

Set 10:

Synchrotron

Lorentz Force

- Acceleration of electrons due to the magnetic field gives rise to synchrotron radiation
- Lorentz force

$$\frac{dP^\mu}{d\tau} = \frac{e}{c} F^\mu{}_\nu U^\nu, \quad F^\mu{}_\nu = \begin{pmatrix} 0 & E_x & E_y & E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{pmatrix}$$

$$\frac{d}{d\tau} \gamma m \begin{pmatrix} c \\ v_x \\ v_y \\ v_z \end{pmatrix} = \frac{e}{c} \begin{pmatrix} 0 & E_x & E_y & E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{pmatrix} \gamma \begin{pmatrix} c \\ v_x \\ v_y \\ v_z \end{pmatrix}$$

Lorentz Force

$$\frac{d}{d\tau} \gamma m \begin{pmatrix} c \\ v_x \\ v_y \\ v_z \end{pmatrix} = \frac{e}{c} \gamma \begin{pmatrix} \mathbf{E} \cdot \mathbf{v} \\ E_x c + B_z v_y - B_y v_z \\ E_y c - B_z v_x + B_x v_z \\ E_z c + B_y v_x - B_x v_y \end{pmatrix}$$

$$\frac{d}{d\tau} \gamma m \begin{pmatrix} c \\ \mathbf{v} \end{pmatrix} = \frac{e}{c} \gamma \begin{pmatrix} \mathbf{E} \cdot \mathbf{v} \\ \mathbf{E} c + \mathbf{v} \times \mathbf{B} \end{pmatrix} \quad [d\tau = dt/\gamma]$$

$$\frac{d}{dt} \gamma m \begin{pmatrix} c \\ \mathbf{v} \end{pmatrix} = \frac{e}{c} \begin{pmatrix} \mathbf{E} \cdot \mathbf{v} \\ \mathbf{E} c + \mathbf{v} \times \mathbf{B} \end{pmatrix}$$

Energy-Momentum Equations

- Energy and momentum equations

$$\frac{d}{dt}\gamma mc^2 = e\mathbf{E} \cdot \mathbf{v}$$

$$\frac{d}{dt}\gamma m\mathbf{v} = e\left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}\right)$$

- Consider a B field with no E field

$$\frac{d}{dt}\gamma mc^2 = 0, \quad \frac{d}{dt}\gamma m\mathbf{v} = e\frac{\mathbf{v}}{c} \times \mathbf{B}$$

where the former ignores radiation energy losses

- Then the velocity is constant

$$\frac{d\gamma}{dt} = 0 \quad \rightarrow \gamma = \text{const} \quad \rightarrow |v| = \text{const}$$

Cyclical Motion

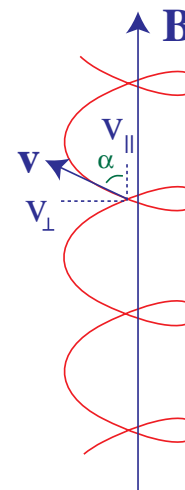
- Define components $\mathbf{v}_{\parallel} \parallel \mathbf{B}$,
 $\mathbf{v}_{\perp} \perp \mathbf{B}$. Momentum equation becomes

$$\frac{d\mathbf{v}_{\parallel}}{dt} = 0 \quad \frac{d\mathbf{v}_{\perp}}{dt} = \frac{e}{\gamma mc} \mathbf{v}_{\perp} \times \mathbf{B}$$

so that $v_{\parallel} = \text{const.}$ $v^2 = v_{\parallel}^2 + v_{\perp}^2 = \text{const.}$
and thus $v_{\perp} = \text{const.}$

- Acceleration is orthogonal to v_{\perp} with

$$|a_{\perp}| = \frac{e}{\gamma mc} |v_{\perp}| B$$



Power

- Equation of motion solved by

$$\mathbf{v}_{\perp}(t) = v_{\perp} \begin{pmatrix} \sin(\omega_B t + \delta) \\ \cos(\omega_B t + \delta) \end{pmatrix}, \quad \omega_B = \frac{eB}{\gamma mc}$$

