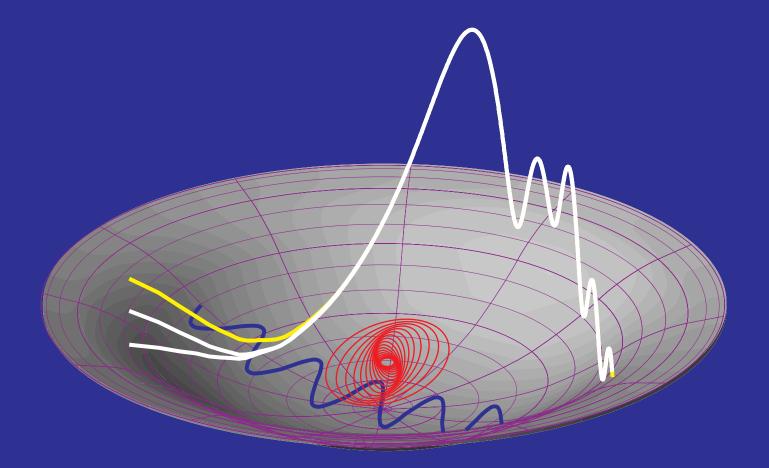
#### Toward a Parameterized Post Friedmann



### Description of Cosmic Acceleration Wayne Hu NYU, November 2006

## **Cosmic Acceleration**

• Cosmic acceleration, like the cosmological constant, can either be viewed as arising from

Missing, or dark energy, with  $w \equiv \bar{p}/\bar{\rho} < -1/3$ 

Modification of gravity on large scales

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- Proof of principle models for both exist: quintessence, k-essence; DGP braneworld acceleration, f(R) modified action
- Compelling models for either explanation lacking
- Dark energy parameterized description on small scales: smooth component with a w(z) that completely defines expansion history
- Parameterized description of modified gravity acceleration?
- Previous ad-hoc attempts violate basic principles like energy-momentum conservation

## Outline

 Constraints imposed by energy-momentum conservation on linear metric fluctuations around an FRW background with a given expansion history

• "Parameterized Post Friedmann" (PPF) description of remaining degrees of freedom

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# Outline

- Constraints imposed by energy-momentum conservation on linear metric fluctuations around an FRW background with a given expansion history
- "Parameterized Post Friedmann" (PPF) description of remaining degrees of freedom
- DGP braneworld acceleration example
- f(R) modified action example
- Collaborators:
  - Dragan Huterer
  - Iggy Sawicki
  - Yong-Seon Song
  - Kendrick Smith

**PPF** Framework

- Parallel treatment of parameterized dark energy beyond a quintessence scalar field
- Demand that the model satisfies (Hu 1998)
  - Given Background Expansion
  - Gauge Invariance
  - **Energy-Momentum Conservation**

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- Demand that the model satisfies (Hu 1998)
  - Given Background Expansion
  - Gauge Invariance
  - **Energy-Momentum Conservation**
  - and the phenomenogically desirable property that the dark energy does not cluster with the dark matter  $\rightarrow$  sound horizon
- Larger scales: energy-momentum conservation requires conservation of the comoving curvature (Bardeen 1980)
- Smaller scales: dark energy spatial perturbations negligible and observable phenomena depend on expansion history only

• Implement with a parameterized model: the sound speed in the dark energy rest frame. Quintessence sound speed  $c_s = 1$ 

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- Parameterization later shown to describe k-essence with modified scalar field kinetic term (Garriga & Mukhanov 1999)

$$\mathcal{L} = F(X, \phi)$$
  $X = -\frac{1}{2} \nabla^{\mu} \phi \nabla_{\mu} \phi$ 

with a sound speed

$$c_s^2 = \frac{\partial F/\partial X}{2(\partial^2 F/\partial X^2)X + (\partial F/\partial X)}$$

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• Beyond single scalar fields: parameterize multiple internal degrees of freedom to allow an evolution across w = -1 phantom divide (Hu 2004)

- Modified gravity models of acceleration
- Demand that the model satisfies
  - Given Background Expansion History
  - Bianchi Identities / (FRW) Metric Theory
  - **Energy-Momentum Conservation**
  - and that modifications reach quasi-static Newtonian limit on small scales: time derivatives neglected compared with spatial gradients
- PPF description can be used to test general relativity on cosmological scales and distinguish modified gravity from smooth dark energy

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  - and that modifications reach quasi-static Newtonian limit on small scales: time derivatives neglected compared with spatial gradients
- PPF description can be used to test general relativity on cosmological scales and distinguish modified gravity from smooth dark energy
- In addition non-linear effects must bring gravity stably back to general relativity on small scales to satisfy solar system tests.
  Beyond the scope of this talk.

