

Astro 321

Set 1: FRW Cosmology

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# FRW Cosmology

- The Friedmann-Robertson-Walker (FRW sometimes Lemaitre, FLRW) cosmology has two elements
  - The FRW geometry or metric
  - The FRW dynamics or Einstein/Friedmann equation(s)
- Same as the two pieces of General Relativity (GR)
  - A metric theory: geometry tells matter how to move
  - Field equations: matter tells geometry how to curve
- Useful to separate out these two pieces both conceptually and for understanding alternate cosmologies, e.g.
  - Modifying gravity while remaining a metric theory
  - Breaking the homogeneity or isotropy assumption under GR

# FRW Geometry

- FRW geometry = homogeneous and isotropic on large scales
- Universe observed to be nearly isotropic (e.g. CMB, radio point sources, galaxy surveys)
- Copernican principle: we're not special, must be isotropic to all observers (all locations)

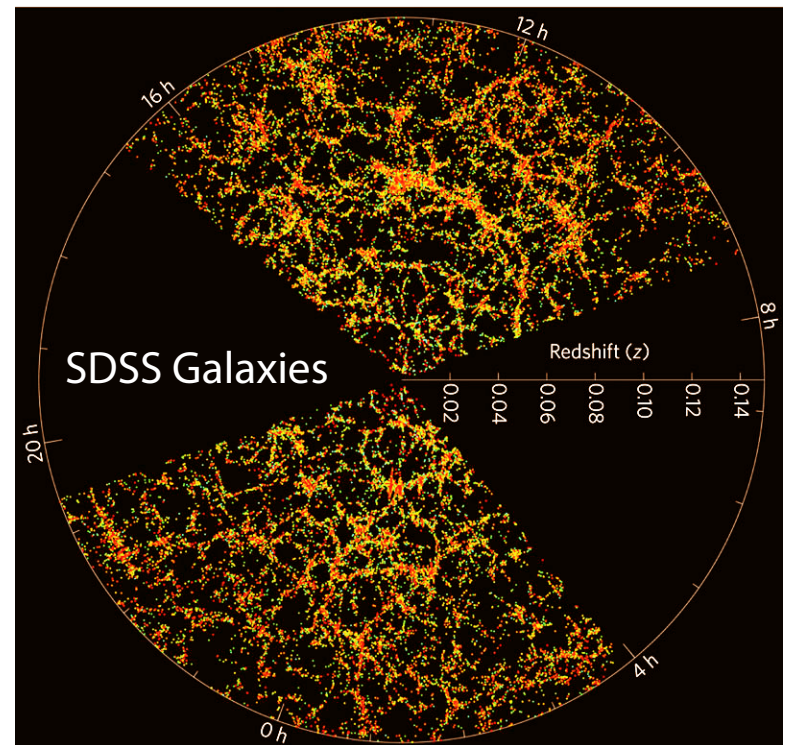
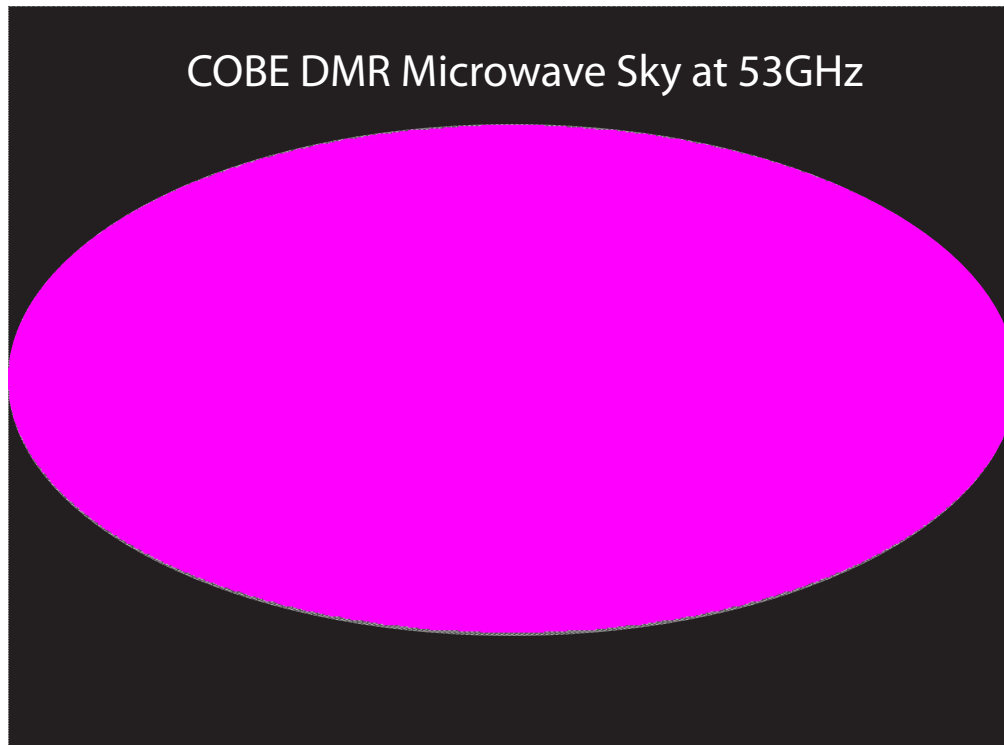
Implies homogeneity

Verified through galaxy redshift surveys

- FRW cosmology (homogeneity, isotropy & field equations) generically implies the expansion of the universe, except for special unstable cases

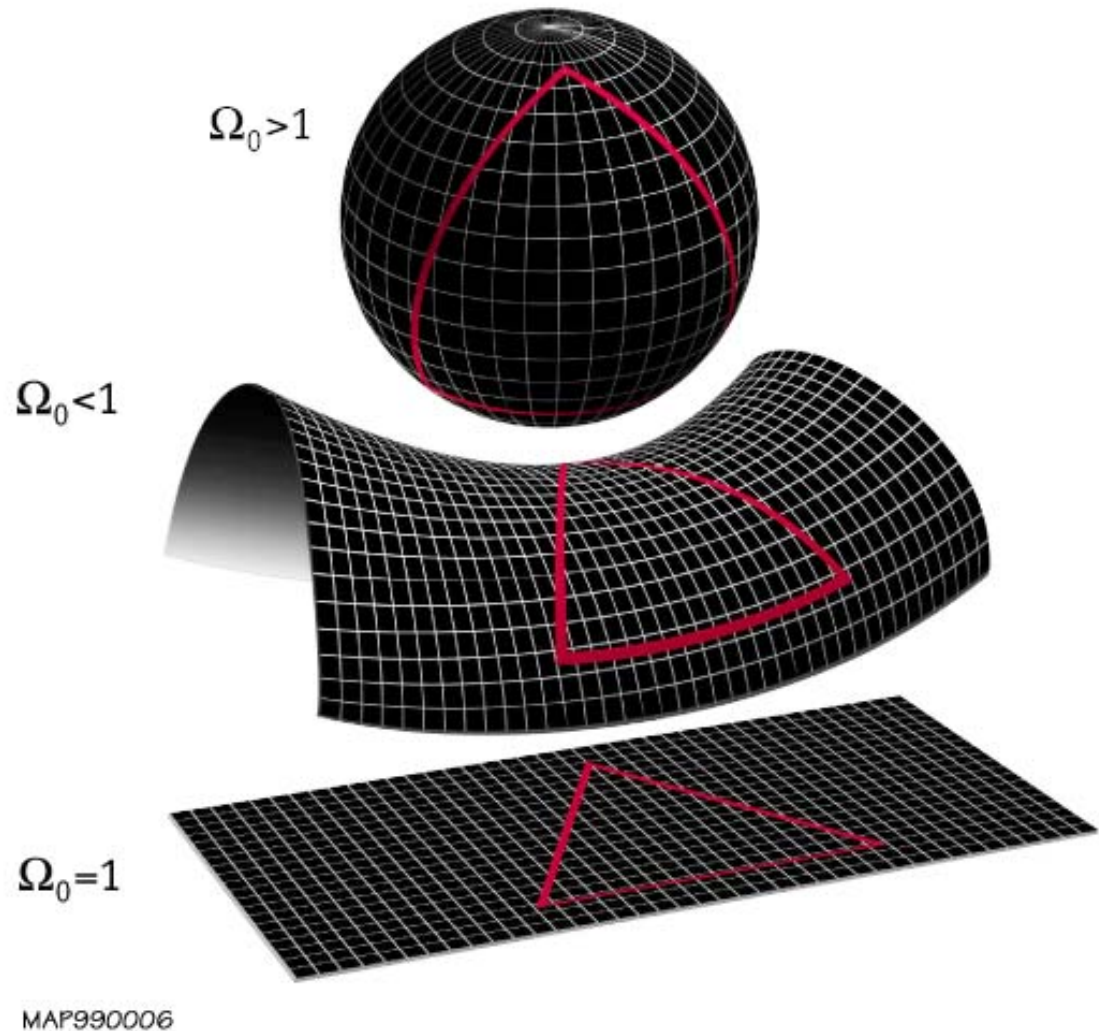
# Isotropy & Homogeneity

- Isotropy: CMB isotropic to  $10^{-3}$ ,  $10^{-5}$  if dipole subtracted
- Redshift surveys show return to homogeneity on the  $>100\text{Mpc}$  scale



