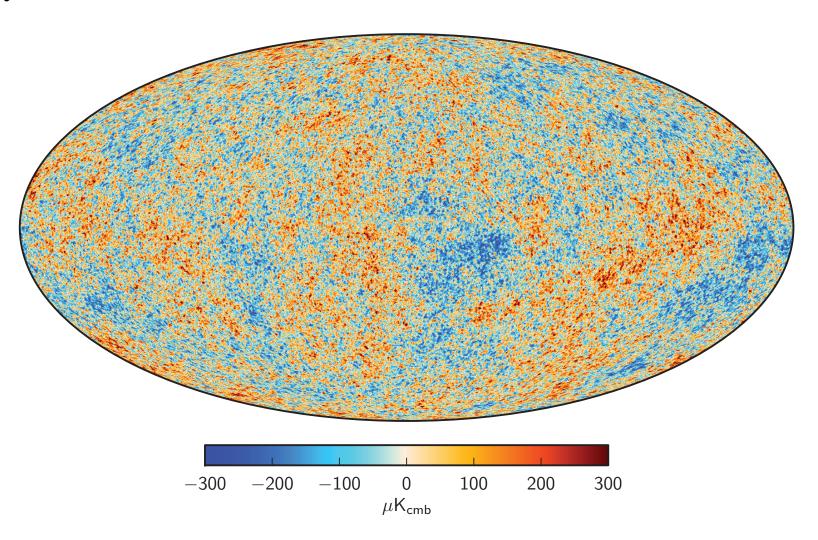
Set 7:

CMB and Large Scale Structure

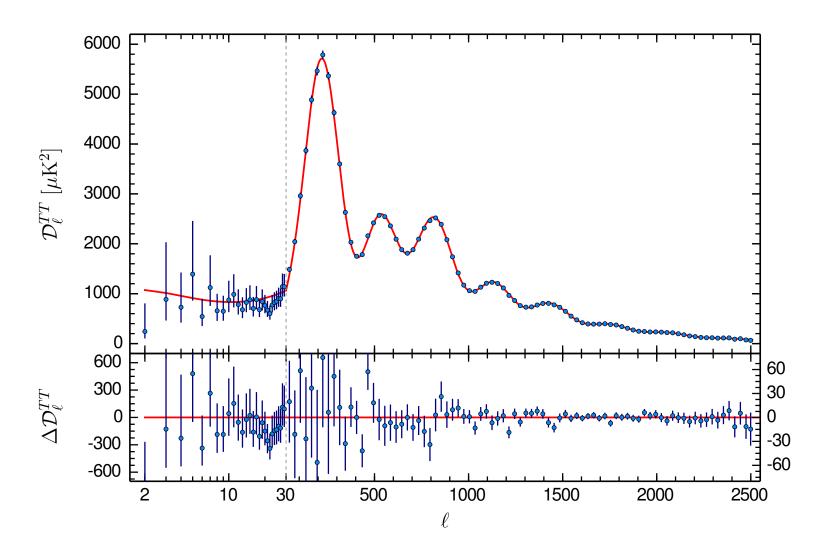
CMB Temperature Anisotropy

• Planck 2015 map of the temperature anisotropy (first discovered by COBE) from recombination:



CMB Temperature Anisotropy

• Power spectrum shows characteristic scales where the intensity of variations peak - reveals geometry and contents of the universe:



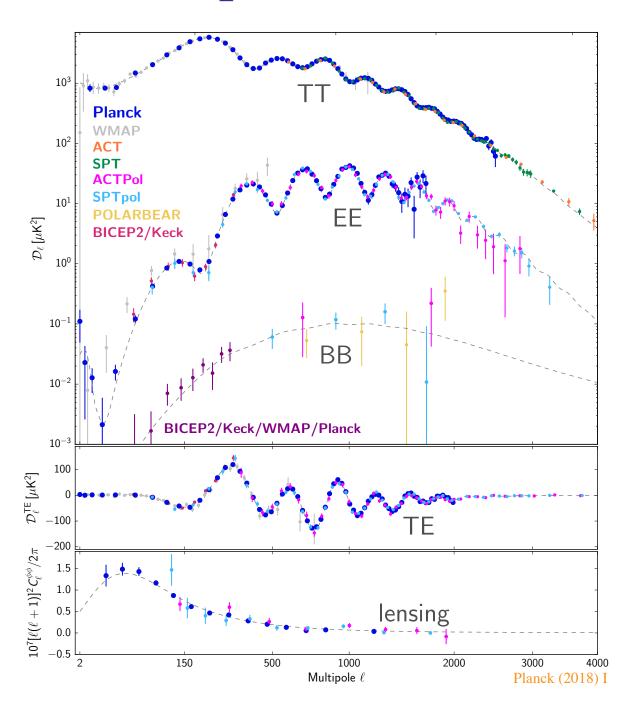
CMB Parameter Inferences

- Spectrum constrains the matter-energy contents of the universe
- Planck 2018 results [arXiv:1807.06209]

Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits	TT,TE,EE+lowE+lensing+BAO 68% limits
$\Omega_b h^2 \ldots \ldots \ldots$	0.02212 ± 0.00022	0.02249 ± 0.00025	0.0240 ± 0.0012	0.02236 ± 0.00015	0.02237 ± 0.00015	0.02242 ± 0.00014
$\Omega_c h^2 \dots \dots \dots$	0.1206 ± 0.0021	0.1177 ± 0.0020	0.1158 ± 0.0046	0.1202 ± 0.0014	0.1200 ± 0.0012	0.11933 ± 0.00091
$100\theta_{MC}$	1.04077 ± 0.00047	1.04139 ± 0.00049	1.03999 ± 0.00089	1.04090 ± 0.00031	1.04092 ± 0.00031	1.04101 ± 0.00029
$\tau \ \dots \dots \dots$	0.0522 ± 0.0080	0.0496 ± 0.0085	0.0527 ± 0.0090	$0.0544^{+0.0070}_{-0.0081}$	0.0544 ± 0.0073	0.0561 ± 0.0071
$ln(10^{10}A_s)\ldots\ldots$	3.040 ± 0.016	$3.018^{+0.020}_{-0.018}$	3.052 ± 0.022	3.045 ± 0.016	3.044 ± 0.014	3.047 ± 0.014
$n_s \dots \dots \dots$	0.9626 ± 0.0057	0.967 ± 0.011	0.980 ± 0.015	0.9649 ± 0.0044	0.9649 ± 0.0042	0.9665 ± 0.0038
$H_0 [km s^{-1} Mpc^{-1}]$	66.88 ± 0.92	68.44 ± 0.91	69.9 ± 2.7	67.27 ± 0.60	67.36 ± 0.54	67.66 ± 0.42

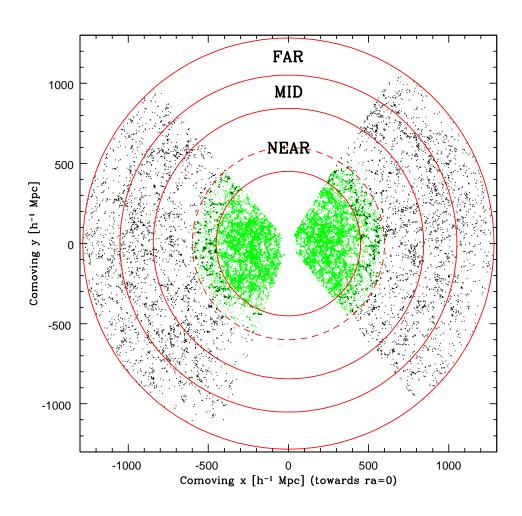
CMB Power Spectra

- Power spectra of CMB
 - temperature
 - polarization
 - lensing



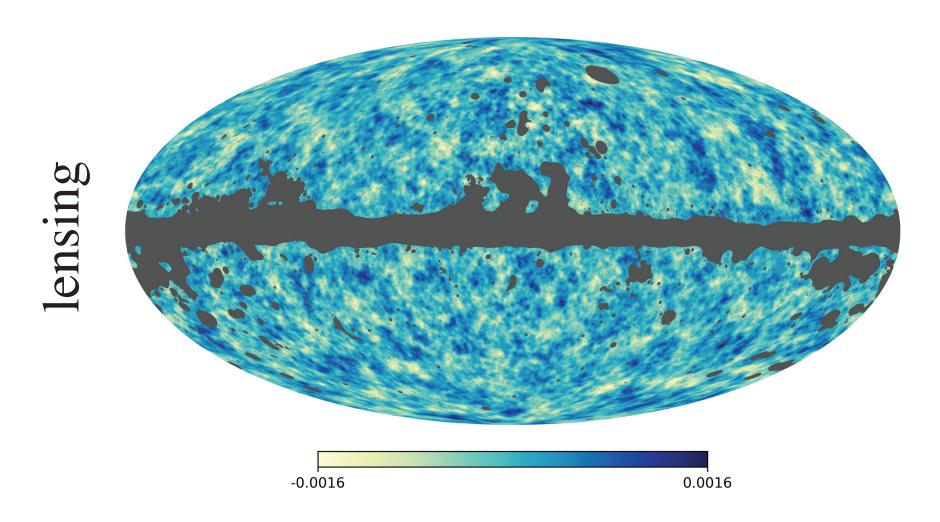
Galaxy Redshift Surveys

• Galaxy redshift surveys (e.g. 2dF and SDSS) measure the three dimensional distribution of galaxies today:



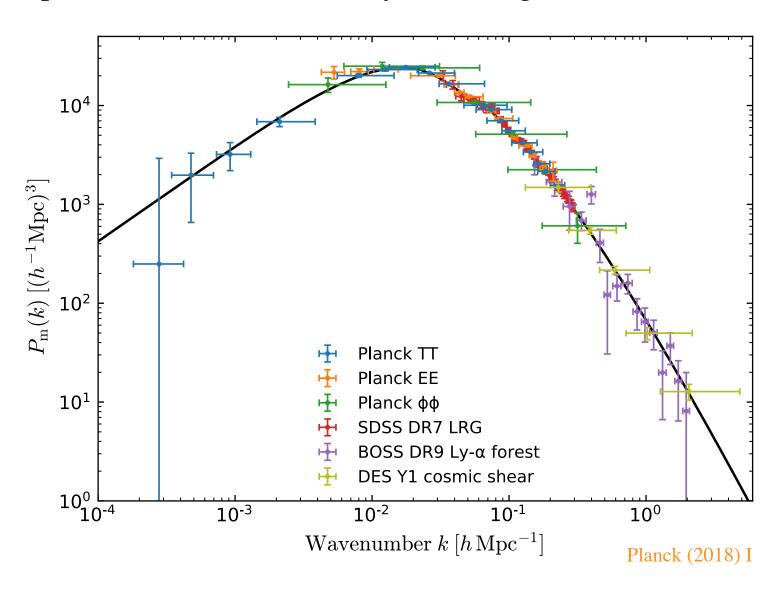
Gravitational Lensing

- Gravitational Lensing measures projected mass
- Planck CMB lensing map



Matter Power Spectrum

• Compilation of Redshift Surveys, Lensing, CMB

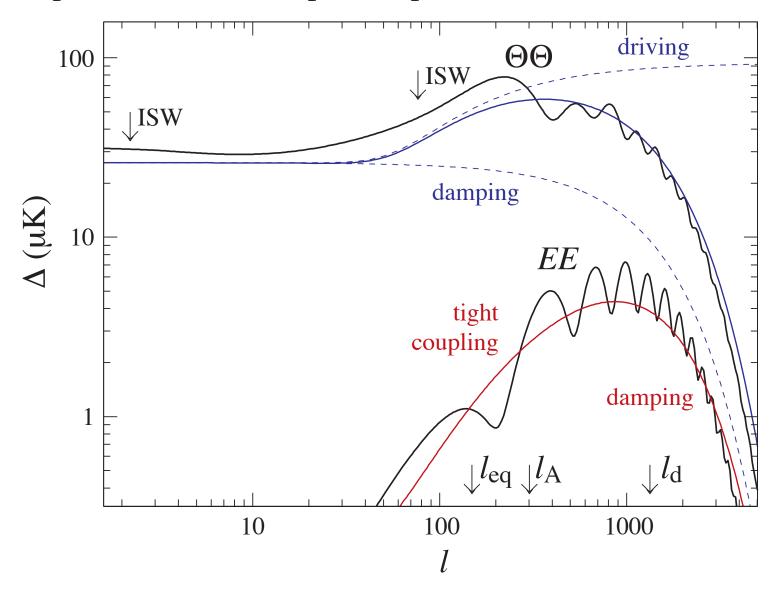


Structure Formation

- Small perturbations from inflation over the course of the 14Gyr life of the universe are gravitationally enhanced into all of the structure seen today
- Cosmic microwave background shows a snapshot at a few hundred thousand years old at recombination
- Discovery in 1992 of cosmic microwave background anisotropy provided the observational breakthrough convincing support for adiabatic initial density fluctuations of amplitude 10^{-5}
- Combine with galaxy clustering large scale structure seen in galaxy surveys right amplitude given cold dark matter

Schematic CMB Spectrum

• Take apart features in the power spectrum



Fluid Approximation

- Thomson scattering of photons and free electrons before recombination is sufficiently rapid that the bayrons and photons are in equilibrium and hence move together
- Mean free path of the photons for $z \approx 10^3$ and $\Omega_b h^2 \approx 0.02$

$$\lambda_C \equiv \frac{1}{n_e \sigma_T a} \sim 2.5 \mathrm{Mpc}$$

small by cosmological standards!

- On scales $\lambda \gg \lambda_C$ photons are tightly coupled to the electrons by Thomson scattering which in turn are tightly coupled to the baryons by Coulomb interactions
- Specifically, their bulk velocities are defined by a single fluid velocity $v_{\gamma}=v_{b}$ and the photons carry no anisotropy in the rest frame of the baryons

Zeroth Order Approximation

- Momentum density of a fluid is $(\rho + p)v$, where p is the pressure
- Neglect the momentum density of the baryons

$$R \equiv \frac{(\rho_b + p_b)v_b}{(\rho_\gamma + p_\gamma)v_\gamma} = \frac{\rho_b + p_b}{\rho_\gamma + p_\gamma} = \frac{3\rho_b}{4\rho_\gamma}$$
$$\approx 0.6 \left(\frac{\Omega_b h^2}{0.02}\right) \left(\frac{a}{10^{-3}}\right)$$

since $\rho_{\gamma} \propto T^4$ is fixed by the CMB temperature $T=2.73(1+z){\rm K}$

- OK substantially before recombination
- Neglect radiation in the expansion

$$\frac{\rho_m}{\rho_r} = 3.6 \left(\frac{\Omega_m h^2}{0.15}\right) \left(\frac{a}{10^{-3}}\right)$$

Neglect gravity

Fluid Equations

• Density $\rho_{\gamma} \propto T^4$ so define temperature fluctuation Θ

$$\delta_{\gamma} = 4 \frac{\delta T}{T} \equiv 4\Theta$$

- Real space continuity eqn.: the local number or energy density of photons changes if there is a divergence of the velocity field a flow inwards or outwards from the volume
- Transformed to Fourier space $\nabla(e^{i\mathbf{k}\cdot\mathbf{x}}) \to i\mathbf{k}(e^{i\mathbf{k}\cdot\mathbf{x}})$ and $\nabla \cdot \mathbf{v} = -kv$

$$\dot{\delta}_{\gamma} = -(1 + w_{\gamma})kv_{\gamma}$$

$$\dot{\Theta} = -\frac{1}{3}kv_{\gamma}$$

Fluid Equations

• Euler equation (neglecting gravity for now): momentum conservation says that pressure gradients generate changes in momentum density $k\delta p_{\gamma}=kc_{s}^{2}\delta\rho_{\gamma}$

$$\dot{v}_{\gamma} = \frac{kc_s^2}{1 + w_{\gamma}} \delta_{\gamma}$$
$$= kc_s^2 \frac{3}{4} \delta_{\gamma} = 3c_s^2 k\Theta$$

where the sound speed $c_s^2 = \delta p/\delta \rho$ is the pressure response to a density fluctuation