- On superhorizon scales, energy momentum conservation and expansion history constrain the evolution of metric fluctuations (Bertschinger 2006)
- For adiabatic perturbations in a flat universe, conservation of comoving curvature applies  $\zeta' = 0$  where  $' \equiv d/d \ln a$  (Bardeen 1980)

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- Gauge transformation to Newtonian gauge

$$ds^{2} = -(1+2\Psi)dt^{2} + a^{2}(1+2\Phi)dx^{2}$$

yields (Hu & Eisenstein 1999)

$$\Phi'' - \Psi' - \frac{H''}{H'} \Phi' - \left(\frac{H'}{H} - \frac{H''}{H'}\right) \Psi = 0$$

• Modified gravity theory supplies the closure relationship between  $\Phi$  and  $\Psi$  and expansion history  $H = \dot{a}/a$  supplies rest.

A Worked Example: DGP Braneworld Acceleration

# A Worked Example: DGP Gravity

• Braneworld acceleration (Dvali, Gabadadze & Porrati 2000)

$$S = \int d^5x \sqrt{-g} \left[ \frac{{}^{(5)}R}{2\kappa^2} + \delta(\chi) \left( \frac{{}^{(4)}R}{2\mu^2} + \mathcal{L}_m \right) \right]$$

with crossover scale  $r_c = \kappa^2/2\mu^2$ 

- Influence of bulk through Weyl tensor anisotropy solve master equation in bulk (Deffayet 2001; see also Sawicki's talk)
- Matter still minimally coupled and conserved
- Satisfies PFF requirements

# A Worked Example: DGP Gravity

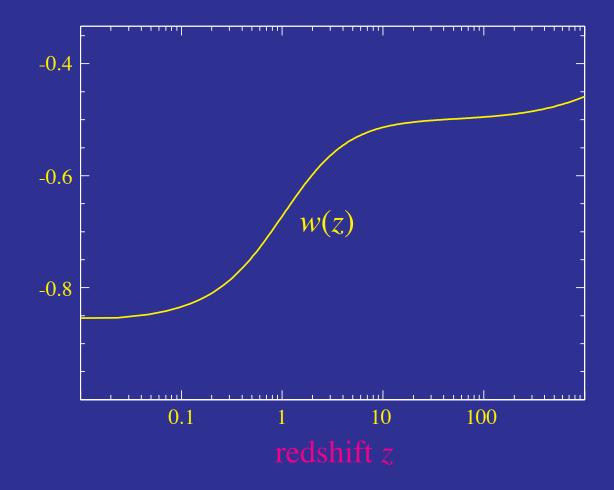
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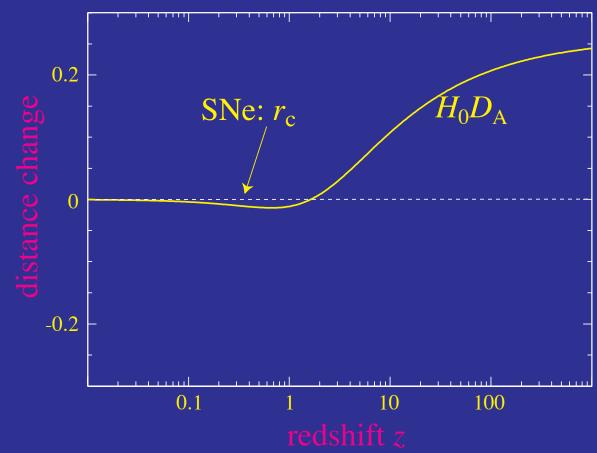
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- Influence of bulk through Weyl tensor anisotropy solve master equation in bulk (Deffayet 2001; see also Sawicki's talk)
- Matter still minimally coupled and conserved
- Satisfies PFF requirements
- Dominance of Weyl tensor anisotropy over other components and matter sets closure relation during self acceleration  $\Psi \to \Phi$
- Transition to this limit leads to enhancement of potential decay and large angle CMB anisotropy

- Matching the DGP expansion history to a dark energy model with the same expansion history
- Effective equation of state w(z) [ $w_0 \sim -0.85$ ,  $w_a \sim 0.35$ ]