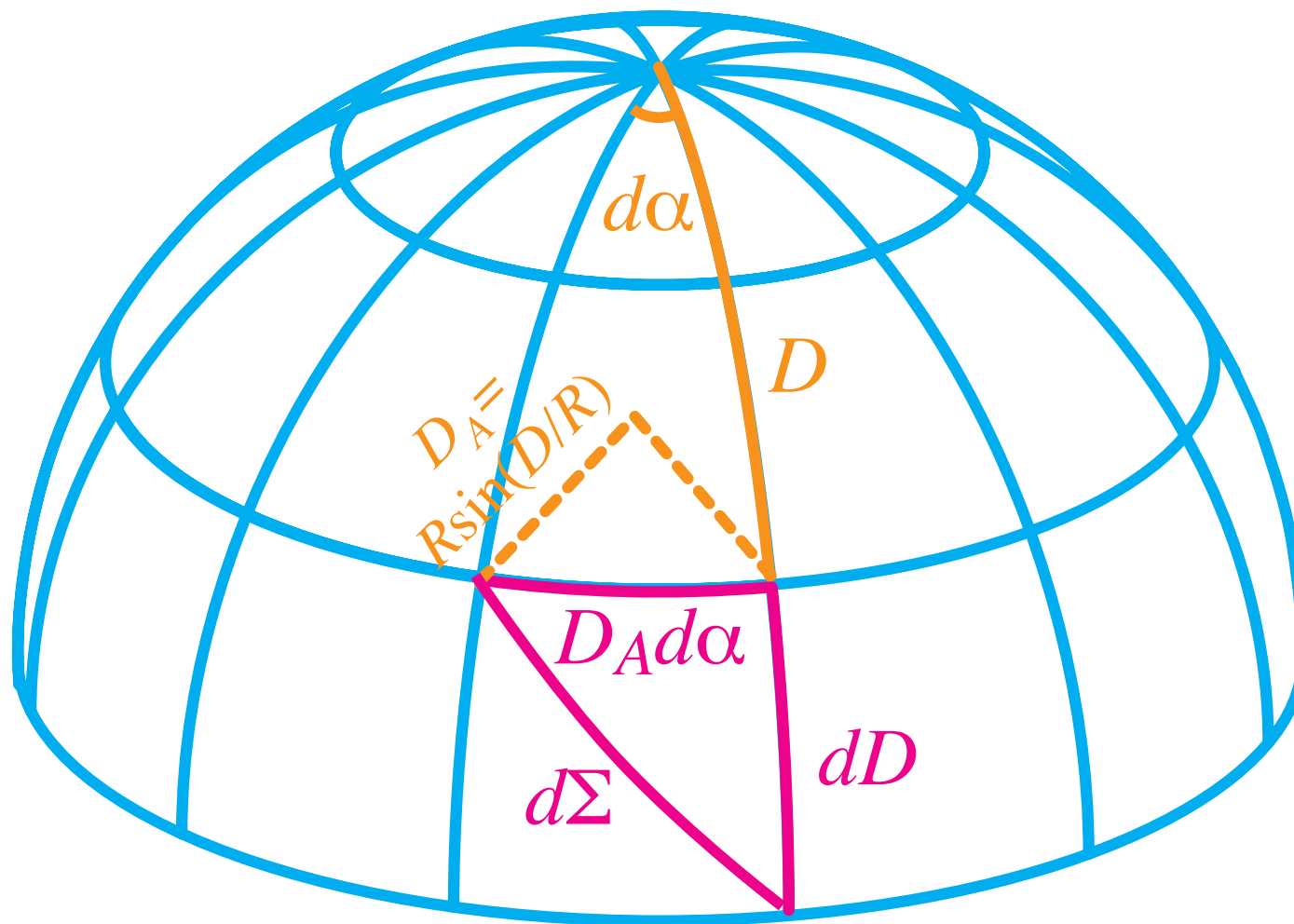
# FRW Geometry

- Spatial geometry is that of a constant curvature  $K = 1/R^2$
- Positive: sphere  
Negative: saddle  
Flat: plane
- Metric tells us how to measure distances on this surface



# FRW Geometry

- Closed geometry of a sphere of radius  $R$
- Suppress 1 dimension  $\alpha$  represents total angular separation  $(\theta, \phi)$



# FRW Geometry

- Two types of distances:

- Radial distance on the arc  $D$

Distance (for e.g. **photon**) traveling along the arc

- Angular diameter distance  $D_A$

Distance inferred by the **angular separation**  $d\alpha$  for a known **transverse** separation (on a constant latitude)  $D_A d\alpha$

Relationship  $D_A = R \sin(D/R)$

- As if background geometry (gravitationally) **lenses** image
- **Positively curved** geometry  $D_A < D$  and objects are **further** than they appear

- **Negatively curved** universe  $R$  is imaginary and

$$R \sin(D/R) = i|R| \sin(D/i|R|) = |R| \sinh(D/|R|)$$

and  $D_A > D$  objects are closer than they appear

# Angular Diameter Distance

- 3D distances restore usual spherical polar angles

$$\begin{aligned} d\Sigma^2 &= dD^2 + D_A^2 d\alpha^2 \\ &= dD^2 + D_A^2 (d\theta^2 + \sin^2 \theta d\phi^2) \end{aligned}$$

- GR allows arbitrary choice of coordinates, alternate notation is to use  $D_A$  as radial coordinate
- $D_A$  useful for describing observables (flux, angular positions)
- $D$  useful for theoretical constructs (causality, relationship to temporal evolution)



# Angular Diameter Distance

- The line element is often also written using  $D_A$  as the coordinate distance

$$dD_A^2 = \left( \frac{dD_A}{dD} \right)^2 dD^2$$

$$\left( \frac{dD_A}{dD} \right)^2 = \cos^2(D/R) = 1 - \sin^2(D/R) = 1 - (D_A/R)^2$$

$$dD^2 = \frac{1}{1 - D_A^2/R^2} dD_A^2$$

and defining the curvature  $K = 1/R^2$  the line element becomes

$$d\Sigma^2 = \frac{1}{1 - D_A^2 K} dD_A^2 + D_A^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

where  $K < 0$  for a negatively curved space

# Volume Element

- Metric also defines the volume element

$$\begin{aligned} dV &= (dD)(D_A d\theta)(D_A \sin \theta d\phi) \\ &= D_A^2 dD d\Omega \end{aligned}$$

where  $d\Omega = \sin \theta d\theta d\phi$  is solid angle

- Most of classical cosmology boils down to these three quantities, (comoving) radial distance, (comoving) angular diameter distance, and volume element
- For example, distance to a high redshift supernova, angular size of the horizon at last scattering and BAO feature, number density of clusters...

# Comoving Coordinates

- Remaining degree of freedom (preserving homogeneity and isotropy) is the **temporal evolution** of overall **scale factor**
- Relates the **geometry** (fixed by the radius of curvature  $R$ ) to **physical coordinates** – a function of time only

$$d\sigma^2 = a^2(t)d\Sigma^2$$

our conventions are that the scale factor today  $a(t_0) \equiv 1$

- Similarly **physical distances** are given by  $d(t) = a(t)D$ ,  
 $d_A(t) = a(t)D_A$ .
- Distances in **upper case** are **comoving**; lower, physical  
Do not change with time  
Simplest coordinates to work out geometrical effects

# Time and Conformal Time

- Proper time (with  $c = 1$ )

$$\begin{aligned}d\tau^2 &= dt^2 - d\sigma^2 \\ &= dt^2 - a^2(t)d\Sigma^2\end{aligned}$$

- Taking out the scale factor in the time coordinate

$$d\tau^2 \equiv a^2(t) (d\eta^2 - d\Sigma^2)$$

$d\eta = dt/a$  defines **conformal time** – useful in that photons travelling radially from observer then obey

$$\Delta D = \Delta\eta = \int \frac{dt}{a}$$

so that **time** and **distance** may be interchanged

# FRW Metric

- Relationship between coordinate differentials and space-time separation defines the **metric**  $g_{\mu\nu}$
- Mostly plus convention  $ds^2 = -d\tau^2$

$$ds^2 \equiv g_{\mu\nu} dx^\mu dx^\nu = a^2(\eta)(-d\eta^2 + d\Sigma^2)$$

Einstein summation - **repeated** lower-upper pairs **summed**

- Usually we will use **comoving coordinates** and **conformal time** as the  $x^\mu$  unless otherwise specified – metric for other choices are related by  $a(t)$

# Horizon

- Distance travelled by a photon in the whole lifetime of the universe defines the **horizon**
- Since  $d\tau = 0$ , the horizon is simply the **elapsed conformal time**

$$D_{\text{horizon}}(t) = \int_0^t \frac{dt'}{a} = \eta(t)$$

- Horizon always **grows with time**
- Always a point in time before which two observers separated by a distance  $D$  could **not** have been in **causal contact**

# Horizon

- **Horizon problem:** why is the universe homogeneous and **isotropic** on large scales especially for objects seen at early times, e.g. CMB, when horizon small

- Intuition: in each doubling (or efolding) of the scale factor, photons travel larger and larger distances

Consequence: horizon is approximately the distance travelled in the last efolding

- To avoid the horizon problem, we want the distance to get smaller and smaller with each efolding
- Quantify by transforming time to efolds through the Hubble parameter

# Hubble Parameter

- Useful to define the expansion rate or Hubble parameter

$$H(t) \equiv \frac{1}{a} \frac{da}{dt} = \frac{d \ln a}{dt}$$

fractional change in the scale factor per unit time -  $\ln a = N$  is also known as the e-folds of the expansion

- Cosmic time becomes

$$t = \int dt = \int \frac{d \ln a}{H(a)}$$

- Conformal time becomes

$$\eta = \int \frac{dt}{a} = \int \frac{d \ln a}{a H(a)}$$



# Horizon Problem Redux

- Does  $aH$  increase or decrease with  $a$ ?
- If  $aH$  decreases then for each successive  $\Delta \ln a$ , a photon travels a larger  $\Delta D$ , total distance dominated by last efold
- If  $aH$  increases then for each successive  $\Delta \ln a$ , a photon travels a smaller  $\Delta D$ , total distance dominated by first efold
- Critical point is when the acceleration of the expansion switches sign

$$\frac{d(aH)}{dt} = \frac{d^2 a}{dt^2}$$

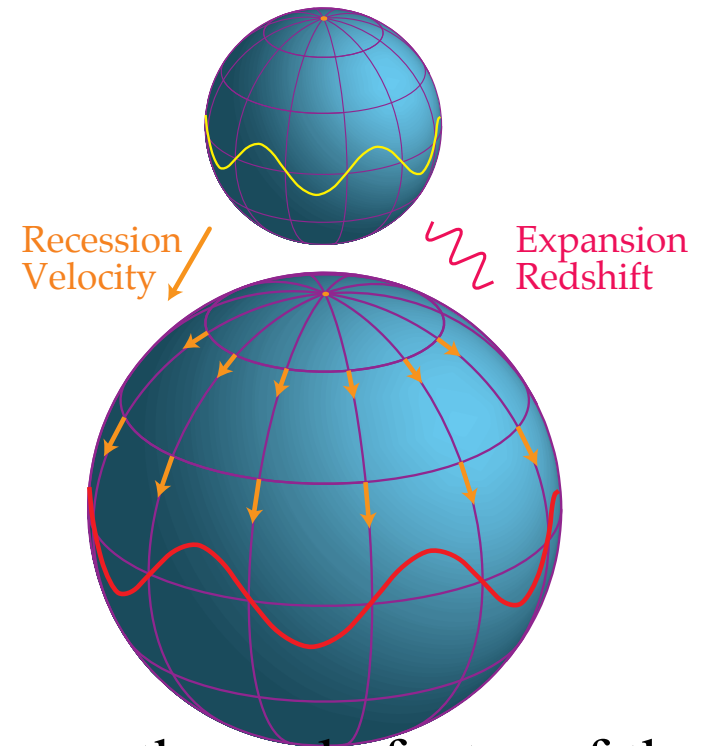
# Redshift

- **Wavelength** of light “stretches” with the scale factor
- The physical wavelength  $\lambda_{\text{emit}}$  associated with an observed wavelength today  $\lambda_{\text{obs}}(a = 1)$  (or comoving=physical units today) is

$$\lambda_{\text{emit}} = a(t) \lambda_{\text{obs}}$$

so that the redshift of spectral lines measures the scale factor of the universe at  $t$ ,  $1 + z = 1/a$ .