- Radiated power

$$P = \frac{2e^2}{3c^3} \gamma^4 a_{\perp}^2 = \frac{2}{3} \frac{e^4 v_{\perp}^2 B^2}{m^2 c^5} \gamma^2 = \frac{1}{4\pi} \frac{v_{\perp}^2 B^2}{c} \gamma^2 \sigma_T, \quad \left[\sigma_T = \frac{8\pi}{3} \frac{e^4}{m^2 c^4} \right]$$

- Synchrotron emission can be viewed as Compton scattering off of virtual B photons. Defining the pitch angle α between \mathbf{B} and \mathbf{v} ,

$$v_{\perp} = v \sin \alpha \quad \langle \sin^2 \alpha \rangle = \int \sin^2 \alpha \frac{d\Omega}{4\pi} = \frac{2}{3}$$

$$\langle P_{\text{synch}} \rangle = \frac{1}{6\pi} \frac{v^2 B^2}{c} \gamma^2 \sigma_T$$

Virtual Photons

- The energy density associated with the magnetic field is

$$u_B = B^2/8\pi$$

$$\langle P_{\text{synch}} \rangle = \frac{4}{3} \sigma_T \frac{v^2}{c} \gamma^2 u_B = \frac{4}{3} \sigma_T c \beta^2 \gamma^2 u_B$$

- Power in Compton scattering

$$\langle P_{\text{comp}} \rangle = \frac{4}{3} \sigma_T c \beta^2 \gamma^2 u_\gamma$$

so that synchrotron power = Compton power with incident energy density replaced by the magnetic field energy density

$$\frac{\langle P_{\text{synch}} \rangle}{\langle P_{\text{comp}} \rangle} = \frac{u_B}{u_\gamma}$$

- Synchrotron radiation produces photons for inverse Compton scattering - removes energy from electrons - self regulation process

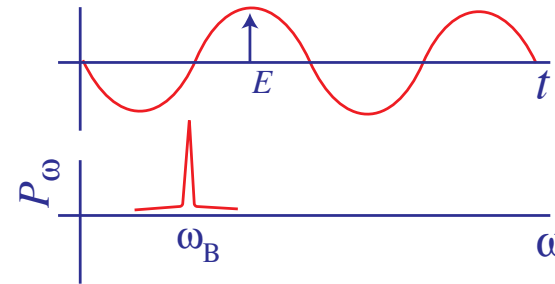
Frequency Spectrum

- The electron gyrates at a frequency

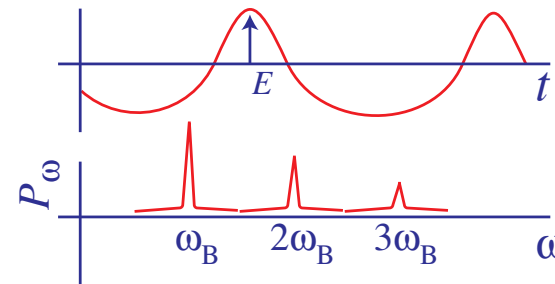
$$\omega_B = \frac{eB}{\gamma mc}$$

- Non relativistic, or cyclotron limit, given by the dipole approximation.

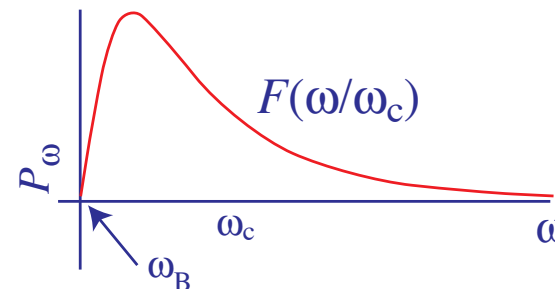
The electric field of the radiation follows a sinusoid:
the frequency structure is a near delta function at ω_B



cyclotron



pulse
sharpening



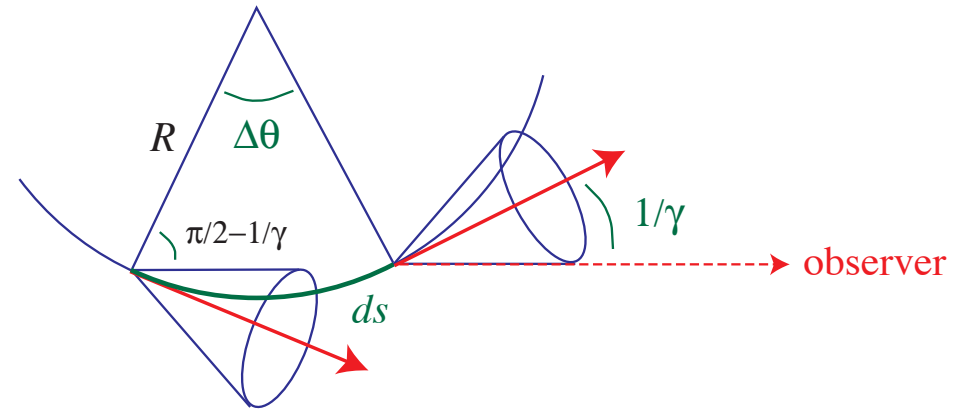
synchrotron
(blended)