 So if you squeeze the photon gas to raise its density, its going to respond with a restoring force by raising the pressure and resisting compression → acoustic oscillations

Oscillator: Take One

Combine these to form the simple harmonic oscillator equation

$$\ddot{\Theta} + c_s^2 k^2 \Theta = 0$$

where the sound speed is adiabatic

$$c_s^2 = \frac{\delta p}{\delta \rho} = \frac{\dot{p}_{\gamma}}{\dot{\rho}_{\gamma}}$$

here $c_s^2 = 1/3$ since we are photon-dominated

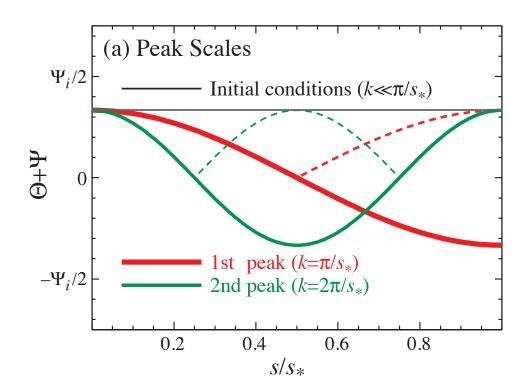
• General solution:

$$\Theta(\eta) = \Theta(0)\cos(ks) + \frac{\dot{\Theta}(0)}{kc_s}\sin(ks)$$

where the sound horizon is defined as $s \equiv \int c_s d\eta$

Harmonic Extrema

- All modes begin at end of inflation and are frozen in at recombination (denoted with a subscript *)
- Temperature perturbations of different amplitude for different modes.



• For the adiabatic (curvature mode) initial conditions

$$\dot{\Theta}(0) = 0$$

So solution

$$\Theta(\eta_*) = \Theta(0)\cos(ks_*)$$

Harmonic Extrema

 Modes caught in the extrema of their oscillation will have enhanced fluctuations

$$k_n s_* = n\pi$$

yielding a fundamental scale or frequency, related to the inverse sound horizon

$$k_A = \pi/s_*$$

and a harmonic relationship to the other extrema as 1:2:3...

Peak Location

• The fundmental physical scale is translated into a fundamental angular scale by simple projection according to the angular diameter distance D_A

$$\theta_A = \lambda_A / D_A$$

$$\ell_A = k_A D_A$$

• In a flat universe, the distance is simply $D_A = D \equiv \eta_0 - \eta_* \approx \eta_0$, the horizon distance, and $k_A = \pi/s_* = \sqrt{3}\pi/\eta_*$ so

$$\theta_A pprox rac{\eta_*}{\eta_0}$$

• In a matter-dominated universe $\eta \propto a^{1/2}$ so $\theta_A \approx 1/30 \approx 2^\circ$ or

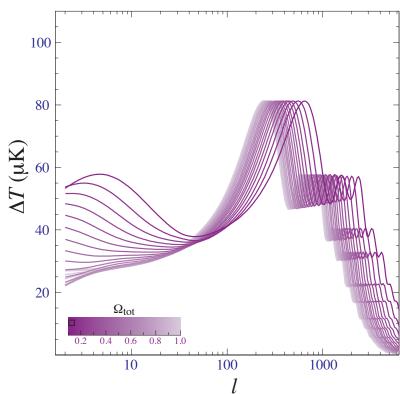
$$\ell_A \approx 200$$

Curvature

- In a curved universe, the apparent or angular diameter distance is no longer the conformal distance $D_A = R \sin(D/R) \neq D$
- Objects in a closed universe are further than they appear! gravitational lensing of the background...
- Curvature scale of the universe must be substantially larger than current horizon

Curvature

- Flat universe indicates critical density and implies missing energy given local measures of the matter density "dark energy"
- D also depends on dark energy density $\Omega_{\rm DE}$ and equation of state $w=p_{\rm DE}/\rho_{\rm DE}$.
- Expansion rate at recombination or matter-radiation ratio enters into calculation of k_A .



Restoring Gravity

- Take a simple photon dominated system with gravity
- Continuity altered since a gravitational potential represents a stretching of the spatial fabric that dilutes number densities – formally a spatial curvature perturbation
- Think of this as a perturbation to the scale factor $a \to a(1 + \Phi)$ so that the cosmogical redshift is generalized to

$$\frac{\dot{a}}{a} \rightarrow \frac{\dot{a}}{a} + \dot{\Phi}$$

so that the continuity equation becomes

$$\dot{\Theta} = -\frac{1}{3}kv_{\gamma} - \dot{\Phi}$$

Restoring Gravity

• Gravitational force in momentum conservation ${\bf F}=-m\nabla\Psi$ generalized to momentum density modifies the Euler equation to

$$\dot{v}_{\gamma} = k(\Theta + \Psi)$$

- General relativity says that Φ and Ψ are the relativistic analogues of the Newtonian potential and that $\Phi \approx -\Psi$.
- In our matter-dominated approximation, Φ represents matter density fluctuations through the cosmological Poisson equation

$$k^2 \Phi = 4\pi G a^2 \rho_m \Delta_m$$

where the difference comes from the use of comoving coordinates for k (a^2 factor), the removal of the background density into the background expansion ($\rho\Delta_m$) and finally a coordinate subtlety that enters into the definition of Δ_m