- Interpreting the redshift as a **Doppler shift**, objects recede in an expanding universe
- More generally the de Broglie wavelength of any particle redshifts in this way



# Distance-Redshift Relation

- Given **atomically known** rest wavelength  $\lambda_{\text{emit}}$ , redshift can be precisely measured from spectra
- Combined with a measure of **distance**, distance-redshift  $D(z) \equiv D(z(a))$  can be inferred - given that photons travel  $D = \Delta\eta$  this tells us how the scale factor of the universe evolves with time.
- Related to the **expansion history** as

$$D(a) = \int dD = \int_a^1 \frac{d \ln a'}{a' H(a')}$$
$$[d \ln a' = -d \ln(1 + z) = -a' dz]$$
$$D(z) = - \int_z^0 \frac{dz'}{H(z')} = \int_0^z \frac{dz'}{H(z')}$$

# Hubble Law

- Note limiting case is the Hubble law

$$\lim_{z \rightarrow 0} D(z) = z/H(z=0) \equiv z/H_0$$

independently of the geometry and expansion dynamics

- Hubble constant usually quoted as as dimensionless  $h$

$$H_0 = 100h \text{ km s}^{-1} \text{ Mpc}^{-1}$$

- Observationally  $h \sim 0.7$  (see below)

# Scale of the Universe

- In natural units of  $\hbar = c = 1$  used here,  $H_0$  sets an length, time, energy, mass scale
- $H_0^{-1} = 9.7778 (h^{-1} \text{ Gyr})$   
e-folding time scale of the expansion (Hubble time), age of (decelerating) universe
- $H_0^{-1} = 2997.9 (h^{-1} \text{ Mpc})$   
Observable length scale (**Hubble scale**), horizon scale of (decelerating) universe
- $H_0 = 2.1332h \times 10^{-33} \text{eV} = m_{\text{de}}$   
Mass scale of explanations of dark energy
- $H_0 = 10^{-6}h \times (2.9979 \text{ kpc})^{-1} = (GM/r) \times r^{-1}$   
Acceleration/MOND scale - order of magnitude at which dark matter in galaxies flatten rotation curve ( $\sim 10^{-10} \text{m s}^{-2}$ )

# Scale of the Universe

- Since  $GM/r$  is dimensionless and  $r$  has inverse  $M$  dimensions, gravity sets a natural mass scale in the reduced Planck mass  
 $M_{\text{Pl}} = 1/\sqrt{8\pi G} = 1.22 \times 10^{19} \text{ GeV}$

$$\begin{aligned} M^4 &\equiv \rho_c = 3H_0^2/8\pi G \\ &= (3.000 \times 10^{-12} \text{ GeV})^4 h^2 = 8.098 \times 10^{-47} h^2 \text{ GeV}^4 \end{aligned}$$

Density scale of the expansion, critical energy density (see below)

- $M/M_{\text{Pl}} = 2.46h^{1/2} \times 10^{-31}$  – seems highly unnatural in natural units! (famous 120 orders of magnitude in density, see below)
- $M = 3^{1/4} \sqrt{m_{\text{de}} M_{\text{Pl}}}$ , geometric mean
- $m_{\text{de}}$  as far from any standard model particle – what protects such a hierarchy? (note that  $M$  is comparable to neutrino masses)

# Measuring $D(z)$

- Standard Ruler: object of known physical size

$$\lambda = a(t)\Lambda$$

subtending an observed angle  $\alpha$  on the sky  $\alpha$

$$\alpha = \frac{\Lambda}{D_A(z)} \equiv \frac{\lambda}{d_A(z)}$$

$$d_A(z) = aD_A(a) = \frac{D_A(z)}{1+z}$$

where, by analogy to  $D_A$ ,  $d_A$  is the physical angular diameter distance

- Since  $D_A \rightarrow D_{\text{horizon}}$  whereas  $(1+z)$  unbounded, angular size of a fixed physical scale at high redshift actually increases with radial distance

# Measuring $D(z)$

- **Standard Candle:** object of known luminosity  $L$  with a measured flux  $F$  (energy/time/area)
  - Comoving surface area  $4\pi D_A^2$
  - Frequency/energy redshifts as  $(1 + z)$
  - Time-dilation or arrival rate of photons (crests)  $dt = a d\eta$  lowered as  $(1 + z)$  vs emission rate:

$$F = \frac{L}{4\pi D_A^2} \frac{1}{(1 + z)^2} \equiv \frac{L}{4\pi d_L^2}$$

- So **luminosity distance**

$$d_L = (1 + z) D_A = (1 + z)^2 d_A$$

- As  $z \rightarrow 0$ ,  $d_L = d_A = D_A$



# Olber's Paradox

- Surface brightness

$$S = \frac{F}{\Delta\Omega} = \frac{L}{4\pi d_L^2} \frac{d_A^2}{\lambda^2}$$

- In a non-expanding geometry (regardless of curvature), surface brightness is conserved  $d_A = d_L$

$$S = \text{const.}$$

- So since each sight line in universe full of stars will eventually end on surface of star, night sky should be as bright as sun (not infinite)
- In an expanding universe

$$S \propto (1 + z)^{-4}$$

# Olber's Paradox

- Second piece: **age finite** so even if stars exist in the early universe, not all sight lines end on stars
- But even as **age** goes to infinity and the number of sight lines goes to 100%, **surface brightness** of distant objects (of fixed physical size) goes to **zero**
  - Angular size increases
  - Redshift of energy and arrival time

# Measuring $D(z)$

- Ratio of fluxes or difference in log flux (magnitude) measurable independent of knowing luminosity

$$m_1 - m_2 = -2.5 \log_{10}(F_1/F_2)$$

related to  $d_L$  by definition by inverse square law

$$m_1 - m_2 = 5 \log_{10}[d_L(z_1)/d_L(z_2)]$$

- If absolute magnitude is known

$$m - M = 5 \log_{10}[d_L(z)/10\text{pc}]$$

absolute distances measured, e.g. at low  $z = z_0$  Hubble constant

$$d_L \approx z_0/H_0 \rightarrow H_0 = z_0/d_L$$

- Also standard ruler whose length, calibrated in physical units

# Measuring $D(z)$

- If absolute calibration of standards unknown, then both standard candles and standard rulers measure relative sizes and fluxes

For **standard candle**, e.g. **compare magnitudes** low  $z_0$  to a high  $z$  object involves

$$\Delta m = m_z - m_{z_0} = 5 \log_{10} \frac{d_L(z)}{d_L(z_0)} = 5 \log_{10} \frac{H_0 d_L(z)}{z_0}$$

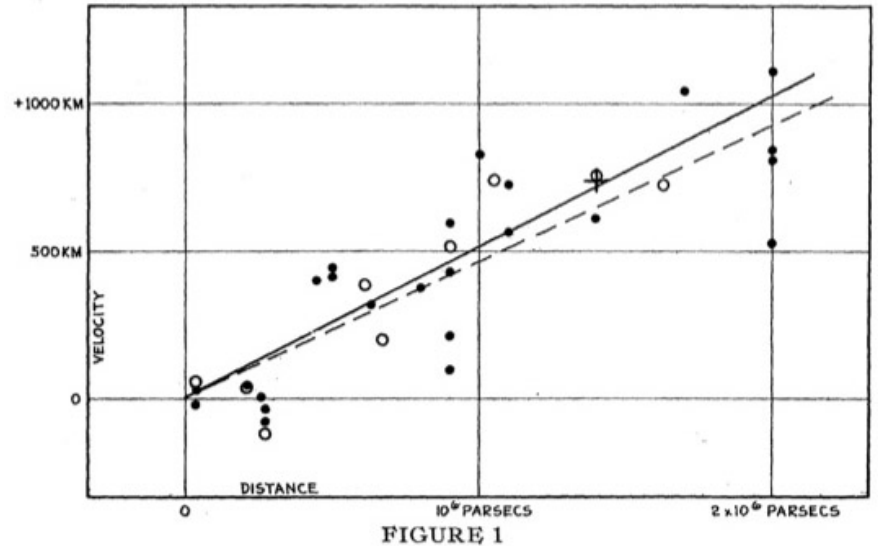
Likewise for a **standard ruler** comparison at the two redshifts

$$\frac{d_A(z)}{d_A(z_0)} = \frac{H_0 d_A(z)}{z_0}$$

- Distances are measured in units of  $h^{-1}$  Mpc.
- **Change** in expansion rate measured as  $H(z)/H_0$

# Hubble Constant

- Hubble in 1929 used the Cepheid period luminosity relation to infer distances to nearby galaxies thereby discovering the expansion of the universe