Frequency Spectrum

- As the velocities become relativistic beaming sharpens the time profile while remaining periodic
- Appears as a series of delta functions at integer multiples of ω_B
- The beaming sets the width of the pulses and hence a cut off in the frequency spectrum of $\omega_c \sim 1/\Delta t$
- In the ultrarelativistic limit, the frequency content extends to $\omega_c \gg \omega_B$ and forms a continuum due to the range of γ 's

Frequency Spectrum

- The profile width is given by beaming by determining the length of time the emission is observable. Take the ultra-relativistic limit where the radiation makes a cone of angle $\alpha \sim 1/\gamma$
- The observer first sees the radiation when the velocity (tangent to the spiral) is $1/\gamma$ from the line of sight and continues until it is $1/\gamma$ on the other side
- These two points and the center form an equilateral triangle



Frequency Cutoff

- The angle along the spiral traversed during the emission + 2 angles

$$\Delta\theta + 2 \left(\frac{\pi}{2} - \frac{1}{\gamma} \right) = \pi$$
$$\Delta\theta = \frac{2}{\gamma}$$

- The arclength traversed given a gyration radius R combined with velocity gives the duration of emission

$$\Delta s = R\Delta\theta = v\Delta t_{\text{em}} \quad \rightarrow \quad \Delta t_{\text{em}} = \frac{R\Delta\theta}{v} = \frac{2R}{\gamma v}$$

Frequency Cutoff

- Now eliminate the gyration radius R by expressing the angular speed as

$$\frac{\Delta\theta}{\Delta t_{\text{em}}} = \frac{v}{R}$$

and extracting the angular speed from the centripetal acceleration

$$\gamma m \frac{\Delta \mathbf{v}}{\Delta t_{\text{em}}} = \gamma m v \frac{\Delta \theta}{\Delta t_{\text{em}}} = \frac{e}{c} \mathbf{v} \times \mathbf{B} = e \frac{v}{c} B \sin \alpha \quad [\Delta v = v \Delta \theta]$$

Frequency Cutoff

- Reexpress using $\Delta\theta/\Delta t_{\text{em}} = v/R$

$$\gamma m v^2 \frac{1}{R} = e \frac{v}{c} B \sin \alpha$$
$$R = \frac{\gamma m v c}{e B \sin \alpha} = \frac{v}{\omega_B \sin \alpha}$$
$$\Delta t_{\text{em}} = \frac{2}{\gamma \omega_B \sin \alpha}$$

- The arrival time is further shortened in that light must travel across the path difference between beginning and end of the emission

$$\Delta t_{\text{trav}} = \frac{\Delta s}{c} = \frac{2}{\gamma c} \frac{v}{\omega_B \sin \alpha}$$

Frequency Distribution

- The emission only beats the particle by a small amount leading to a smaller observed pulse duration

$$\Delta t = \Delta t_{\text{em}} - \Delta t_{\text{trav}} = \frac{2}{\gamma \omega_B \sin \alpha} \left(1 - \frac{v}{c}\right)$$
$$\Delta t \approx \frac{1}{\gamma^3 \omega_B \sin \alpha} = \frac{mc}{\gamma^2 e B \sin \alpha} \quad \left[1 - \frac{v}{c} \approx \frac{1}{2\gamma^2}\right]$$

- Define a critical frequency

$$\omega_c = \frac{3}{2\Delta t} = \frac{3}{2} \gamma^2 \frac{eB}{mc} \sin \alpha$$

Frequency Spectrum

- Scaling relations: the frequency spectrum a function of ω/ω_c

$$P_\omega = CF(\omega/\omega_c)$$

and the total power is known

$$P = \frac{\sigma_T}{4\pi} \frac{v^2 B^2}{c} \gamma^2 \sin^2 \alpha = \int_0^\infty P_\omega d\omega$$

