity from line infer physical size

$$a = v^2/R, \quad R = d_A\theta$$

Maser-Cepheid-SN Distance Ladder

- Calibrate Cepheid period-luminosity relation in same galaxy
- SHOES project then calibrates SN distance in galaxies with Cepheids
 - Also: consistent with recent HST parallax determinations of 10 galactic Cepheids (8% distance each) with $\sim 20\%$ larger H_0 error bars - normal metallicity as opposed to LMC Cepheids.
- Measure SN at even larger distances out into the Hubble flow
- Riess et al (2019) $H_0 = 74.03 \pm 1.42$ km/s/Mpc more precise (1.9%) than the HST Key Project calibration (11%).
- As of Spring 2019, this differs from the CMB/BAO distance ladder $H_0 = 67.66 \pm 0.42$ working from high redshifts at 4σ