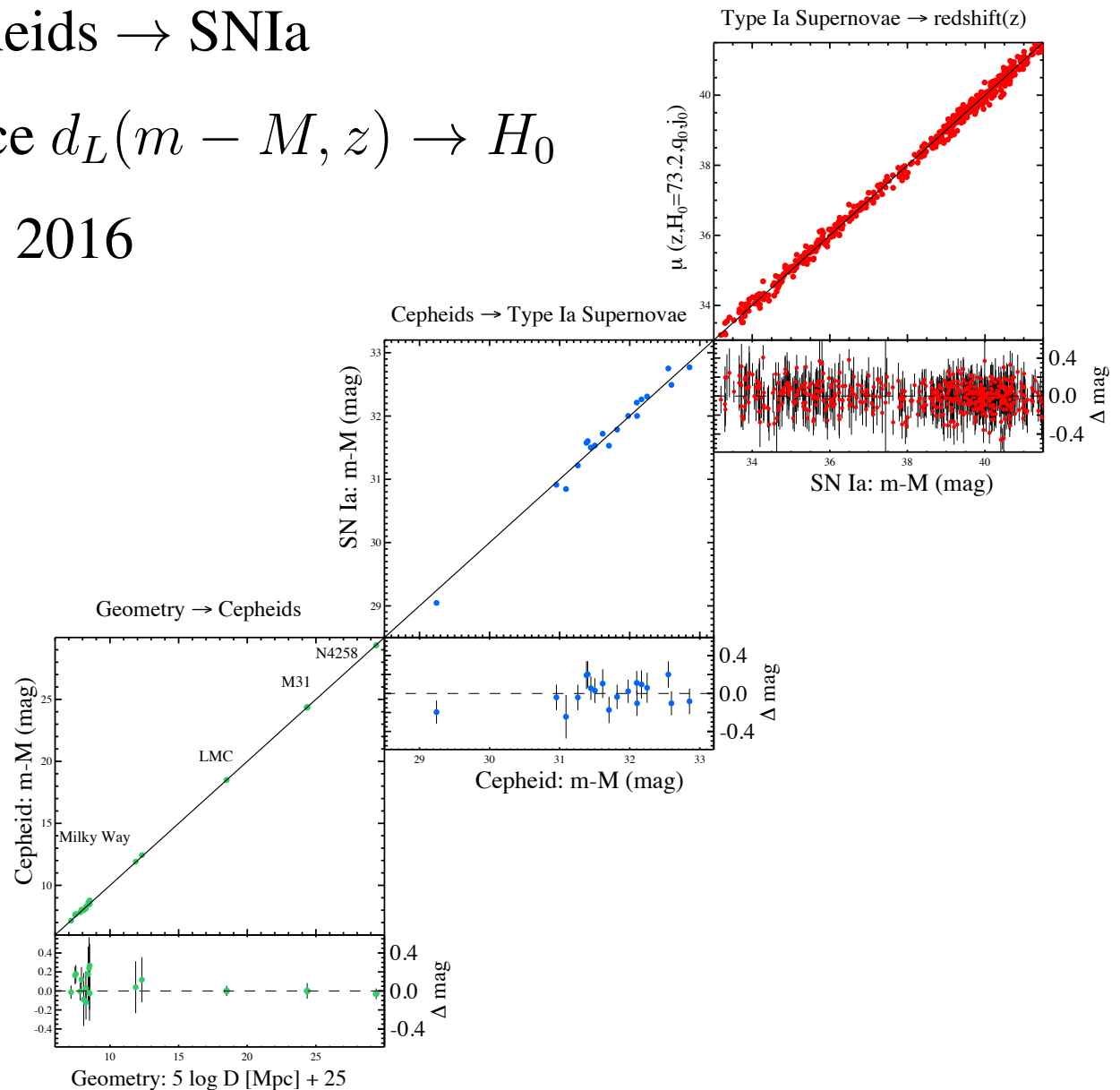
- Hubble actually inferred too large a Hubble constant of  $H_0 \sim 500 \text{ km/s/Mpc}$
- Miscalibration of the Cepheid distance scale - absolute measurement hard, checkered history

# Hubble Constant History

- Took 70 years to settle on this value with a factor of 2 discrepancy persisting until late 1990's
- Difficult measurement since local galaxies where individual Cepheids can be measured have peculiar motions and so their velocity is not entirely due to the “Hubble flow”
- A “distance ladder” of cross calibrated measurements
- Primary distance indicators cepheids, novae planetary nebula, tip of red giant branch, AGN water maser
- GAIA will soon improve geometric calibration of galactic cepheids with parallax measurements
- More luminous secondary distance indicators into the Hubble flow: Tully-Fisher, fundamental plane, surface brightness fluctuations, Type 1A supernova

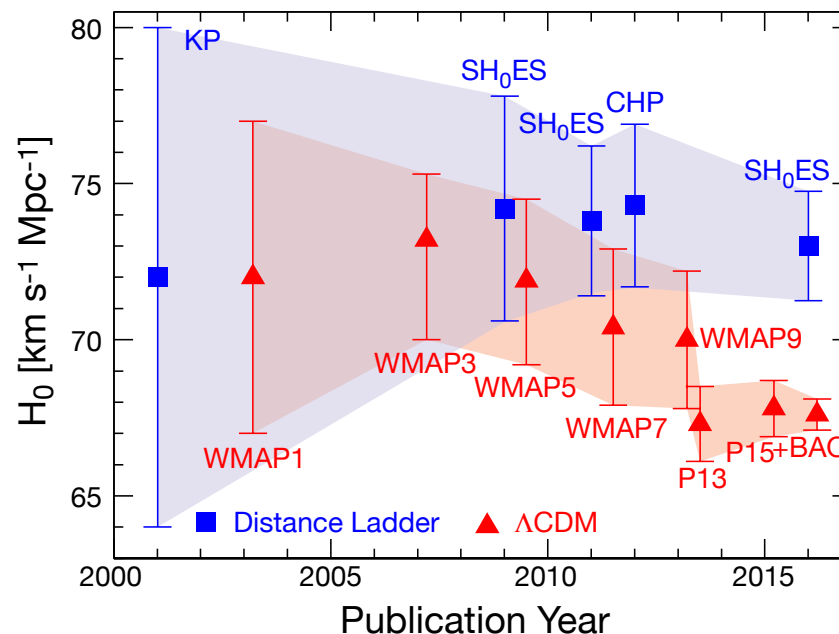
# Modern Distance Ladder

- Geometry  $\rightarrow$  Cepheids  $\rightarrow$  SNIa
- Luminosity distance  $d_L(m - M, z) \rightarrow H_0$
- SH0ES, Riess et al 2016



# Hubble Constant

- $H_0$  now measured as  $73.24 \pm 1.74 \text{ km/s/Mpc}$  by SH0ES calibrating SNIa off cepheids off AGN **water maser** as well as the local distance ladder.
- Comparable precision from Carnegie-Chicago Hubble Program
- Inverse distance ladder: standard ruler CMB calibration at  $z \sim 10^3$  to BAO to SNIa
- Assuming the  $\Lambda\text{CDM}$  model the inverse distance ladder gives:  
 $H_0 = 67.6 \pm 0.5 \text{ km/s/Mpc}$



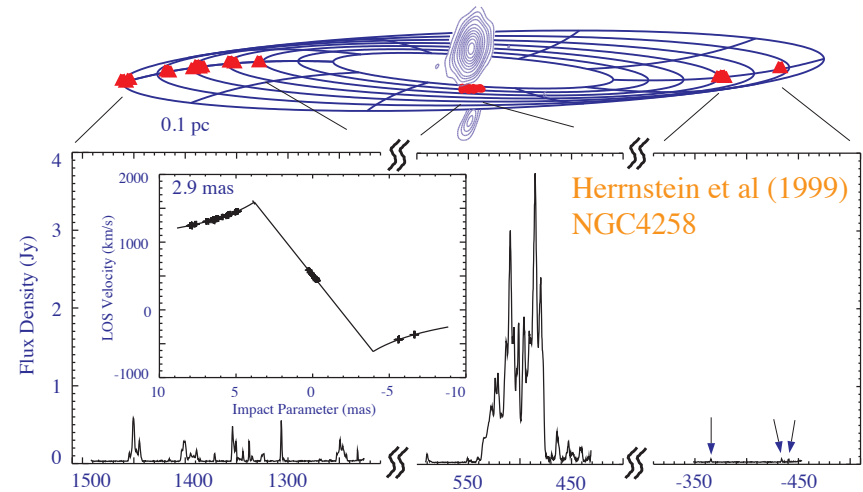


# Hubble Constant

- Given the history and difficulty of connecting these ladders, this agreement is actually quite impressive – but not within the quoted errors
- Resolution remains to be seen: must ensure that both of these precise measurements are *accurate* in the presence of systematics.

# Maser-Cepheid-SN Distance Ladder

- Water maser around AGN, gas in Keplerian orbit
- Measure proper motion, radial velocity, acceleration of orbit
- Method 1: radial velocity plus orbit infer tangential velocity = distance  $\times$  angular proper motion



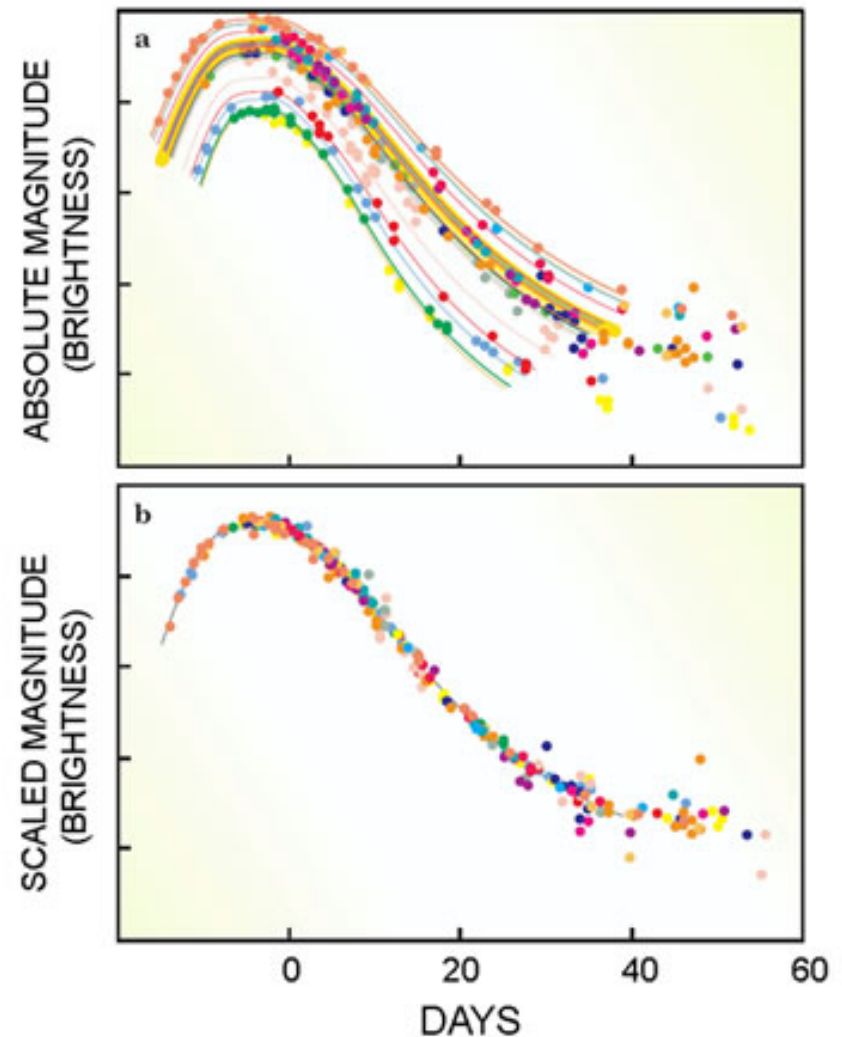
$$v_t = d_A(d\alpha/dt)$$

- Method 2: centripetal acceleration and radial velocity from line infer physical size

