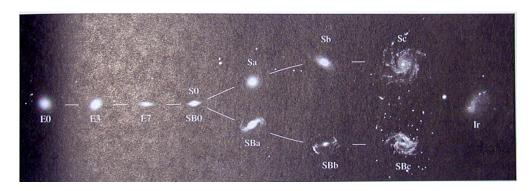
# Set 2: Nature of Galaxies

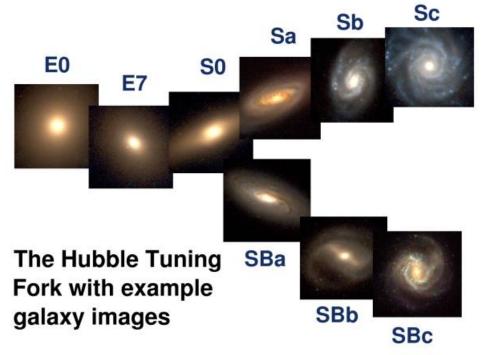
# Great Shapley-Curtis Debate

- History: as late as the early 1920's it was not known that the "spiral nebula" were galaxies like ours
- Debate between Shapley (galactic objects) and Curtis (extragalactic, or galaxies) in 1920 highlighted the difficulties distances in astrophysics difficult to measure - Shapley's inferences based on star counts without extinction and too large a galaxy, novae as standard candles, proper motion
- Hubble in 1923 used Cepheids to establish that Andromeda (M31) is extragalactic at 285kpc - modern measurements say it is 770kpc from the sun.
- Our galaxy is just one of many. Copernican principle in cosmology
  we do not occupy a special place in the universe

# Galaxy Zoology

- Hubble's tuning fork classification of galaxies
- A sequence going from ellipticals En, through regular S0 and barred SB0 lenticulars, to normal S and barred spirals SB ending in irregulars

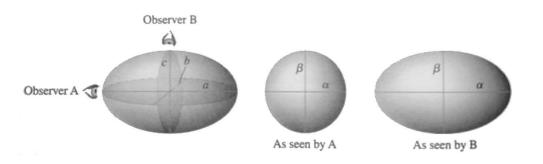




# Galaxy Zoology

• Ellipticals

are further distinguished by the degree of projected ellipticity: the projected major  $\alpha$  and minor  $\beta$  axes



$$\frac{n}{10} = \epsilon \equiv 1 - \beta/\alpha$$

Classification does not necessarily correspond to physical distinctions!

# Galaxy Zoology: Ellipticals

- The actual ellipticity is 3 dimensional and the three axes ordered as
   a ≥ b ≥ c determine the degree of oblateness: a sphere has
   a = b = c, perfectly oblate has a = b, perfectly prolate b = c
- In projection, a strongly prolate or oblate elliptical can have vanishingly small ellipticities
- Ellipticals are often called "early type" and spirals "late type" despite the fact that mergers of spirals can result in ellipticals

# Galaxy Zoology: Ellipticals

- Ellipticals vary widely in physical properties from giants to dwarfs
- Absolute *B* magnitude from -8 to -23
- Total mass from  $10^7 M_{\odot}$  to  $10^{13} M_{\odot}$
- Diameters from few tenths of kpc to hundreds of kpc
- Further classication

cD: high mass, high luminosity, high mass to light, in clusters Normal elliptical: B = -15 to -23,  $M = 10^8 - 10^{13} M_{\odot}$ 

Dwarf ellipticals: low surface brightness for a given B = -13 to  $-19, M = 10^7 - 10^9 M_{\odot}$ 

Dwarf spheroidal: extremely low luminosity B = -8 to -15 and surface brightness can only be detected locally

Blue compact dwarf: small with vigorous star formation B = -14 to -17 and  $M \sim 10^9$ .

# Galaxy Zoology: Spiral NGC4414



# Galaxy Zoology: Spirals

- Spirals are subdivided a, ab, b, bc, c in order of bulge prominance, tightly wound spiral arms, smoothest distribution of stars
- The presence of a central bar is indicated with *B*
- Milky Way is a SBbc, M31 is an Sb
- S(B)a c smaller range of physical properties compared with ellipticals (table)

TABLE 25.1 Characteristics of Early Spiral Galaxies.