Constant Potentials

- In the matter dominated epoch potentials are constant because infall generates velocities as $v_m \sim k \eta \Psi$
- Velocity divergence generates density perturbations as $\Delta_m \sim -k\eta v_m \sim -(k\eta)^2 \Psi$
- And density perturbations generate potential fluctuations

$$\Phi = \frac{4\pi G a^2 \rho \Delta}{k^2} \approx \frac{3}{2} \frac{H^2 a^2}{k^2} \Delta \sim \frac{\Delta}{(k\eta)^2} \sim -\Psi$$

keeping them constant. Note that because of the expansion, density perturbations must grow to keep potentials constant.

Constant Potentials

- More generally, if stress perturbations are negligible compared with density perturbations ($\delta p \ll \delta \rho$) then potential will remain roughly constant
- More specifically a variant called the Bardeen or comoving curvature is strictly constant

$$\mathcal{R} = \text{const} \approx \frac{5 + 3w}{3 + 3w} \Phi$$

where the approximation holds when $w \approx \text{const.}$

Oscillator: Take Two

• Combine these to form the simple harmonic oscillator equation

$$\ddot{\Theta} + c_s^2 k^2 \Theta = -\frac{k^2}{3} \Psi - \ddot{\Phi}$$

• In a CDM dominated expansion $\dot{\Phi} = \dot{\Psi} = 0$. Also for photon domination $c_s^2 = 1/3$ so the oscillator equation becomes

$$\ddot{\Theta} + \ddot{\Psi} + c_s^2 k^2 (\Theta + \Psi) = 0$$

Solution is just an offset version of the original

$$[\Theta + \Psi](\eta) = [\Theta + \Psi](0) \cos(ks)$$

ullet $\Theta+\Psi$ is also the observed temperature fluctuation since photons lose energy climbing out of gravitational potentials at recombination

Effective Temperature

- Photons climb out of potential wells at last scattering
- Lose energy to gravitational redshifts
- Observed or effective temperature

$$\Theta + \Psi$$

- Effective temperature oscillates around zero with amplitude given by the initial conditions
- Note: initial conditions are set when the perturbation is outside of horizon, need inflation or other modification to matter-radiation FRW universe.
- GR says that initial temperature is given by initial potential

Sachs-Wolfe Effect and the Magic 1/3

• A gravitational potential is a perturbation to the temporal coordinate [formally a gauge transformation]

$$\frac{\delta t}{t} = \Psi$$

• Convert this to a perturbation in the scale factor, in a matter dominated expansion $a \propto t^{2/3}$ so

$$\frac{\delta a}{a} = \frac{2}{3} \frac{\delta t}{t}$$

• CMB temperature is cooling as $T \propto a^{-1}$ so

$$\Theta + \Psi \equiv \frac{\delta T}{T} + \Psi = -\frac{\delta a}{a} + \Psi = \frac{1}{3}\Psi$$

Sachs-Wolfe Normalization

- Use measurements of $\Delta T/T \approx 10^{-5}$ in the Sachs-Wolfe effect to infer $\Delta_{\mathcal{R}}^2$
- Recall in matter domination $\Psi = -3\mathcal{R}/5$ and so $\Delta T/T = -\mathcal{R}/5$
- So that the amplitude of initial curvature fluctuations is $\Delta_R \approx 5 \times 10^{-5}$
- This then determines the amplitude of the inflationary power spectrum $A_S = \Delta_R^2$ in the previous lecture set

Baryon Loading

- Baryons add extra mass to the photon-baryon fluid
- Controlling parameter is the momentum density ratio:

$$R \equiv \frac{p_b + \rho_b}{p_\gamma + \rho_\gamma} \approx 30\Omega_b h^2 \left(\frac{a}{10^{-3}}\right)$$

of order unity at recombination

Momentum density of the joint system is conserved

$$(\rho_{\gamma} + p_{\gamma})v_{\gamma} + (\rho_b + p_b)v_b \approx (p_{\gamma} + p_{\gamma} + \rho_b + \rho_{\gamma})v_{\gamma}$$
$$= (1 + R)(\rho_{\gamma} + p_{\gamma})v_{\gamma b}$$

Momentum density ratio enters as

$$[(1+R)v_{\gamma b}] = k\Theta + (1+R)k\Psi$$

New Euler Equation

Photon continuity remains the same

$$\dot{\Theta} = -\frac{k}{3}v_{\gamma b} - \dot{\Phi}$$

Modification of oscillator equation

$$\frac{d}{d\eta}[(1+R)\dot{\Theta}] + \frac{1}{3}k^2\Theta = -\frac{1}{3}k^2(1+R)\Psi - \frac{d}{d\eta}[(1+R)\dot{\Phi}]$$