$$a = v^2/R, \quad R = d_A\theta$$

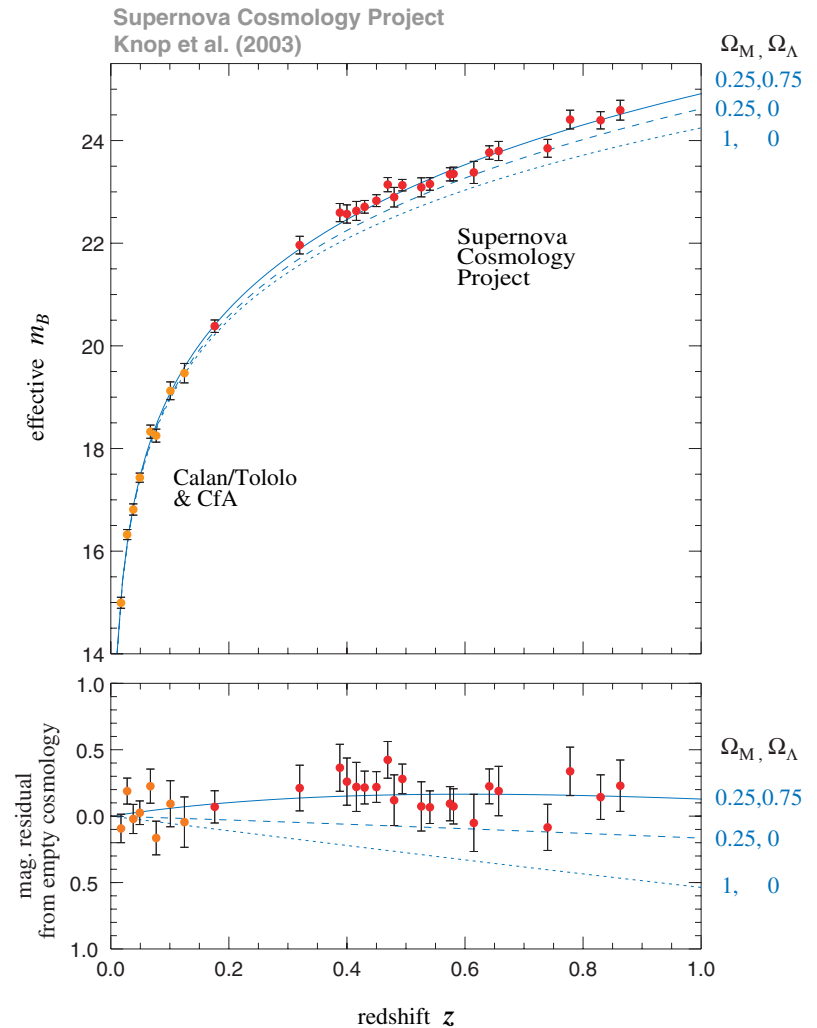
# Supernovae as Standard Candles

- Type 1A supernovae are **white dwarfs** that reach **Chandrasekar mass** where electron degeneracy pressure can no longer support the star, hence a **very regular explosion**
- Moreover, the scatter in absolute magnitude is correlated with the **shape** of the light curve - the rate of decline from peak light, empirical “**Phillips relation**”
- Higher  $^{56}\text{Ni}$ , **brighter** SN, higher opacity, **longer** light curve duration



# Beyond Hubble's Law

- Type 1A are therefore “standardizable” candles leading to a very low scatter  $\delta m \sim 0.15$  and visible out to high redshift  $z \sim 1$
- Two groups in 1999 found that SN more distant at a given redshift than expected
- Cosmic acceleration



# Beyond Hubble's Law

- Using SN as a **relative indicator** (independent of absolute magnitude), comparison of low and high  $z$  gives

$$H_0 D(z) = \int dz \frac{H_0}{H}$$

more distant implies that  $H(z)$  not increasing at expected rate, i.e. is more constant

- Take the limiting case where  $H(z)$  is a **constant** (a.k.a. **de Sitter expansion**)

$$H = \frac{1}{a} \frac{da}{dt} = \text{const}$$

$$\frac{dH}{dt} = \frac{1}{a} \frac{d^2 a}{dt^2} - H^2 = 0$$

$$\frac{1}{a} \frac{d^2 a}{dt^2} = H^2 > 0$$

# Beyond Hubble's Law

- Indicates that the **expansion** of the universe is **accelerating**
- Intuition tells us (FRW dynamics shows) **ordinary matter** decelerates expansion since gravity is **attractive**
- **Ordinary expectation** is that

$$H(z > 0) > H_0$$

so that the Hubble parameter is higher at high redshift

- Or equivalently that **expansion rate decreases** as it **expands**

# FRW Dynamics

- This is as far as we can go on FRW geometry alone - we still need to know how the **scale factor**  $a(t)$  evolves given **matter-energy content**
- **General relativity**: matter tells geometry how to curve, **scale factor** determined by **content**
- Build the **Einstein tensor**  $G_{\mu\nu}$  out of the **metric** and use **Einstein equation** (overdots conformal time derivative)

$$G_{\mu\nu} (= R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R) = 8\pi GT_{\mu\nu}$$

- Easier to work with mixed upper and lower indices since the metric  $g^\mu{}_\nu = \delta^\mu{}_\nu$

# Einstein Equations

- For the FRW metric

$$G^0_0 = -3 \left( H^2 + \frac{K}{a^2} \right)$$
$$G^i_j - G^0_0 \frac{\delta^i_j}{3} = -\frac{2}{a^2} \left( \frac{\ddot{a}}{a} - a^2 H^2 \right) \delta^i_j = -\frac{2}{a} \frac{d^2 a}{dt^2} \delta^i_j,$$

where recall the curvature  $K = 1/R^2$  and overdots are  $d/d\eta$

- Likewise **isotropy** demands that the **stress-energy tensor** take the form

$$T^0_0 = -\rho, \quad T^i_j = p\delta^i_j \quad \rightarrow \quad T^i_j - T^0_0 \frac{\delta^i_j}{3} = p + \rho/3$$

where  $\rho$  is the **energy density** and  $p$  is the **pressure**

- It is **not** necessary to assume that the content is a **perfect fluid** - consequence of **FRW symmetry**



# Friedmann Equations

- Einstein equations given the FRW symmetries become the Friedmann equations

$$H^2 + \frac{K}{a^2} = \frac{8\pi G}{3}\rho$$
$$\frac{1}{a} \frac{d^2 a}{dt^2} = -\frac{4\pi G}{3}(\rho + 3p)$$

- Acceleration source is  $\rho + 3p$  requiring  $p < -\rho/3$  for positive acceleration
- Curvature as an effective energy density component

$$\rho_K = -\frac{3}{8\pi G} \frac{K}{a^2} \propto a^{-2}$$

Positive curvature gives negative effective energy density

# Critical Density

- Friedmann equation for  $H$  then reads

$$H^2(a) = \frac{8\pi G}{3}(\rho + \rho_K) \equiv \frac{8\pi G}{3}\rho_c$$

defining a **critical density** today  $\rho_c$  in terms of the expansion rate

- In particular, its value today is given by the Hubble constant as

$$\rho_c(z = 0) = \frac{3H_0^2}{8\pi G} = 1.8788 \times 10^{-29} h^2 \text{ g cm}^{-3}$$

- Energy density today is given as a **fraction of critical**

$$\Omega \equiv \frac{\rho}{\rho_c(z = 0)}$$

- Note that **physical** energy density  $\propto \Omega h^2$  ( $\text{g cm}^{-3}$ )

# Critical Density

- Likewise radius of **curvature** then given by

$$\Omega_K = (1 - \Omega) = -\frac{1}{H_0^2 R^2} \rightarrow R = (H_0 \sqrt{\Omega - 1})^{-1}$$

- If  $\Omega \approx 1$ , then **true density** is near **critical**  $\rho \approx \rho_c$  and

$$\rho_K \ll \rho_c \leftrightarrow H_0 R \ll 1$$

Universe is **flat** across the Hubble distance

- $\Omega > 1$  **positively** curved

$$D_A = R \sin(D/R) = \frac{1}{H_0 \sqrt{\Omega - 1}} \sin(H_0 D \sqrt{\Omega - 1})$$

- $\Omega < 1$  **negatively** curved

$$D_A = R \sin(D/R) = \frac{1}{H_0 \sqrt{1 - \Omega}} \sinh(H_0 D \sqrt{1 - \Omega})$$

# Newtonian Energy Interpretation

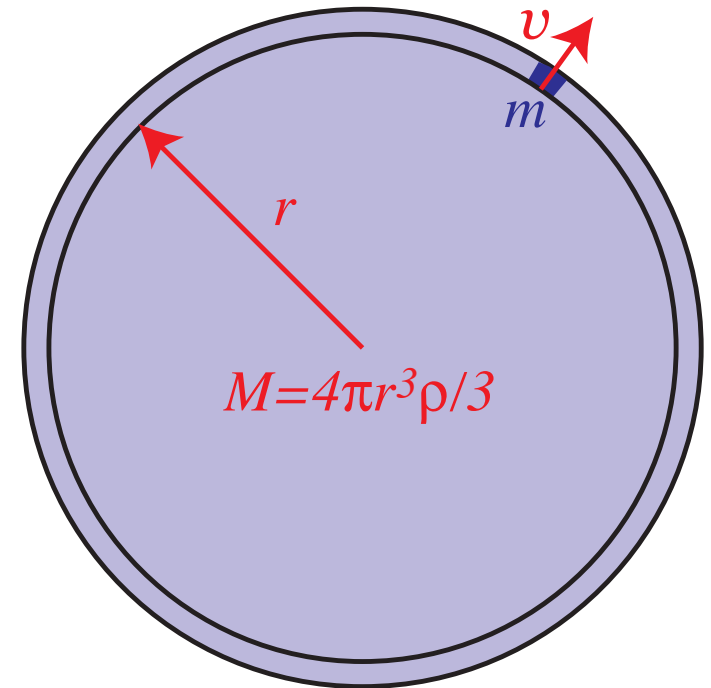
- Consider a **test particle** of mass  $m$  as part of expanding **spherical shell** of radius  $r$  and total mass  $M$ .
- **Energy conservation**

$$E = \frac{1}{2}mv^2 - \frac{GMm}{r} = \text{const}$$

$$\frac{1}{2} \left( \frac{dr}{dt} \right)^2 - \frac{GM}{r} = \text{const}$$

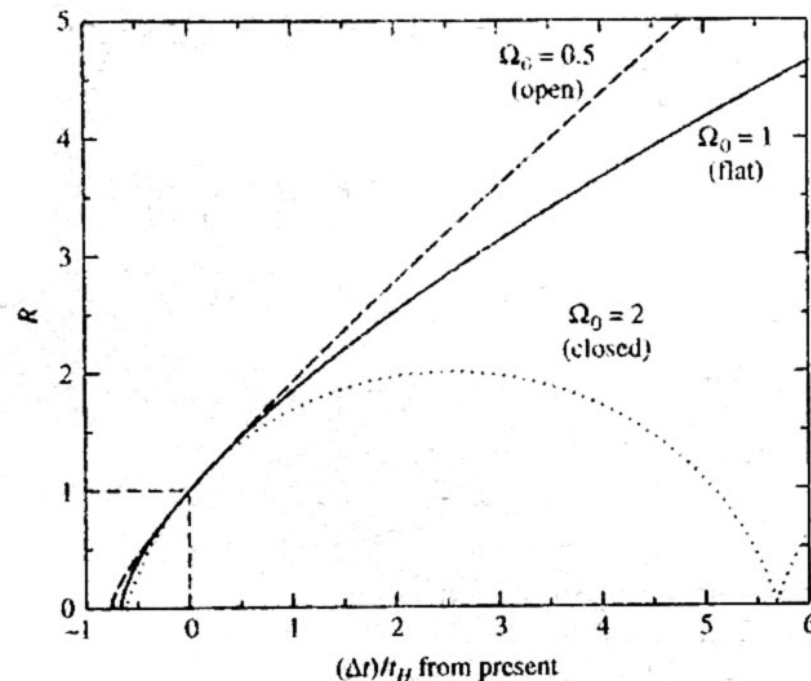
$$\frac{1}{2} \left( \frac{1}{r} \frac{dr}{dt} \right)^2 - \frac{GM}{r^3} = \frac{\text{const}}{r^2}$$

$$H^2 = \frac{8\pi G\rho}{3} - \frac{\text{const}}{a^2}$$



# Newtonian Energy Interpretation

- Constant determines whether the system is **bound** and in the Friedmann equation is associated with **curvature** – not general since **neglects pressure**
- Nonetheless Friedmann equation is the **same** with pressure - but **mass-energy** within expanding shell is **not constant**
- Instead, rely on the fact that **gravity** in the weak field regime is **Newtonian** and **forces** unlike energies are unambiguously defined **locally**.