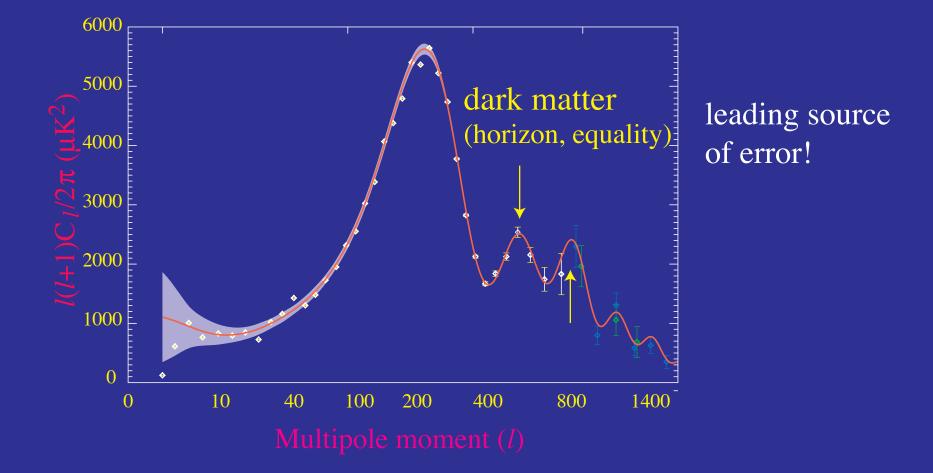


• Crossover scale  $r_c$  fit to SN relative distance from z=0:  $H_0D_A$ 



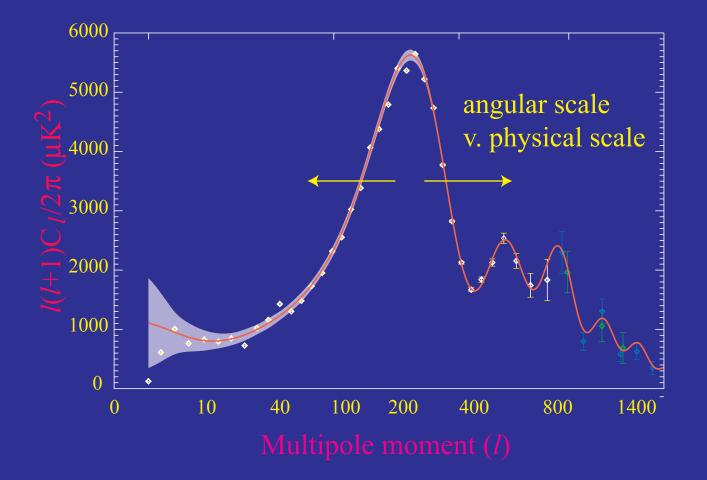
# Leveraging the CMB

• Relative heights of the first 3 peaks calibrates sound horizon and matter radiation equality horizon: measures  $\Omega_m h^2$  currently 8%

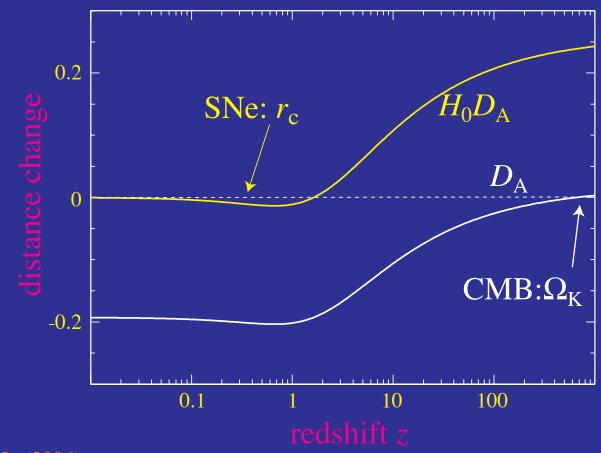


#### Standard Ruler

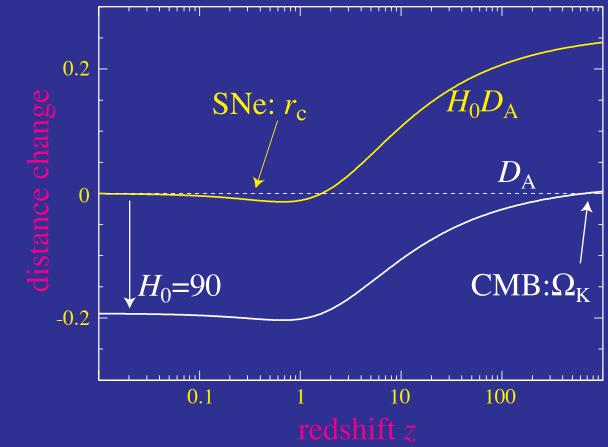
• Standard ruler used to measure the angular diameter distance to recombination (z~1100; currently 2%) or any redshift for which acoustic phenomena observable



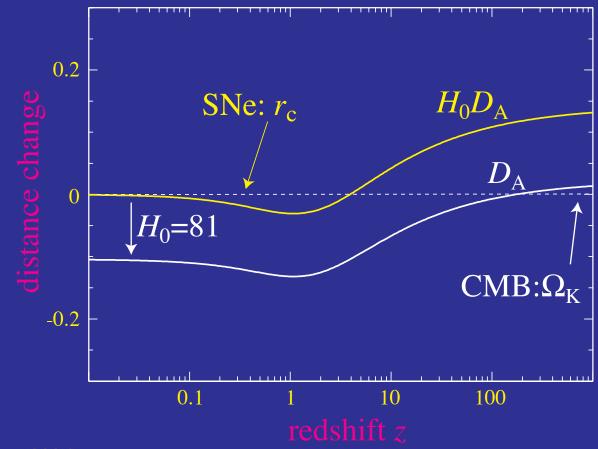
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- Crossover scale  $r_c$  fit to SN relative distance from z=0:  $H_0D_A$
- Mismatch to CMB absolute distance  $D_A$  requires curvature
- Difference in expansion history appears as a change in local distances or the Hubble constant:  $H_0$