	Sa	Sb	Sc
M <sub>B</sub>	-17 to -23	-17 to -23	-16 to -22
$M (M_{\odot})$	109-1012	109-1012	10 <sup>9</sup> -10 <sup>12</sup>
$\langle L_{\rm bulge}/L_{\rm total} \rangle_{\rm B}$	0.3	0.13	0.05
Diameter (D25, kpc)	5-100	5-100	5-100
$\langle M/L_B \rangle (M_{\odot}/L_{\odot})$	$6.2 \pm 0.6$	$4.5 \pm 0.4$	$2.6 \pm 0.2$
$\langle V_{\rm max} \rangle$ (km s <sup>-1</sup> )	299	222	175
$V_{\rm max}$ range (km s <sup>-1</sup> )	163367	144330	99-304
pitch angle	$\sim 6^{\circ}$	$\sim 12^{\circ}$	$\sim 18^{\circ}$
$\langle B-V\rangle$	0.75	0.64	0.52
(Mgas/Mtotal)	0.04	0.08	0.16
$\left(M_{\rm H_2}/M_{\rm H~I}\right)$	$2.2 \pm 0.6$ (Sab)	$1.8 \pm 0.3$	$0.73 \pm 0.13$
$\langle S_N \rangle$	$1.2 \pm 0.2$	$1.2 \pm 0.2$	$0.5 \pm 0.2$

TABLE 25.2 Characteristics of Late Spiral and Irregular Galaxies.

	Sd/Sm	Im/Ir
MB	-15 to -20	-13 to -18
<i>M</i> (M <sub>☉</sub> )	10 <sup>8</sup> -10 <sup>10</sup>	10 <sup>8</sup> -10 <sup>10</sup>
Diameter $(D_{25}, kpc)$	0.5-50	0.5-50
$\langle M/L_B \rangle (M_{\odot}/L_{\odot})$	$\sim 1$	$\sim 1$
$V_{\rm max}$ range (km s <sup>-1</sup> )	80-120	50-70
(B-V)	0.47	0.37
$(M_{\rm gas}/M_{\rm total})$	0.25 (Scd)	0.5-0.9
$\left(M_{\rm H_2}/M_{\rm H~I}\right)$	0.03-0.3	~ 0
$\langle S_N \rangle$	$0.5 \pm 0.2$	$0.5\pm0.2$

# Galaxy Zoology: Irregulars

- Irregulars classed as *IrrI* if there is any organized structure such as spiral arms
- Otherwise *IrrII* otherwise
- Examples: Large Magellanic Clouds (LMC) is *IrrI* and Small Magellanic Clouds (SMC) is *IrrII*
- Physical properties: tend to be small and faint
- Absolute B magnitude from -13 to -20
- Masses from  $10^8 M_{\odot}$  to  $10^{10} M_{\odot}$

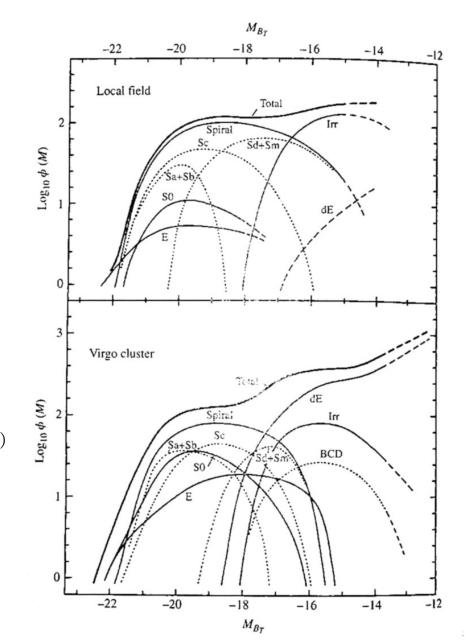
## Galaxy Properties: Luminosity Function

#### • Abundance

as a function of luminosity is called the "luminosity function". Number of galaxies in dL around L and has a rough shape of a "Schechter function"

$$\phi_L dL \propto L^{\alpha} e^{-L/L_*} dL$$
  

$$\phi_M dM \propto 10^{-0.4(\alpha+1)M} e^{-10^{0.4(M_*-M)}}$$
  
with  $\alpha \approx -1$   
and  $L_*$  from  $M_* = -21$  in  $B$ 



# Galaxy Properties: Luminosity Function

- Luminosity function is to galaxies what the distribution in magnitudes of stars is to star counts
- Galaxy counts probe the galaxy number density as a function of angular position (and redshift) to a limiting magnitude (a "redshift" survey)
- Luminosity function (determined locally) tells you how to interpret the observed counts in terms of a 3D distribution of galaxies