# Newtonian Force Interpretation

- An alternate, more general Newtonian derivation, comes about by realizing that locally around an observer, **gravity** must look **Newtonian**.
- Given Newton's **iron sphere** theorem, the **gravitational acceleration** due to a spherically symmetric distribution of mass outside some radius  $r$  is zero (**Birkhoff theorem** in GR)
- We can determine the acceleration simply from the **enclosed mass**

$$\nabla^2 \Psi_N = 4\pi G(\rho + 3p)$$

$$\nabla \Psi_N = \frac{4\pi G}{3}(\rho + 3p)r = \frac{GM_N}{r^2}$$

where  $\rho + 3p$  reflects the **active gravitational mass** provided by pressure.

# Newtonian Force Interpretation

- Hence the gravitational acceleration

$$\frac{\ddot{r}}{r} = -\frac{1}{r}\nabla\Psi_N = -\frac{4\pi G}{3}(\rho + 3p)$$

- We'll come back to this way of viewing the effect of the expansion on spherical collapse including the dark energy.

# Conservation Law

- The two **Friedmann equation** are redundant in that one can be derived from the other via **energy conservation**

- (Consequence of **Bianchi** identities in GR:  $\nabla^\mu G_{\mu\nu} = 0$ )

$$d\rho V + p dV = 0$$

$$d\rho a^3 + p da^3 = 0$$

$$\dot{\rho} a^3 + 3 \frac{\dot{a}}{a} \rho a^3 + 3 \frac{\dot{a}}{a} p a^3 = 0$$

$$\frac{\dot{\rho}}{\rho} = -3(1 + w) \frac{\dot{a}}{a}$$

- Time evolution depends on “**equation of state**”  $w(a) = p/\rho$
- If  $w = \text{const.}$  then the **energy density** depends on the scale factor as  $\rho \propto a^{-3(1+w)}$ .



# Multicomponent Universe

- Special cases:
  - nonrelativistic matter  $\rho_m = mn_m \propto a^{-3}$ ,  $w_m = 0$
  - ultrarelativistic radiation  $\rho_r = En_r \propto \nu n_r \propto a^{-4}$ ,  $w_r = 1/3$
  - curvature  $\rho_K \propto a^{-2}$ ,  $w_K = -1/3$
  - (cosmological) constant energy density  $\rho_\Lambda \propto a^0$ ,  $w_\Lambda = -1$
  - total energy density summed over above

$$\rho(a) = \sum_i \rho_i(a) = \rho_c(a=1) \sum_i \Omega_i a^{-3(1+w_i)}$$

- If constituent  $w$  also evolve (e.g. massive neutrinos)

$$\rho(a) = \rho_c(a=1) \sum_i \Omega_i e^{-\int d \ln a 3(1+w_i)}$$

# Multicomponent Universe

- Friedmann equation gives Hubble parameter evolution in  $a$

$$H^2(a) = H_0^2 \sum \Omega_i e^{-\int d \ln a \, 3(1+w_i)}$$

- In fact we can always define a critical equation of state

$$w_c = \frac{p_c}{\rho_c} = \frac{\sum w_i \rho_i - \rho_K/3}{\sum_i \rho_i + \rho_K}$$

- Critical energy density obeys energy conservation

$$\rho_c(a) = \rho_c(a=1) e^{-\int d \ln a \, 3(1+w_c(a))}$$

- And the Hubble parameter evolves as

$$H^2(a) = H_0^2 e^{-\int d \ln a \, 3(1+w_c(a))}$$

# Acceleration Equation

- Time derivative of (first) Friedmann equation

$$\begin{aligned}\frac{dH^2}{dt} &= \frac{8\pi G}{3} \frac{d\rho_c}{dt} \\ 2H \left[ \frac{1}{a} \frac{d^2 a}{dt^2} - H^2 \right] &= \frac{8\pi G}{3} H [-3(1 + w_c)\rho_c] \\ \left[ \frac{1}{a} \frac{d^2 a}{dt^2} - 2 \frac{4\pi G}{3} \rho_c \right] &= -\frac{4\pi G}{3} [3(1 + w_c)\rho_c] \\ \frac{1}{a} \frac{d^2 a}{dt^2} &= -\frac{4\pi G}{3} [(1 + 3w_c)\rho_c] \\ &= -\frac{4\pi G}{3} (\rho + \rho_K + 3p + 3p_K) \\ &= -\frac{4\pi G}{3} (1 + 3w)\rho\end{aligned}$$

- Acceleration equation says that universe decelerates if  $w > -1/3$

# Expansion Required

- Friedmann equations “predict” the expansion of the universe.  
Non-expanding conditions  $da/dt = 0$  and  $d^2a/dt^2 = 0$  require

$$\rho = -\rho_K \quad \rho = -3p$$

i.e. a positive curvature and a total equation of state

$$w \equiv p/\rho = -1/3$$

- Since matter is known to exist, one can in principle achieve this by adding a balancing cosmological constant

$$\rho = \rho_m + \rho_\Lambda = -\rho_K = -3p = 3\rho_\Lambda$$

$$\rho_\Lambda = -\frac{1}{3}\rho_K, \quad \rho_m = -\frac{2}{3}\rho_K$$

Einstein introduced  $\rho_\Lambda$  for exactly this reason – “biggest blunder”; but note that this balance is unstable:  $\rho_m$  can be perturbed but  $\rho_\Lambda$ , a true constant cannot

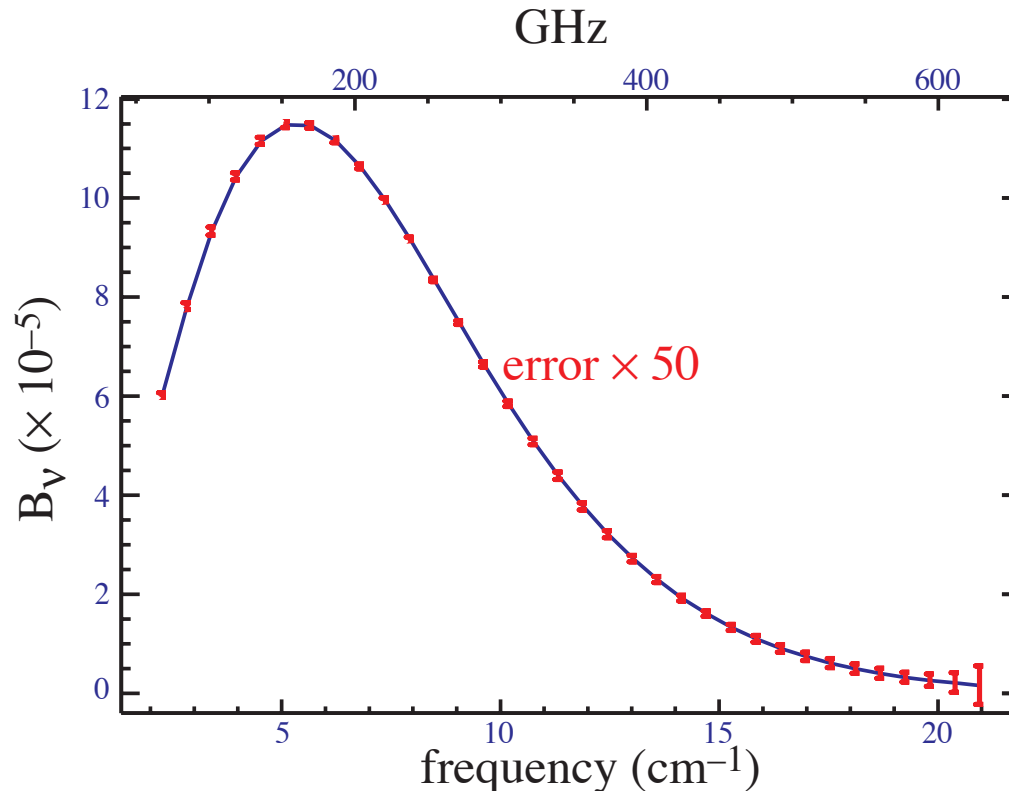
# Cosmic Microwave Radiation

- Existence of a  $\sim 10\text{K}$  radiation background predicted by Gamow and Alpher in 1948 based on the formation of light elements in a hot big bang (BBN)
- Peebles, Dicke, Wilkinson & Roll in 1965 independently predicted this background and proceeded to build instrument to detect it
- Penzias & Wilson 1965 found unexplained excess isotropic noise in a communications antennae and learning of the Peebles et al calculation announced the discovery of the blackbody radiation
- Thermal radiation proves that the universe began in a hot dense state when matter and radiation was in equilibrium - ruling out a competing steady state theory

# Cosmic Microwave Radiation

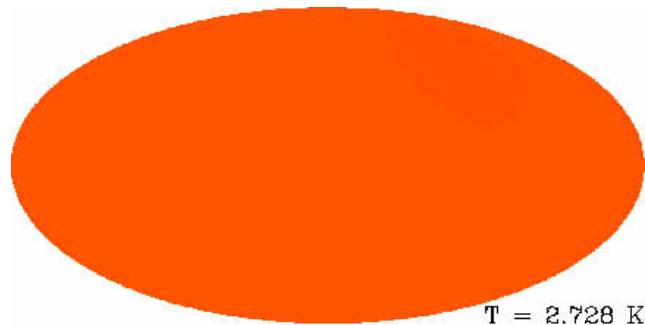
- Modern measurement from COBE satellite of **blackbody spectrum**.