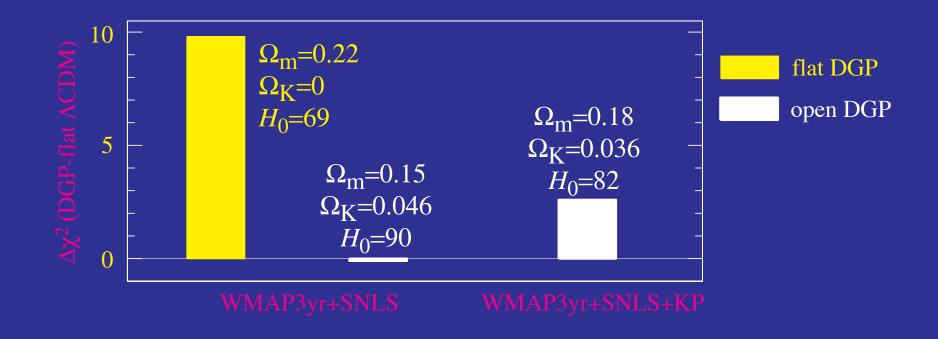


- Crossover scale  $r_c$  fit to SN relative distance from z=0:  $H_0D_A$
- Mismatch to CMB absolute distance  $D_A$  requires curvature
- Compromise between SN and *H*<sub>0</sub> measures



# **DGP** Example

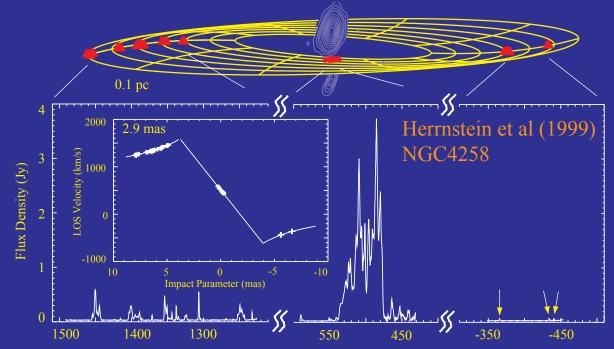
- DGP modified gravity is in tension with distance measures alone: CMB & SNe distances cannot be jointly satisfied in a flat universe
- Even fitting out curvature, Hubble constant is too high for Key Project measurement (and baryon oscillations)
- Joint maximization leads to a poorer fit even with extra curvature parameter



Song, Sawicki & Hu (2006) [also Fairbairn & Goobar 2005, Maartens & Majerotto 2006]

### Prospects for Percent H<sub>0</sub>

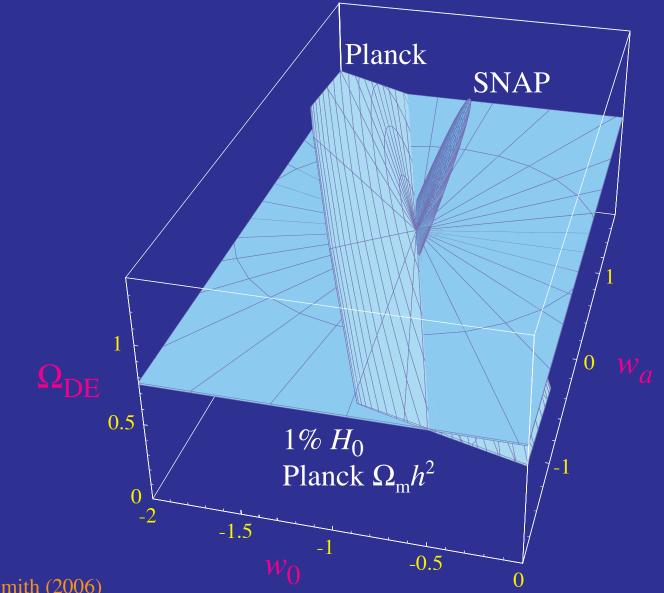
- Improving the distance ladder (~3-5%, Riess 2005; Macri et al 2006)
- Water maser proper motion, acceleration (~3%, VLBA Condon & Lo 2005; ~1% SKA, Greenhill 2004)



- Gravity wave sirens (~2% 3x Adv. LIGO + GRB sat, Dalal et al 2006)
- Combination of dark energy tests: e.g. SNIa relative distances:  $H_0D(z)$  and baryon acoustic oscillations D(z)