- Solve for C up to an order unity coefficient $\int F(x)dx = N$

$$P = C\omega_c \int_0^\infty F(x)dx = CN \frac{3}{2} \gamma^2 \frac{eB}{mc} \sin \alpha$$

$$C = \frac{\sigma_T}{4\pi} \frac{v^2 B^2}{c} \frac{\gamma^2 \sin^2 \alpha}{N \frac{3}{2} \gamma^2 \frac{eB}{mc} \sin \alpha} = \frac{\sigma_T}{6\pi} \frac{v^2 B}{c} \frac{mc}{Ne} \sin \alpha$$

$$= \frac{4}{9} B \sin \alpha \frac{e^3}{mc^2} \frac{1}{N} \quad \left[\sigma_T = \frac{8\pi}{3} \frac{e^4}{m^2 c^4} \right]$$

Electron Distribution

- Put it together

$$P_\omega = \frac{4}{9} B \sin \alpha \frac{e^3}{mc^2} \frac{1}{N} F(\omega/\omega_c) \quad \alpha = \text{polar angle wrt } \mathbf{B}$$

- Power law electron distribution $n_{e,\gamma} \propto \gamma^{-p}$

$$j_\nu \propto \int d\gamma n_{e,\gamma} P_\omega \propto \int n_{e,\gamma} F(\omega/\omega_c) d\gamma \propto \int F(\omega/\omega_c) \gamma^{-p} d\gamma$$

- Transform to $x = \omega/\omega_c = \omega/A\gamma^2$

$$dx = -\frac{2\omega}{A\gamma^3} d\gamma \rightarrow \gamma^{-p} d\gamma = \frac{\gamma^{-p+3} A}{2\omega} dx = \left(\frac{\omega}{Ax}\right)^{(-p+3)/2} \frac{A}{2\omega} dx$$

$$j_\nu \propto \omega^{(-p+1)/2}$$

- Just like inverse Compton scattering, the spectrum of radiation has a power law ω^{-s} with $s = (p - 1)/2$

Full Calculation

- A detailed integration of the orbits similar to bremsstrahlung yields

$$\frac{dW}{dt d\omega} \equiv P_\omega = \frac{\sqrt{3}}{2\pi} \frac{e^3 B \sin \alpha}{mc^2} F(\omega/\omega_c)$$
$$F(x) = x \int_x^\infty K_{5/3}(y) dy$$

with asymptotic behavior

$$\lim_{x \ll 1} F(x) = \frac{4\pi}{\sqrt{3}\Gamma(1/3)} \left(\frac{x}{2}\right)^{1/3}$$
$$\lim_{x \gg 1} F(x) = \left(\frac{\pi}{2}\right)^{1/2} e^{-x} x^{1/2}$$

and a maximum at $x = 0.29$

Polarization

- Polarization in direction perpendicular to \mathbf{B} vs complement with a component parallel

$$\Pi = \frac{P_{\perp,\omega} - P_{\parallel,\omega}}{P_{\perp,\omega} + P_{\parallel,\omega}} = \frac{G(x)}{F(x)} \quad G(x) = xK_{2/3}(x)$$

- The polarization is linear (when integrated over pitch angles) and integrated over frequencies for a single γ is 75%
- For a power law distribution of electron energies

$$\Pi = \frac{p + 1}{p + \frac{7}{3}}$$

Radiative Transfer

- Can think of the emission and absorption as the sum over a discrete set of states or lines with $h\nu = E_2 - E_1$

$$j_\nu = \frac{h\nu}{4\pi} \sum_{E_2} \sum_{E_1} n_2 A_{21} \phi_{21}(\nu) = \frac{1}{4\pi} \sum_{E_2} P_\nu n_2 = \frac{1}{4\pi} \sum_{E_2} 2\pi P_\omega n_2$$
$$h\nu \sum_{E_1} A_{21} \phi_{21}(\nu) = 2\pi P_\omega$$

where $\phi_{21}(\nu)$ the profile enforces energy conservation
 $h\nu = E_2 - E_1$

$$\alpha_\nu = \frac{h\nu}{4\pi} \sum_{E_1} \sum_{E_2} [n_1 B_{12} - n_2 B_{21}] \phi_{21}(\nu)$$
$$\alpha_\nu = \frac{h\nu}{4\pi} \sum_{E_1} \sum_{E_2} [n_1 - n_2] \frac{c^2}{2h\nu^3} A_{21} \phi_{21}(\nu)$$