# Galaxy Properties: Surface Brightness

- Surface brightness profile defines the effective scale of the bulge and disk components
- Surface brightness  $\mu$  measured in *B*-mag arcsec<sup>-2</sup>
- Define  $r_e$  as the radius within which 1/2 the light emitted.
- Bulges of spirals and ellipticals follow a Sersic profile where the surface brightness in mag scales as a power law at  $r \gg r_e$

$$\mu(r) = \mu_e + 8.3268 \left[ \left(\frac{r}{r_e}\right)^{1/n} - 1 \right]$$

where n = 4 is the de Vaucouleurs profile and  $\mu_e$  is the surface brightness at  $r_e$ 

### Galaxy Properties: Surface Brightness

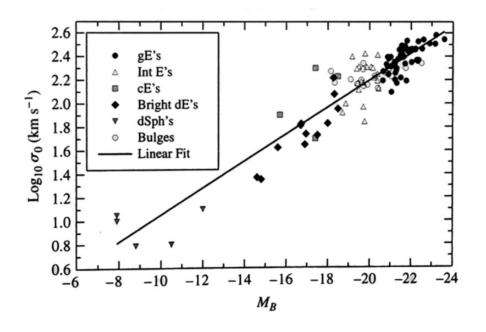
• Disks follow an exponential which in mag scales as

$$\mu(r) = \mu_0 + 1.09 \left(\frac{r}{h_r}\right)$$

where  $h_r$  is the characteristic scale length

### Galaxy Properties: Fundamental Plane

- Faber Jackson correlation between luminosity and velocity dispersion of stars (measured from the width of lines from aggregate unresolved stars)  $L \propto \sigma_0^4$
- Expected if mass to light and surface brightness a constant.
   Consider virial theorem



$$-2\langle K\rangle = \langle U\rangle, \qquad -2\sum_{i}^{N} \frac{1}{2}m_{i}v_{i}^{2} = U$$

7 7

#### Galaxy Properties: Fundamental Plane

• Simplify as equal mass objects composing M

$$-\frac{m}{N}\sum_{i}^{N}v_{i}^{2}=\frac{U}{N}$$

• Sum is the average  $v^2$  and is an observable assuming that radial velocities reflect total  $\langle v^2 \rangle = 3 \langle v_r^2 \rangle \equiv 3\sigma_r^2$ 

$$-3m\sigma_r^2 = \frac{U}{N}$$

• Potential energy for a constant density spherical distribution of mass M = Nm and radius R

$$\frac{U}{N} = -\frac{3}{5} \frac{GM^2}{NR}, \qquad M_{\rm vir} = \frac{5R\sigma_r^2}{G}$$

# Galaxy Properties: Fundamental Plane

• Eliminate R by assuming constant surface brightness  $L/R^2 = C_{SB}$  eliminate  $R = (L/C_{SB})^{1/2}$ 

$$M_{\rm vir} = \frac{5\sigma_r^2}{G} \left( L/C_{SB} \right)^{1/2}$$

• Eliminate  $M_{\text{vir}}$  by assuming constant mass to light  $M/L = 1/C_{ML}$ 

$$L = C_{ML} \frac{5\sigma_r^2}{G} \left( L/C_{SB} \right)^{1/2}$$
$$L \propto \sigma_r^4$$

• A tighter relation is obtained by introducing a second observable - either the effective radius

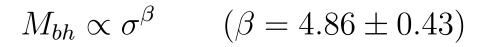
$$L \propto \sigma_r^{2.65} r_e^{0.65}$$

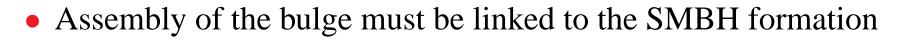
which defines the fundamental plane of ellipticals

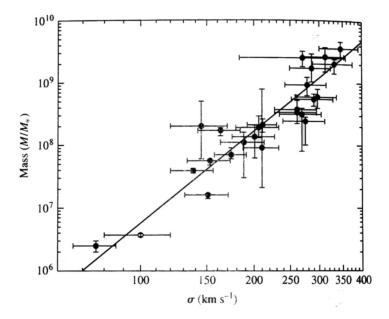
# Galaxy Properties: SMBH

- A similar argument

   is used to measure the mass
   of the central black hole from the
   velocity dispersion of stars around
   it in both spirals and ellipticals
- The inferred mass is also correlated with the velocity dispersion much further out in the bulge