• In a CDM dominated expansion $\dot{\Phi} = \dot{\Psi} = 0$ and the adiabatic approximation where the sound speed evolves slowly

$$c_s = \sqrt{\frac{1}{3} \frac{1}{1+R}}$$

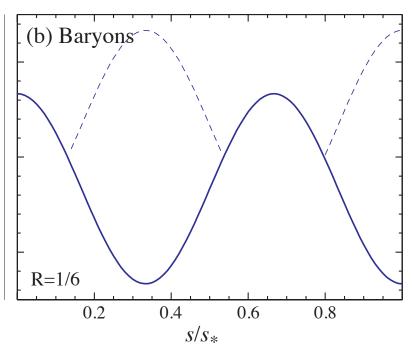
$$[\Theta + (1 + \mathbf{R})\Psi](\eta) = [\Theta + (1 + \mathbf{R})\Psi](0)\cos(k\mathbf{s})$$

Baryon Peak Phenomenology

- Photon-baryon ratio enters in three ways
- Overall larger amplitude:

$$[\Theta + (1 + R)\Psi](0) = \frac{1}{3}(1 + 3R)\Psi(0)$$

• Even-odd peak modulation of effective temperature



$$[\Theta + \Psi]_{\text{peaks}} = [\pm (1 + 3R) - 3R] \frac{1}{3} \Psi(0)$$
$$[\Theta + \Psi]_1 - [\Theta + \Psi]_2 = [-6R] \frac{1}{3} \Psi(0)$$

• Shifting of the sound horizon down or ℓ_A up

$$\ell_A \propto \sqrt{1+R}$$

Photon Baryon Ratio Evolution

- Actual effects smaller since R evolves
- Oscillator equation has time evolving mass

$$c_s^2 \frac{d}{d\eta} (c_s^{-2} \dot{\Theta}) + c_s^2 k^2 \Theta = 0$$

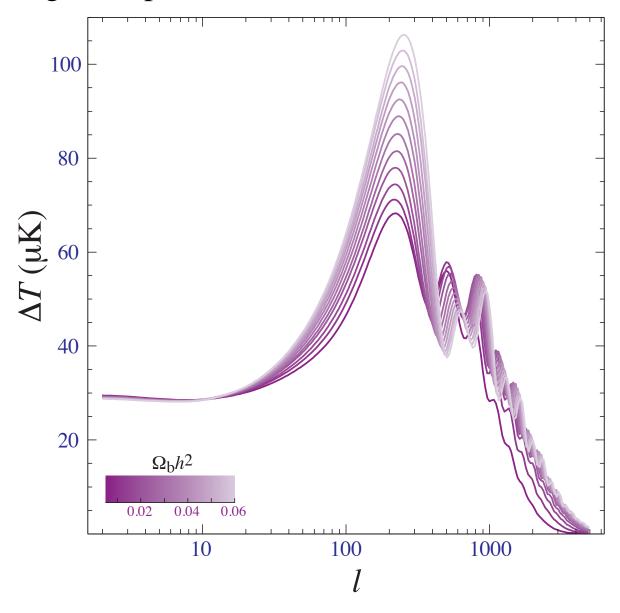
- Effective mass is is $m_{\text{eff}} = 3c_s^{-2} = (1 + R)$
- Adiabatic invariant

$$\frac{E}{\omega} = \frac{1}{2} m_{\text{eff}} \omega A^2 = \frac{1}{2} 3 c_s^{-2} k c_s A^2 \propto A^2 (1 + R)^{1/2} = const.$$

• Amplitude of oscillation $A \propto (1+R)^{-1/4}$ decays adiabatically as the photon-baryon ratio changes

Baryons in the Power Spectrum

• Relative heights of peaks



Oscillator: Take Three and a Half

• The not-quite-so simple harmonic oscillator equation is a forced harmonic oscillator

$$c_s^2 \frac{d}{d\eta} (c_s^{-2} \dot{\Theta}) + c_s^2 k^2 \Theta = -\frac{k^2}{3} \Psi - c_s^2 \frac{d}{d\eta} (c_s^{-2} \dot{\Phi})$$

changes in the gravitational potentials alter the form of the acoustic oscillations

- If the forcing term has a temporal structure that is related to the frequency of the oscillation, this becomes a driven harmonic oscillator
- ullet Term involving Ψ is the ordinary gravitational force
- Term involving Φ involves the $\dot{\Phi}$ term in the continuity equation as a (curvature) perturbation to the scale factor

Potential Decay

Matter-to-radiation ratio

$$\frac{\rho_m}{\rho_r} \approx 24\Omega_m h^2 \left(\frac{a}{10^{-3}}\right)$$

of order unity at recombination in a low Ω_m universe

Radiation is not stress free and so impedes the growth of structure

$$k^2\Phi = 4\pi G a^2 \rho_r \Delta_r$$

 $\Delta_r \sim 4\Theta$ oscillates around a constant value, $\rho_r \propto a^{-4}$ so the Netwonian curvature decays.

• General rule: potential decays if the dominant energy component has substantial stress fluctuations, i.e. below the generalized sound horizon or Jeans scale

Radiation Driving

 Decay is timed precisely to drive the oscillator - close to fully coherent

$$|[\Theta + \Psi](\eta)| = |[\Theta + \Psi](0) + \Delta \Psi - \Delta \Phi|$$

$$= |\frac{1}{3}\Psi(0) - 2\Psi(0)| = |\frac{5}{3}\Psi(0)|$$

$$\Psi_{i}$$

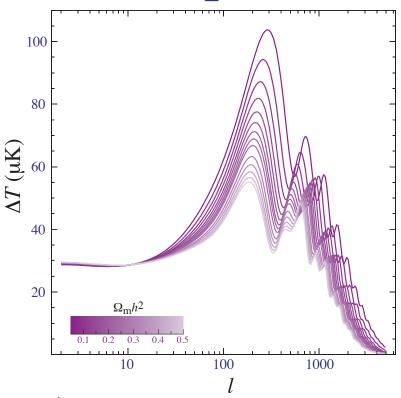
$$\Psi_{i}$$

$$\Phi_{i}$$

• 5× the amplitude of the Sachs-Wolfe effect!