#### Flat Universe Precision

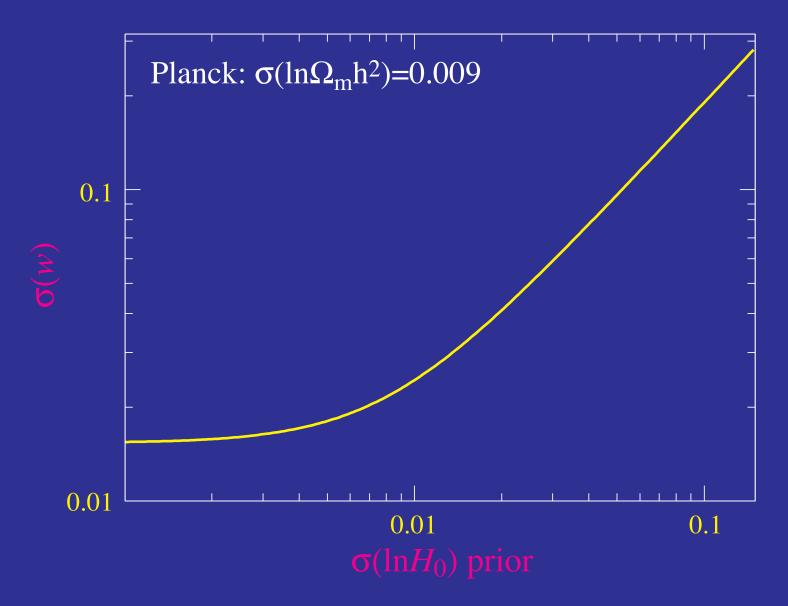
• Planck acoustic peaks,  $1\% H_0$ , SNAP SNe to z=1.7 in a flat universe



Hu, Huterer & Smith (2006)

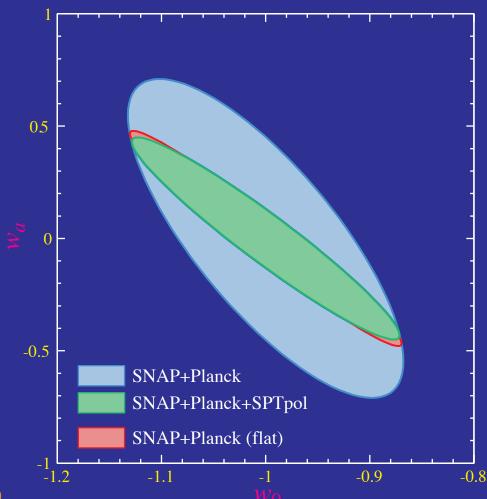
#### Forecasts for $CMB+H_0$

• To complement CMB observations with  $\Omega_m h^2$  to 1%, an  $H_0$  of ~1% enables constant *w* measurement to ~2% in a flat universe



## Dark Energy Equation of State

- Marginalizing curvature degrades 68% CL area by 4.8
- CMB lensing information from SPTpol (~3% B-mode power) fully restores constraints

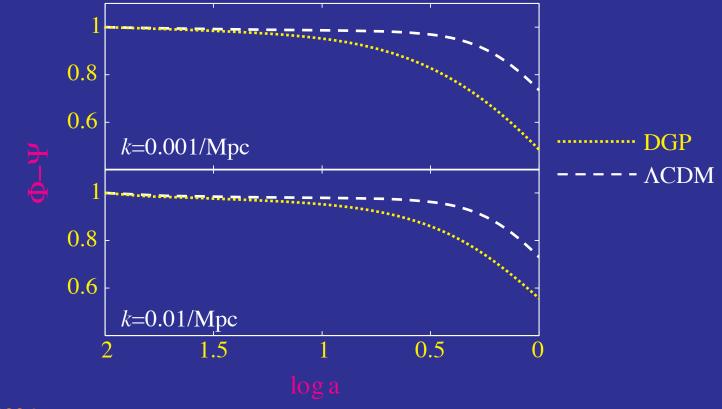


Hu, Huterer & Smith (2006)

### **DGP** Metric Evolution

#### **DGP** Potential Evolution

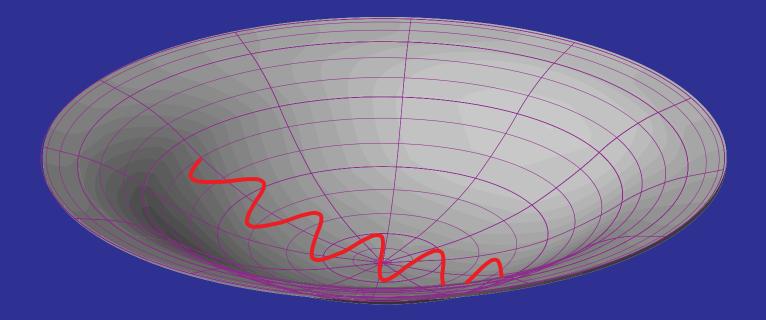
- Difference in expansion history gives excess decay of grav. potential on subhorizon scales (Lue, Scoccimarro, Starkmann 2004; Koyama & Maartins 2005)
- Energy-momentum conservation and dominance of Weyl anisotropy leads to further decay



Sawicki, Song & Hu 2006

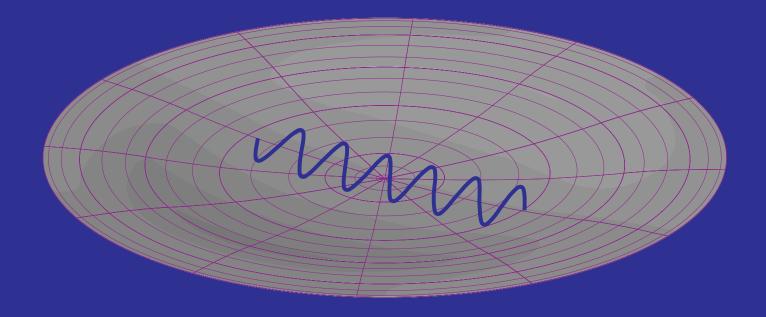
### Integrated Sachs-Wolfe Effect

- CMB photons transit gravitational potentials of large-scale structure
- If potential decays during transit, gravitational blueshift of infall not cancelled by gravitational redshift of exit
- Spatial curvature of gravitational potential leads to additional effect  $\Delta T/T = -\Delta(\Phi \Psi)$