Synchrotron Self Absorption

- So that

$$\alpha_\nu = \frac{c^2}{4h\nu^3} \sum_{E_2} [n_1 - n_2] P_\omega$$

- Now switch to the continuous phase space distribution

$$\sum_E n \rightarrow g_e \int \frac{d^3q}{(2\pi\hbar)^3} f_e$$

$$\alpha_\nu = \frac{c^2}{4h\nu^3} g_e \int \frac{d^3q}{(2\pi\hbar)^3} [f_e(E - h\nu) - f_e(E)] P_\omega$$

$$j_\nu = \frac{1}{2} g_e \int \frac{d^3q}{(2\pi\hbar)^3} f_e P_\omega = \frac{1}{4\pi} g_e \int \frac{d^3q}{(2\pi\hbar)^3} f_e P_\nu$$

Synchrotron Self Absorption

- Check: for a thermal distribution $f_e = e^{-(E-\mu)/kT}$ then

$$f_e(E - h\nu) - f_e(E) = e^{-(E-\mu)/kT} (e^{h\nu/kT} - 1)$$

$$\frac{j_\nu}{\alpha_\nu} = \frac{c^2}{2h\nu^3} \frac{1}{e^{h\nu/kT} - 1} = B_\nu$$

- In the limit that $h\nu \ll E$ then

$$\alpha_\nu \approx -\frac{c^2}{4\nu^2} g_e \int \frac{d^3q}{(2\pi\hbar)^3} \frac{\partial f_e}{\partial E} P_\omega$$

- For a power law distribution of relativistic electrons

$$n_e = \int d\gamma n_{e,\gamma} = g_e \int \frac{d^3q}{(2\pi\hbar)^3} f_e = g_e \int \frac{q^2 dq}{2\pi^2 \hbar^3} f_e \propto \int \gamma^2 d\gamma f_e$$

- Thus $f_e \propto \gamma^{-2} n_{e,\gamma}$

Synchrotron Self Absorption

- So that for a power law $n_{e,\gamma} \propto \gamma^{-p}$ and $\partial f_e / \partial \gamma \propto \gamma^{-3-p}$
- Combine with phase space integration
 $q^2 dq \propto \gamma^2 d\gamma$ and $P_\omega \propto F(x)$ and $x = \omega / \omega_c = \omega / A \gamma^2$

$$\alpha_\nu \propto \nu^{-2} \int d\gamma \gamma^{-p-1} F(\omega / \omega_c) \propto \nu^{-p/2-2}$$

- Compare to emission integral

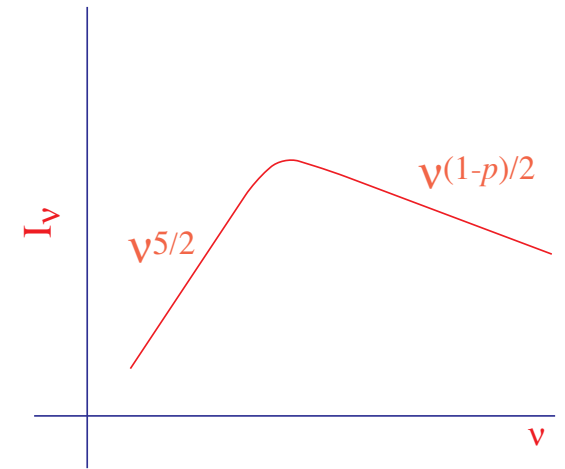
$$j_\nu \propto \int n_{e,\gamma} F(\omega / \omega_c) d\gamma \propto \int F(\omega / \omega_c) \gamma^{-p} d\gamma \propto \nu^{-p/2+1/2},$$