Matter-Radiation in the Power Spectrum

- Coherent approximation is exact for a photon-baryon fluid but reality is reduced to $\sim 4\times$ because neutrino contribution is free streaming not fluid like
- Neutrinos drive the oscillator less efficiently and also slightly change the phase of the oscillation



- Actual initial conditions are $\Theta + \Psi = \Psi/2$ for radiation domination but comparison to matter dominated SW correct
- With 3 peaks, it is possible to solve for both the baryons and dark matter densities, providing a calibration for the sound horizon
- Higher peaks check consistency with assumptions: e.g. extra relativistic d.o.f.s

Damping

- Tight coupling equations assume a perfect fluid: no viscosity, no heat conduction
- Fluid imperfections are related to the mean free path of the photons in the baryons

$$\lambda_C = \dot{\tau}^{-1}$$
 where $\dot{\tau} = n_e \sigma_T a$

is the conformal opacity to Thomson scattering

• Dissipation is related to the diffusion length: random walk approximation

$$\lambda_D = \sqrt{N}\lambda_C = \sqrt{\eta/\lambda_C}\,\lambda_C = \sqrt{\eta\lambda_C}$$

the geometric mean between the horizon and mean free path

• $\lambda_D/\eta_* \sim$ few %, so expect the peaks :> 3 to be affected by dissipation

Near Perfection in 6 Numbers

 All this precision data described by
 6 ΛCDM parameters

 $-\Omega_c h^2$: CDM

 $-\Omega_b h^2$: baryons

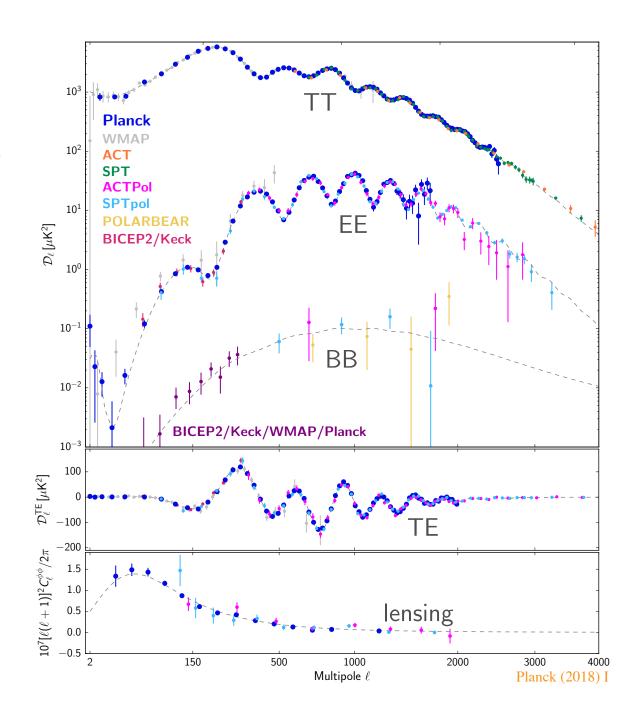
 $-\theta_s$: sound scale

 $-A_s$: amplitude

 $-n_s$: tilt

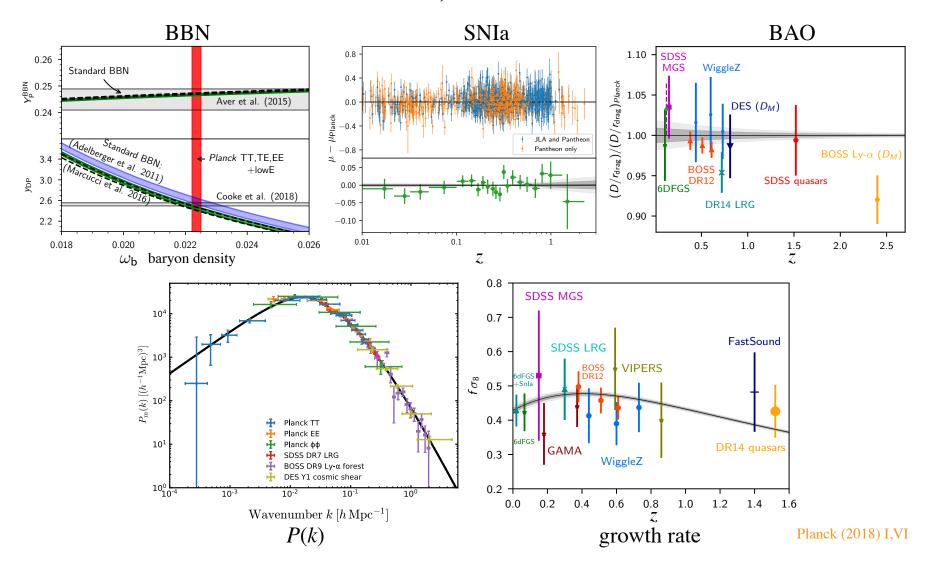
 $-\tau$: reionization

Measured
 to sub percent
 precision (except τ)