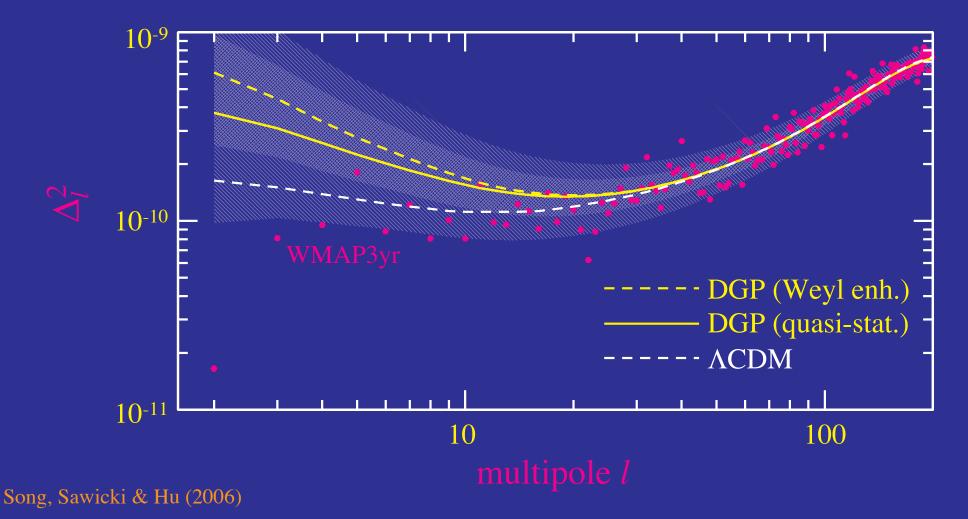
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# **DGP** Example

- Excess decay leads to enhanced large angle CMB anisotropy
- Requires either breaking of initial scale invariance or missing physics beyond Weyl tensor at  $\sim r_c/10$  to be compatible with observations



A Worked Example: f(R) Modified Action Acceleration

# A Worked Example: f(R) Gravity

• Modify the Einstein-Hilbert action (Starobinsky 1980; Carroll et al 2004)

$$S = \int d^4x \sqrt{-g} \left[ \frac{R + f(R)}{2\mu^2} + \mathcal{L}_{\rm m} \right]$$

- In the Jordan frame, gravity becomes 4th order but matter remains minimally coupled and separately conserved
- Satisifies PPF requirements

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- Satisifies PPF requirements
- Expansion history parameterization: Friedmann equation becomes

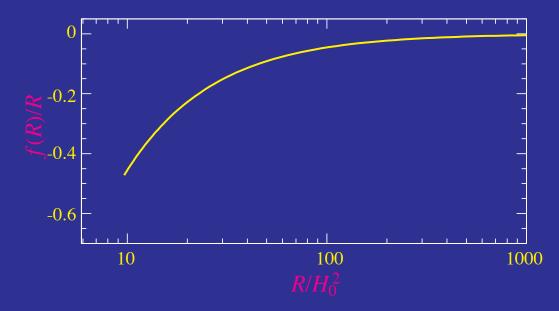
$$H^{2} - f_{R}(HH' + H^{2}) + \frac{1}{6}f + H^{2}f_{RR}R' = \frac{\mu^{2}\rho}{3}$$

where  $f_R = df/dR$ ,  $f_{RR} = d^2 f/dR^2$ 

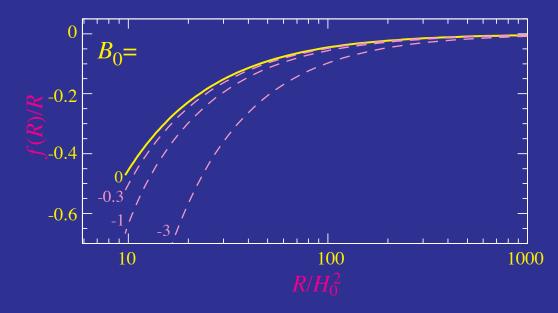
• For any desired H, solve a 2nd order diffeq to find f(R)

**PPF** Functions

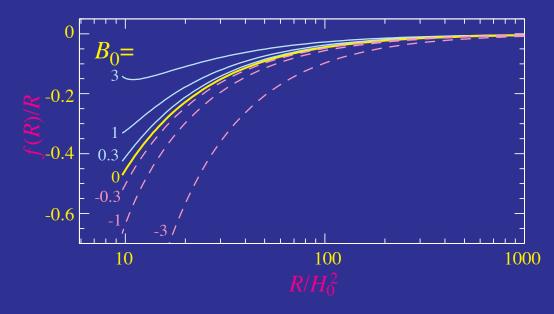
- Each expansion history, matched by dark energy model  $[w(z), \Omega_{DE}, H_0]$  corresponds to a family of f(R) models due to its 4th order nature
- Parameterized by  $B \propto f_{RR} = d^2 f/dR^2$  evaluated at z=0