$$T = 2.725\text{K}, \rho_\gamma = (\pi^2/15)T^4 \text{ giving } \Omega_\gamma h^2 = 2.471 \times 10^{-5}$$

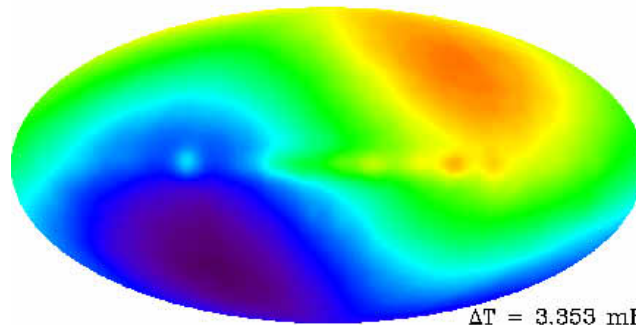


# Cosmic Microwave Radiation

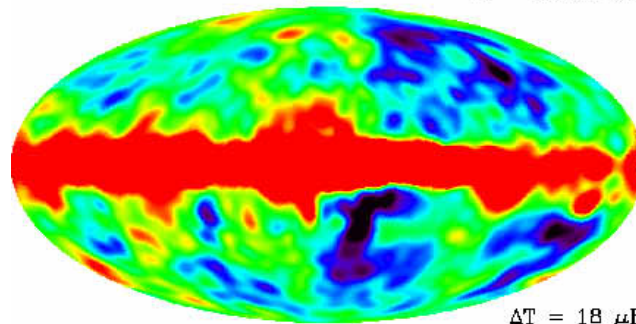
- Radiation is isotropic to  $10^{-5}$  in temperature  $\rightarrow$  horizon problem



$T = 2.728 \text{ K}$



$\Delta T = 3.353 \text{ mK}$



$\Delta T = 18 \text{ } \mu\text{K}$

# Total Radiation

- Adding in **neutrinos** to the radiation gives the **total radiation** (next lecture set) content as  $\Omega_r h^2 = 4.15 \times 10^{-5}$
- Since radiation redshifts faster than matter by one factor of  $1 + z$  even this small radiation content will **dominate** the total energy density at sufficiently **high redshift**
- Matter-radiation **equality**

$$1 + z_{\text{eq}} = \frac{\Omega_m h^2}{\Omega_r h^2}$$

$$1 + z_{\text{eq}} = 3130 \frac{\Omega_m h^2}{0.13}$$



# Dark Matter

- Since **Zwicky** in the 1930's non-luminous or **dark matter** has been known to dominate over luminous matter in stars (and hot gas)
- Arguments based on **internal motion** holding system up against **gravitational force**
- Equilibrium requires a balance **pressure** of **internal motions**
  - rotation velocity** of spiral galaxies
  - velocity dispersion** of galaxies in clusters
  - gas pressure** or **thermal motion** in clusters
  - radiation pressure** in CMB acoustic oscillations

# Classical Argument

- Classical argument for measuring total amount of dark matter
- Assuming that the object is somehow typical in its non-luminous to luminous density: “mass-to-light ratio”
- Convert to dark matter density as  $M/L \times \text{luminosity density}$
- From galaxy surveys the luminosity density in solar units is

$$\rho_L = 2 \pm 0.7 \times 10^8 h L_\odot \text{Mpc}^{-3}$$

( $h$ 's:  $L \propto F d^2$  so  $\rho_L \propto L/d^3 \propto d^{-1}$  and  $d$  in  $h^{-1}$  Mpc)

- Critical density in solar units is

$$\rho_c = 2.7754 \times 10^{11} h^2 M_\odot \text{Mpc}^{-3}$$

so that the critical mass-to-light ratio in solar units is

$$M/L \approx 1400h$$

# Dark Matter: Rotation Curves

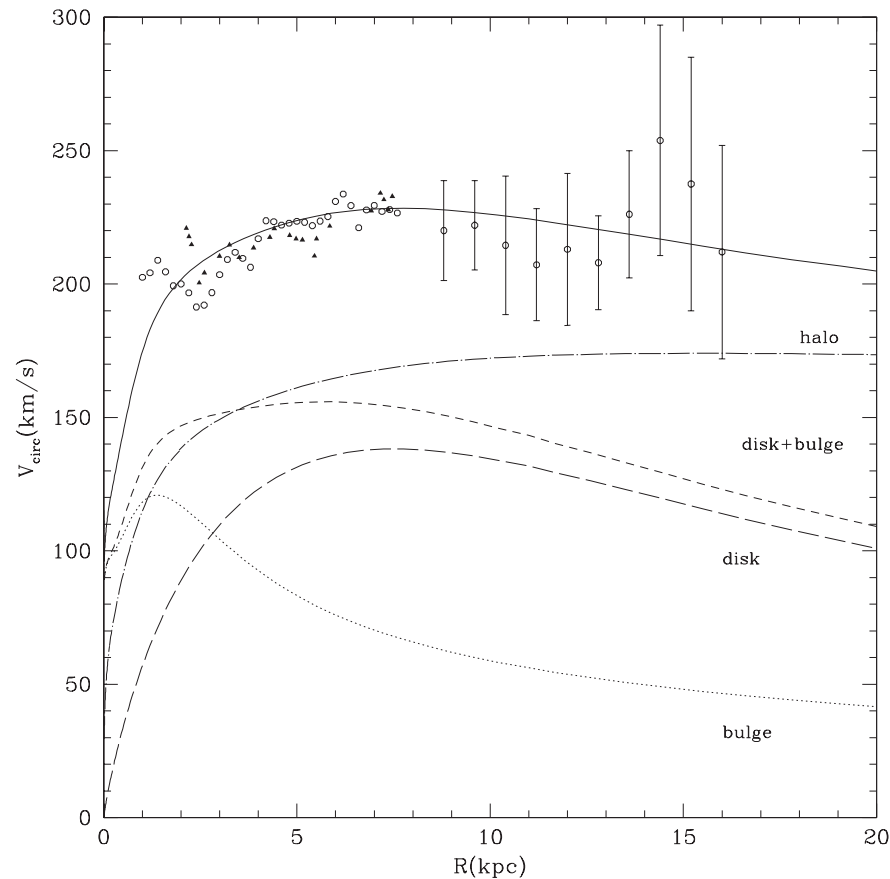
- Flat rotation curves:

$$GM/r^2 \approx v^2/r$$

$$M \approx v^2 r / G$$

so  $M \propto r$  out to tens of kpc

- Implies  $M/L > 30h$   
and perhaps more –  
closure if flat out to  $\sim 1$  Mpc.
- Mass required to keep rotation  
curves flat much larger than implied by stars and gas.
- Hence “dark” matter

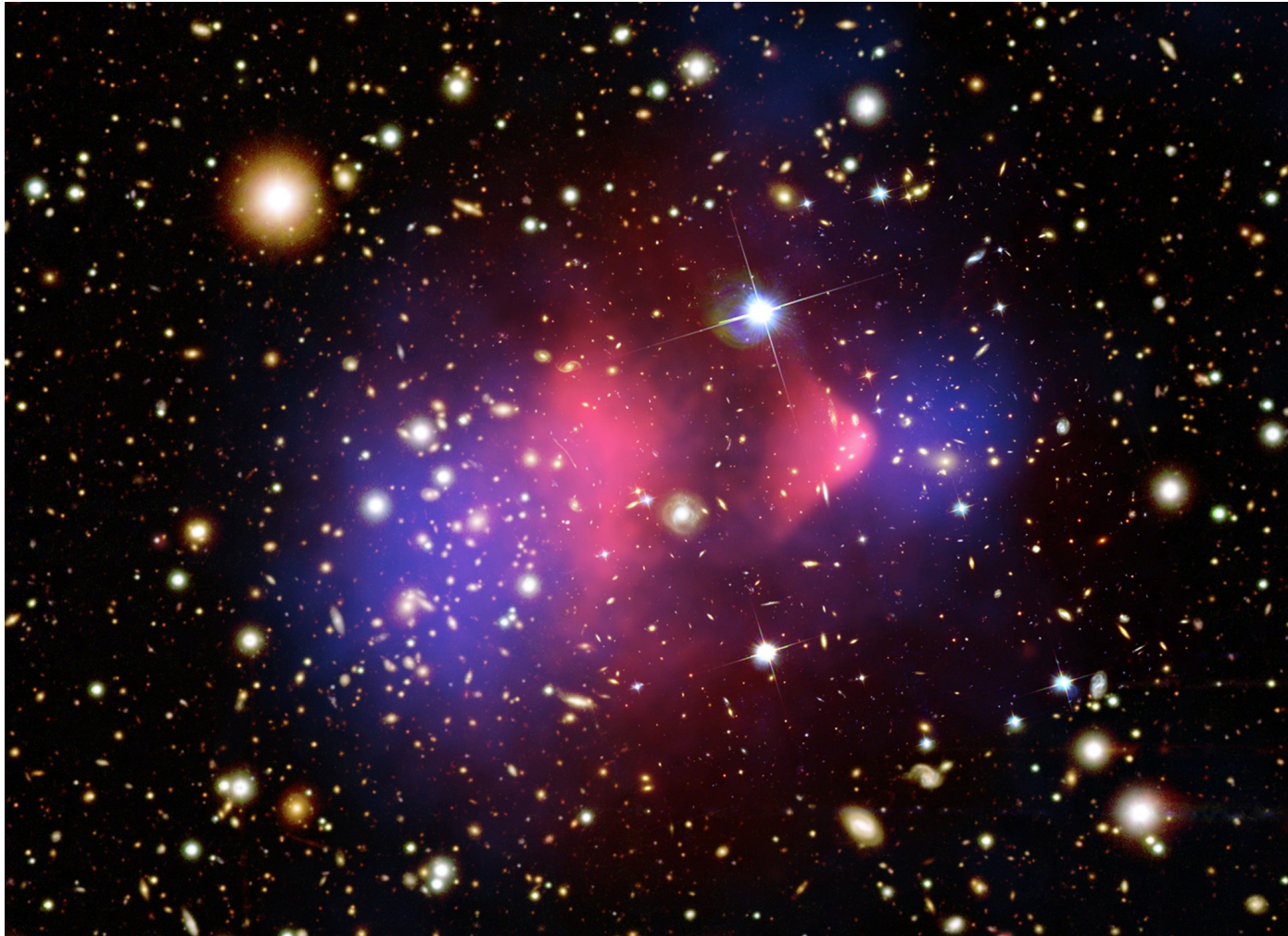


# Dark Matter: Clusters

- Similar argument holds in clusters of galaxies
- Velocity dispersion replaces circular velocity
- Centripetal force is replaced by a “pressure gradient”  $T/m = \sigma^2$   
and  $p = \rho T/m = \rho \sigma^2$
- Zwicky got  $M/L \approx 300h$ .
- Generalization to the gas distribution also gives evidence for dark matter