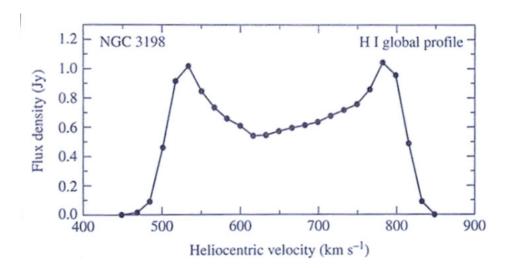






# Galaxy Properties: $v_{max}$

- The maximum velocity in a rotation curve is a robust observable
- The 21 cm line of the disk as a whole reflects the Doppler shifts of the HI participating in the rotation
- Line has a double peaked profile with the peaks near  $v_{\max}$  since much of the gas is in the flat part of the rotation curve near the peak

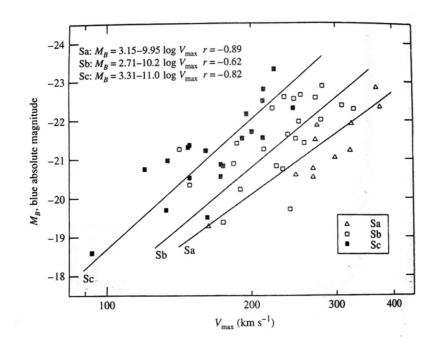


#### Galaxy Properties: Tully Fisher relation

 Correcting for the inclination from the observed radial velocity v<sub>r</sub>

$$\frac{\Delta\lambda}{\lambda_{\rm rest}} = \frac{v_r}{c} = \frac{v_{\rm max}}{c}\sin i$$

• Tully and Fisher established that  $v_{\rm max}$  is correlated with *B* band luminosity as approximately  $L_B \propto v_{\rm max}^4$ 



# Galaxy Properties: Tully Fisher relation

- Tully-Fisher relationship is expected if galaxies have a constant mass to light ratio and constant surface brightness
- Enclosed mass

$$M = \frac{v_{\max}^2 R}{G}$$

• Mass to light ratio  $M/L = 1/C_{ML}$ 

$$L = C_{ML} \frac{v_{\max}^2 R}{G}$$

• Surface brightness  $L/R^2 = C_{SB}$  eliminate  $R = (L/C_{SB})^{1/2}$ 

$$L = C_{ML} \frac{v_{\max}^2}{G} \left(\frac{L}{C_{SB}}\right)^{1/2}$$

$$L \propto v_{\rm max}^4$$

# Galaxy Properties: Tully Fisher relation

• In absolute magnitude

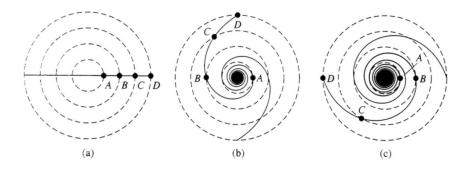
$$M_B = -2.5 \log_{10} L_B + \text{const}$$

$$M_B = -2.5 \log_{10} v_{\rm max}^4 + {\rm const}$$

 $M_B = -10\log_{10}v_{\rm max} + {\rm const}$ 

- Tully Fisher relation is even tighter in IR bands such as *H* band less extinction and late type giant stars are better tracers of overall luminoisity
- Tully Fisher relation can be used to measure distances: measure  $v_{\text{max}}$ , infer absolute magnitude and compare to apparent magnitude

• Winding problem: if spiral structure were physical structures, a flat rotation curve would cause the arms to wind up tightly



- Lin-Shu density wave theory: spiral arms are quasistatic density waves - bunching is like cars in a traffic jam
- Stars pass through the wave/jam and do not cause a winding problem

Consider the orbital motion of a star in cylindrical coordinates
 (R, φ, z) where z is the coordinate out of the disk

$$\frac{d^2\mathbf{r}}{dt^2} = -\nabla\Phi$$

where  $\Phi$  is the gravitational potential.

• Assuming axial symmetry for the potential this yields 3 equations for the three directions

$$\ddot{R} - R\dot{\phi}^2 = -\frac{\partial\Phi}{\partial R}$$
$$\frac{1}{R}\frac{\partial(R^2\dot{\phi})}{\partial t} = 0$$
$$\ddot{z} = -\frac{\partial\Phi}{\partial t}$$