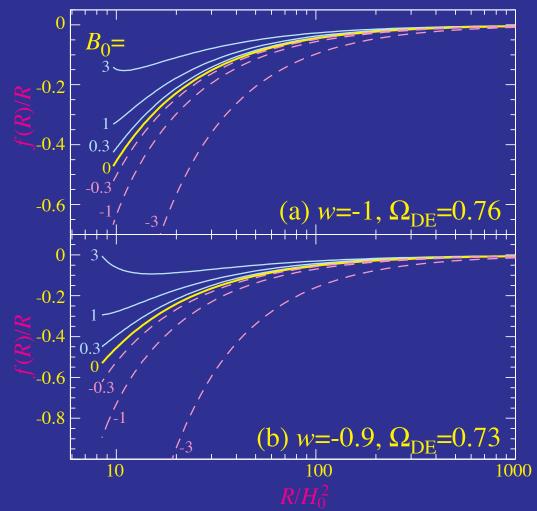
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f(R) Metric Evolution

### **Deviation Parameter**

• Express the 4th order nature of equations as a deviation parameter

$$\Phi'' - \Psi' - \frac{H''}{H'} \Phi' - \left(\frac{H'}{H} - \frac{H''}{H'}\right) \Psi = \left(\frac{k}{aH}\right)^2 B\epsilon$$

• Einstein equation become a second order equation for  $\epsilon$ 

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• In high redshift, high curvature R limit this is

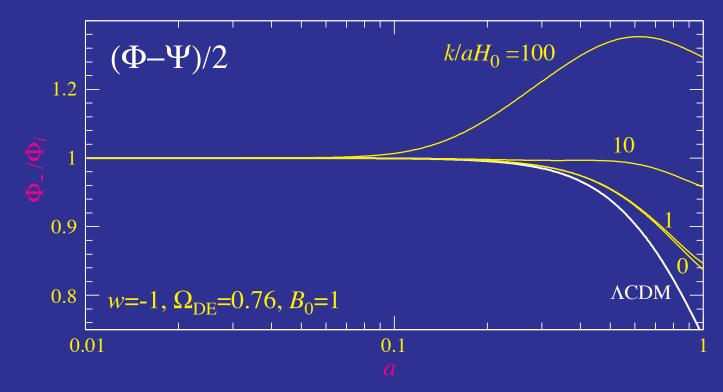
$$\epsilon'' + \left(\frac{7}{2} + 4\frac{B'}{B}\right)\epsilon' + \frac{2}{B}\epsilon = \frac{1}{B} \times \text{ metric sources}$$
$$B = \frac{f_{RR}}{1 + f_R}R'\frac{H}{H'}$$

R→∞, B→ 0 and for B < 0 short time-scale tachyonic instability appears making previous models not cosmologically viable</li>

$$f(R) = -M^{2+2n}/R^n$$

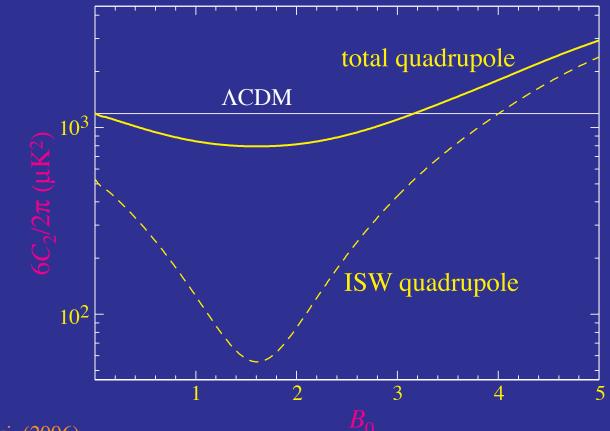
### Potential Growth

- On the stable *B*>0 branch, potential evolution reverses from decay to growth as a function of scale B<sup>1/2</sup>(k/aH)
- Newton constant *G* rescaled by  $1+f_R$  leading to different density and potential growth functions
- On small scales, quasistatic equilibrium reached in linear theory with  $\Psi = -2\Phi$  requiring non-linear effects restore PPN expectations