Predictive Power

• Predicts all other observables, which direct measurements test



• Good agreement, even weak lensing, clusters, and yes H_0 (< 10%)

Transfer Function

• Transfer function transfers the initial Newtonian curvature to its value today (linear response theory)

$$T(k) = \frac{\Phi(k, a = 1)}{\Phi(k, a_{\text{init}})} \frac{\Phi(k_{\text{norm}}, a_{\text{init}})}{\Phi(k_{\text{norm}}, a = 1)}$$

- Conservation of Bardeen curvature: Newtonian curvature is a constant when stress perturbations are negligible: above the horizon during radiation and dark energy domination, on all scales during matter domination
- When stress fluctuations dominate, perturbations are stabilized by the Jeans mechanism
- Hybrid Poisson equation: Newtonian curvature, comoving density perturbation $\Delta \equiv (\delta \rho/\rho)_{\rm com}$ implies Φ decays

$$(k^2 - 3K)\Phi = 4\pi Ga^2 \rho \Delta \sim \eta^{-2} \Delta$$

Transfer Function

• Freezing of Δ stops at $\eta_{\rm eq}$

$$\Phi \sim (k\eta_{\rm eq})^{-2}\Delta_H \sim (k\eta_{\rm eq})^{-2}\Phi_{\rm init}$$

• Transfer function has a k^{-2} fall-off beyond $k_{\rm eq} \sim \eta_{\rm eq}^{-1}$

$$\eta_{\text{eq}} = 15.7(\Omega_m h^2)^{-1} \left(\frac{T}{2.7K}\right)^2 \text{Mpc}$$

- Small correction since growth with a smooth radiation component is logarithmic not frozen
- Transfer function is a direct output of an Einstein-Boltzmann code

Fitting Function

• Alternately accurate fitting formula exist, e.g. pure CDM form:

$$T(k(q)) = \frac{L(q)}{L(q) + C(q)q^2}$$

$$L(q) = \ln(e + 1.84q)$$

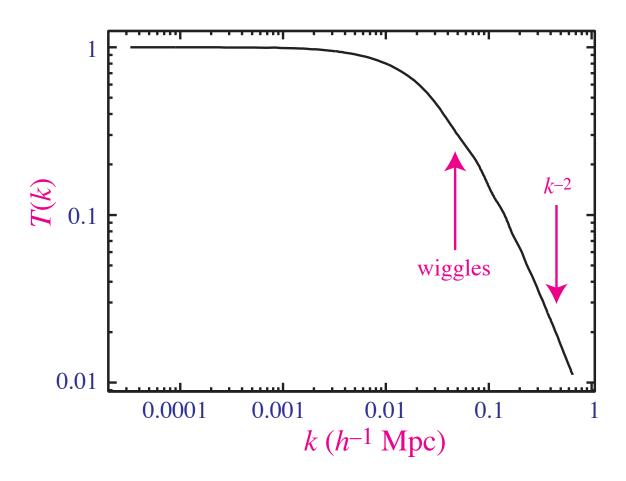
$$C(q) = 14.4 + \frac{325}{1 + 60.5q^{1.11}}$$

$$q = k/\Omega_m h^2 \text{Mpc}^{-1} (T_{\text{CMB}}/2.7K)^2$$

• In h Mpc⁻¹, the critical scale depends on $\Gamma \equiv \Omega_m h$ also known as the shape parameter

Transfer Function

• Numerical calculation

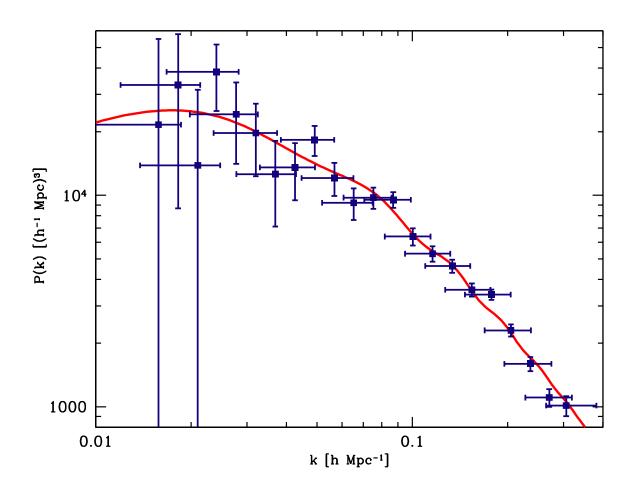


Baryon Acoustic Oscillations

- Baryons caught up in the acoustic oscillations of the CMB and impart acoustic oscillations to the transfer function. Density enhancements are produced kinematically through the continuity equation $\delta_b \sim (k\eta)v_b$ and hence are out of phase with CMB temperature peaks
- Dissipation of the acoustic oscillations eliminates both the CMB and baryon perturbations known as Silk damping for the baryons. This suppression and the general fact that baryons are caught up with photons was one of the main arguments for CDM
- Detected first in the SDSS LRG survey.
- An excellent standard ruler for angular diameter distance $D_A(z)$ since it does not evolve in redshift in linear theory
- Radial extent of BAO gives H(z)

Power Spectrum

• SDSS data



• Power spectrum defines large scale structure observables: galaxy clustering, velocity field, Ly α forest clustering, cosmic shear