 $\partial z$ 

• Second equation says that there is no force in the azimuthal direction or torque  $\tau = \mathbf{r} \times \mathbf{F}$ 

$$L_z = MRv_\phi = MR^2\dot{\phi} = \text{const}$$

where M is the mass of the star. Defining  $J_z = L_z/M = R^2 \dot{\phi}$  the angular momentum per unit mass

$$R\dot{\phi}^2 = \frac{J_z^2}{R^3}$$

• Radial equation becomes

$$\ddot{R} = -\frac{\partial \Phi}{\partial R} + \frac{J_z^2}{R^3}$$

• The second term is an angular momentum barrier against radial infall or equivalently the centripetal acceleration required to keep R constant  $v_{\phi}^2$ . It can be absorbed into an effective potential

$$\Phi_{\rm eff} = \Phi + \frac{J_z^2}{2R^2}$$

so that the equations of motion becomes  $(J_z \text{ is a constant in } z$ 

$$\ddot{R} = -\frac{\partial \Phi_{\text{eff}}}{\partial R}$$
$$\ddot{z} = -\frac{\partial \Phi_{\text{eff}}}{\partial z}$$

• Structure of  $\Phi_{\text{eff}}(R, z)$  determines motion. In z minimum is at the midplane. In R, minimum forms from the competition of gravity and angular momentum

• Minimum found by seeing where slope vanishes (or equivalently where the gravitational and centripetal acceleration match)

$$\frac{\partial \Phi_{\text{eff}}}{\partial R} = \frac{\partial \Phi}{\partial R} - \frac{J_z^2}{R^3} = 0$$

• Orbits near this stable minimum m oscillate around it:  $\rho \equiv R - R_m$ 

$$\Phi_{\rm eff} \approx \Phi_{\rm eff,m} + \frac{1}{2}\kappa^2\rho^2 + \frac{1}{2}\nu^2 z^2$$

where  $\kappa^2 = \partial^2 \Phi_{\text{eff}} / \partial R^2 |_m$  and  $\nu^2 = \partial^2 \Phi_{\text{eff}} / \partial z^2 |_m$ 

• Equations of motion

$$\ddot{\rho} = -\kappa^2 \rho$$

$$\ddot{z} = -\nu^2 z$$

• Star executes simple harmonic motion around minimum:

 $\rho(t) = A_R \sin \kappa t$ 

$$z(t) = A_z \sin(\nu t + \zeta)$$

where  $\zeta$  is a phase factor and we have defined t = 0 to eliminate the other phase factor

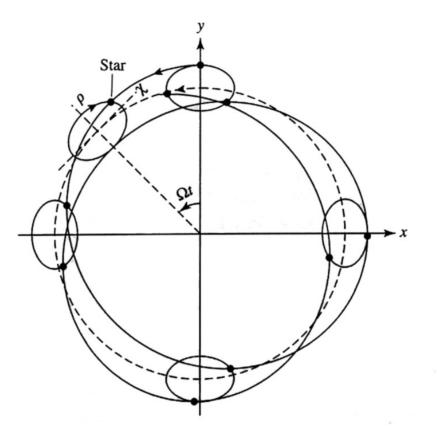
• Azimuthal coordinate given in terms of radial motion

$$\dot{\phi} = \frac{J_z}{R^2} \approx \frac{J_z}{R_m^2} \left( 1 - 2\frac{\rho(t)}{R_m} \right)$$

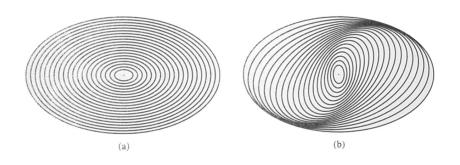
$$\phi(t) = \phi_0 + \Omega t + \frac{2\Omega}{\kappa R_m} A_R \cos \kappa t$$

where the unperturbed angular frequency  $\Omega = J_z/R_m^2$ 

- Star executes epicyclic motion or rosette
- κ also known
   as epicyclic frequency
- Relative to the unperturbed orbit (corrotating with the local angular speed Ω, star executes a simple retrograde closed orbit around R<sub>m</sub>



 If the epicyclic frequency κ/Ω = m/n integer ratio then the orbit is closed in the fixed frame: star executes m epicycles during n orbits



• More generally, can always go into a rotating frame "local pattern speed"  $\Omega_{lp}$  where this condition is true and orbits are closed

$$m(\Omega - \Omega_{lp}) = n\kappa$$

 An (n = 1, m = 2) is shown for a case where Ω<sub>lp</sub> is independent of R: if axis of orbit ovals are aligned then bar structure, if rotated then a two armed bar.

- Only pattern is stationary stars are continuously orbiting and piling up in the arms
- Non-constancy of the  $\Omega_{lp}$  will still cause winding but of the pattern and typically at a slower rate for (1, 2).
- Where the local pattern speed matches the global pattern speed Lindblad resonances occur where the epicyclic amplitude increases due to forcing from the local density enhancement - can destroy spiral pattern.
- N-body simulations show formation of transient m = 2 arm patterns and long lived bar instability.