- Source function has extra $\nu^{5/2}$: ν^2 from absorption formula and $\nu^{1/2}$ from extra γ

$$S_\nu = \frac{j_\nu}{\alpha_\nu} \propto \nu^{5/2}$$

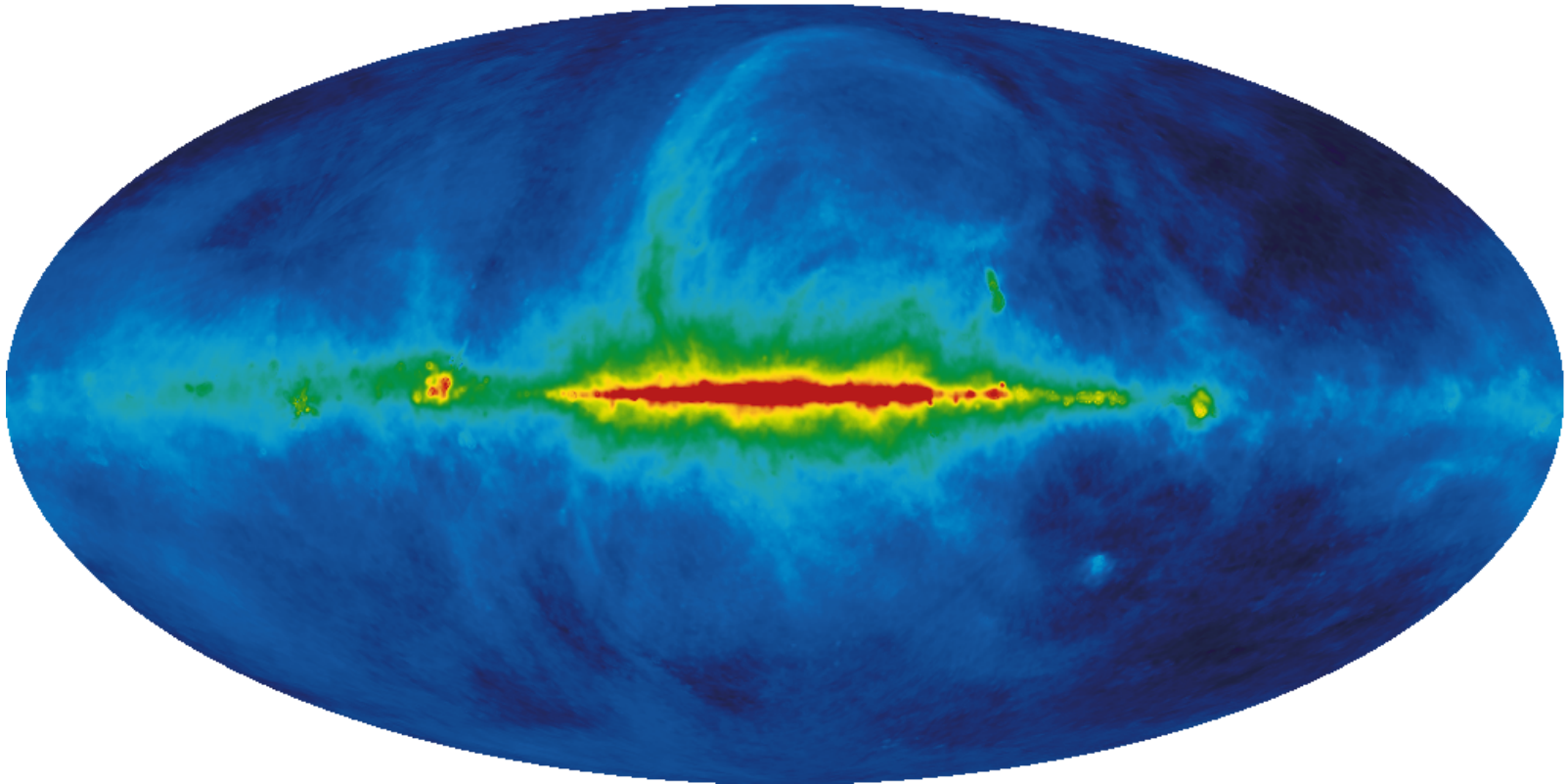
Synchrotron Self Absorption

- At high synchrotron absorption optical depth (low frequencies) $I_\nu \rightarrow S_\nu \propto \nu^{5/2}$
- At low optical depth $I_\nu \propto j_\nu \propto \nu^{-(p-1)/2}$
- The synchrotron spectrum has a characteristic turn over between the regimes



Galactic Synchrotron

- HASLAM 408MHz synchrotron map



Galactic Synchrotron

- WMAP 23GHz emission and polarization map