# Dark Matter: Bullet Cluster

- Merging clusters: gas (visible matter) collides and shocks (X-rays), dark matter measured by gravitational lensing passes through



# Hydrostatic Equilibrium

- Evidence for dark matter in  $X$ -ray clusters also comes from direct hydrostatic equilibrium inference from the gas: balance radial pressure gradient with gravitational potential gradient
- Infinitesimal volume of area  $dA$  and thickness  $dr$  at radius  $r$  and interior mass  $M(r)$ : pressure difference supports the gas

$$[p_g(r) - p_g(r + dr)]dA = \frac{GmM}{r^2} = \frac{G\rho_g M}{r^2}dV$$
$$\frac{dp_g}{dr} = -\rho_g \frac{d\Phi}{dr}$$

with  $p_g = \rho_g T_g / m$  becomes

$$\frac{GM}{r} = -\frac{T_g}{m} \left( \frac{d \ln \rho_g}{d \ln r} + \frac{d \ln T_g}{d \ln r} \right)$$

- $\rho_g$  from X-ray luminosity;  $T_g$  sometimes taken as isothermal

# CMB Hydrostatic Equilibrium

- Same argument in the CMB with radiation pressure in the gas balancing the gravitational potential gradients of linear fluctuations
- Best measurement of the dark matter density to date (Planck 2015):  $\Omega_c h^2 = 0.1188 \pm 0.0010$ ,  $\Omega_b h^2 = (2.23 \pm 0.014) \times 10^{-2}$ .
- Unlike other techniques, measures the physical density of the dark matter rather than contribution to critical since the CMB temperature sets the physical density and pressure of the photons

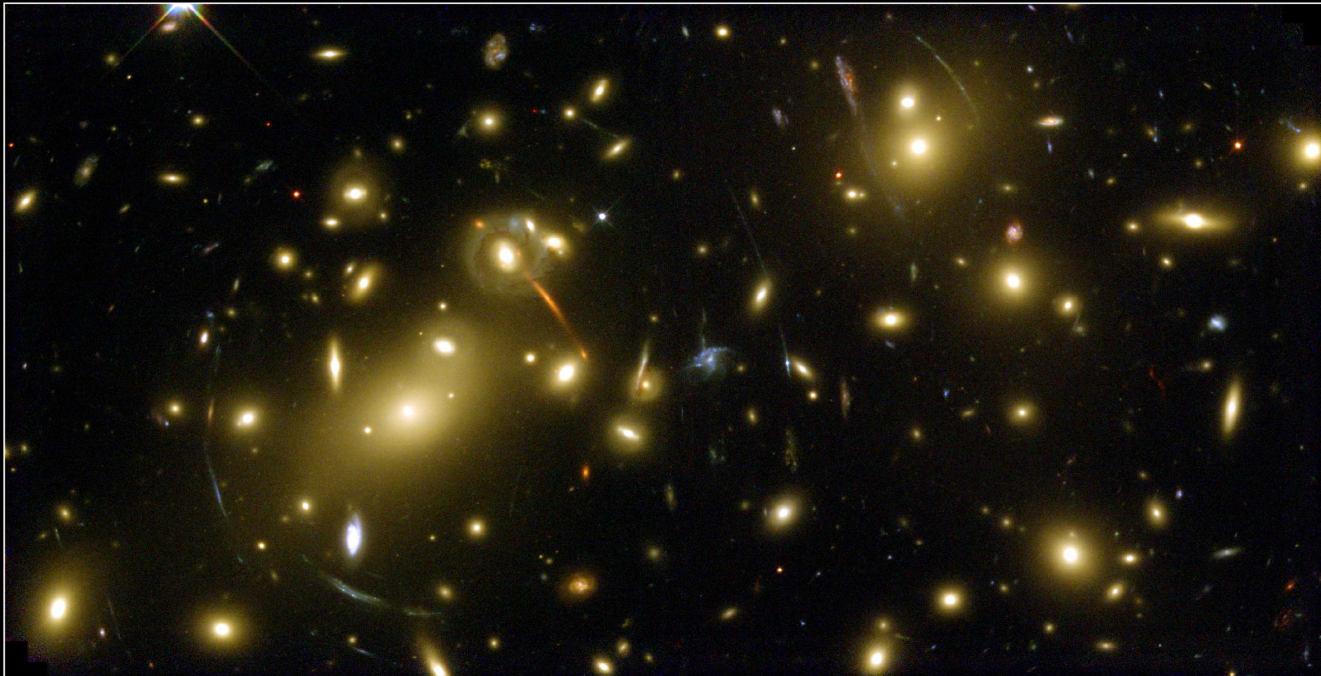
# Gravitational Lensing

- Mass can be directly measured in the gravitational lensing of sources behind the cluster
- Strong lensing (giant arcs) probes central region of clusters
- Weak lensing (1-10%) elliptical distortion to galaxy image probes outer regions of cluster and total mass



# Giant Arcs

- Giant arcs in galaxy clusters: galaxies, source; cluster, lens

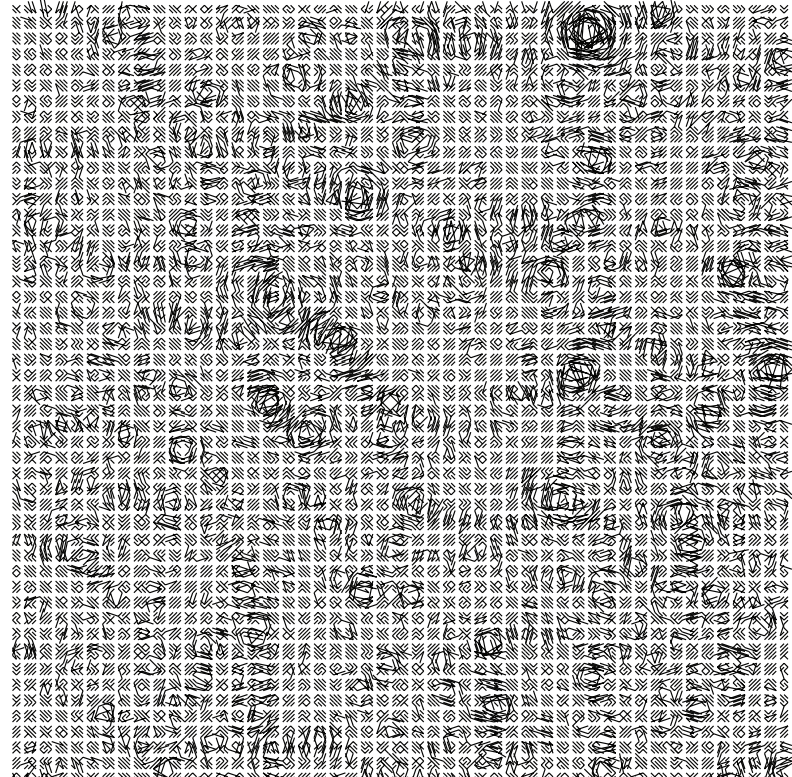
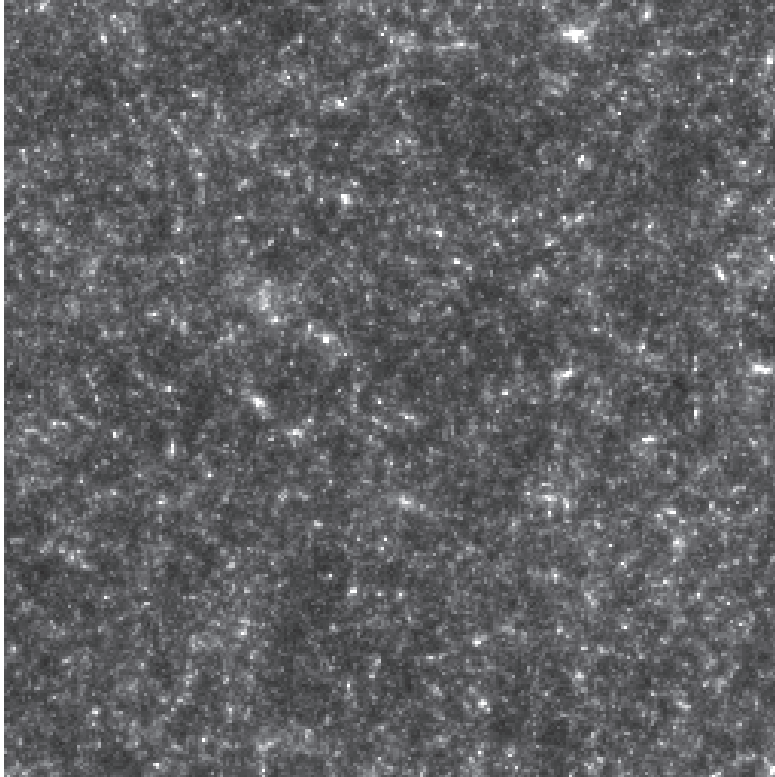


**Galaxy Cluster Abell 2218**  
Hubble Space Telescope • WFPC2

NASA, A. Fruchter and the ERO Team (STScI) • STScI-PRC00-08

# Cosmic Shear

- On even larger scales, the large-scale structure weakly shears background images: weak lensing



# Dark Energy

- Distance redshift relation depends on energy density components

$$H_0 D(z) = \int dz \frac{H_0}{H(a)}$$

- SN dimmer, distance further than in a matter dominated epoch
- Hence  $H(a)$  must be smaller than expected in a matter only  $w_c = 0$  universe where it increases as  $(1+z)^{3/2}$

$$H_0 D(z) = \int dz e^{\int d \ln a \frac{3}{2}(1+w_c(a))}$$

- Distant supernova Ia as standard candles imply that  $w_c < -1/3$  so that the expansion is accelerating
- Consistent with a cosmological constant that is  $\Omega_\Lambda \approx 0.70$
- Coincidence problem: different components of matter scale differently with  $a$ . Why are two components comparable today?

# Cosmic Census

- With  $h = 0.68$  and CMB  $\Omega_m h^2 = 0.14$ ,  $\Omega_m = 0.30$  - consistent with other, less precise, dark matter measures
- CMB provides a test of  $D_A \neq D$  through the standard rulers of the acoustic peaks and shows that the universe is close to flat  $\Omega \approx 1$
- Consistency has lead to the standard model for the cosmological matter budget:
  - 70% dark energy
  - 30% non-relativistic matter (with 84% of that in dark matter)
  - 0% spatial curvature