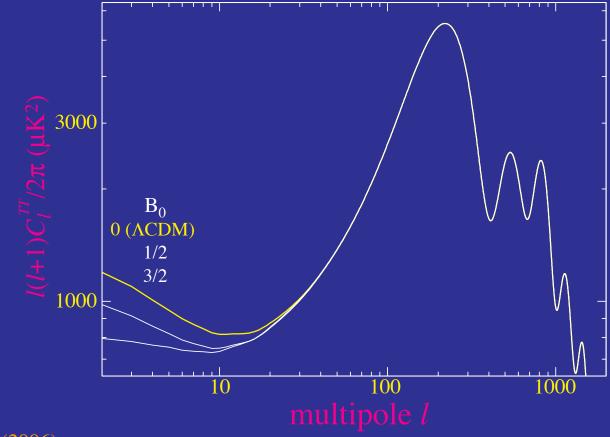
# ISW Quadrupole

- Reduction of potential decay can eliminate the ISW effect at the quadrupole for  $B_0 \sim 3/2$
- In conjunction with a change in the initial power spectrum can also bring the total quadrupole closer in ensemble average to the observed quadrupole



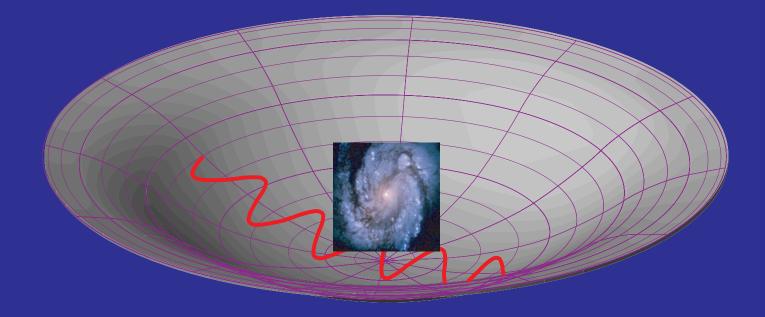
# ISW Quadrupole

- Reduction of large angle anisotropy for  $B_0 \sim 1$  for same expansion history and distances as  $\Lambda CDM$
- Well-tested small scale anisotropy unchanged



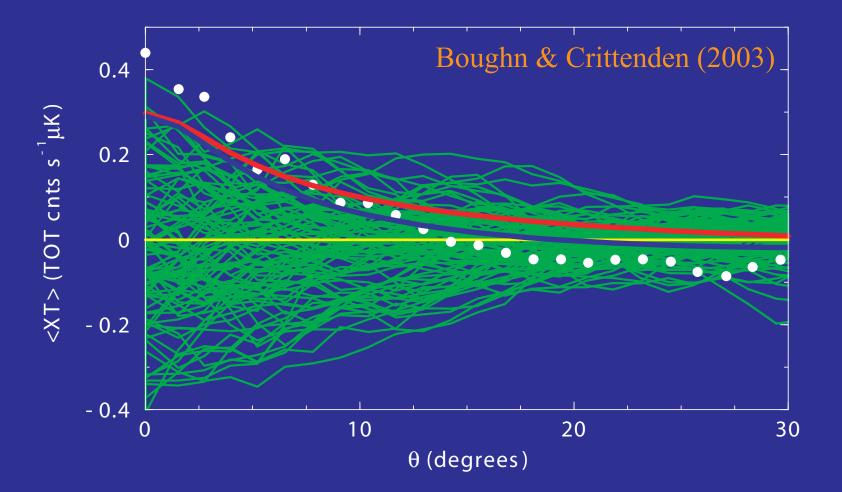
# **ISW-Galaxy** Correlation

- Decaying potential: galaxy positions correlated with CMB
- Growing potential: galaxy positions anticorrelated with CMB
- Observations indicate correlation



### **ISW Galaxy Correlation**

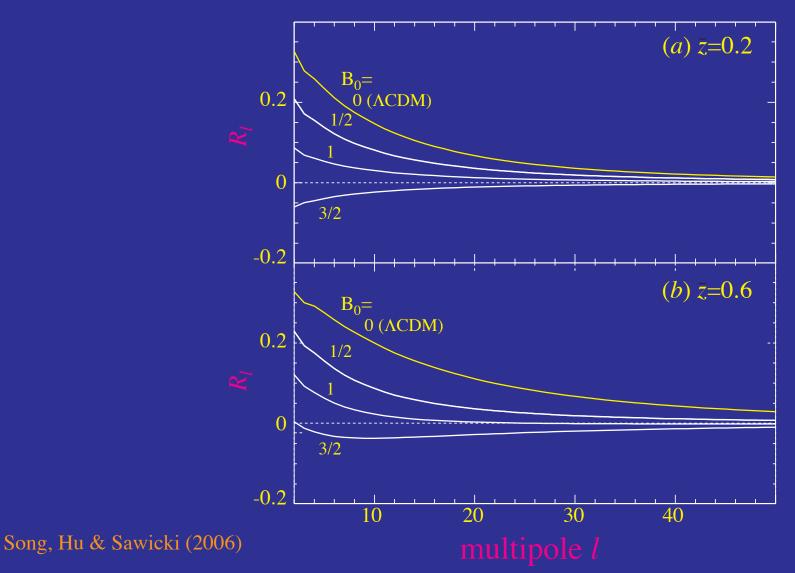
• A 2-3σ detection of the ISW effect through galaxy correlations



Boughn & Crittenden (2003); Nolte et al (2003); Fosalba & Gaztanaga (2003); Fosalba et al (2003); Afshordi et al (2003)

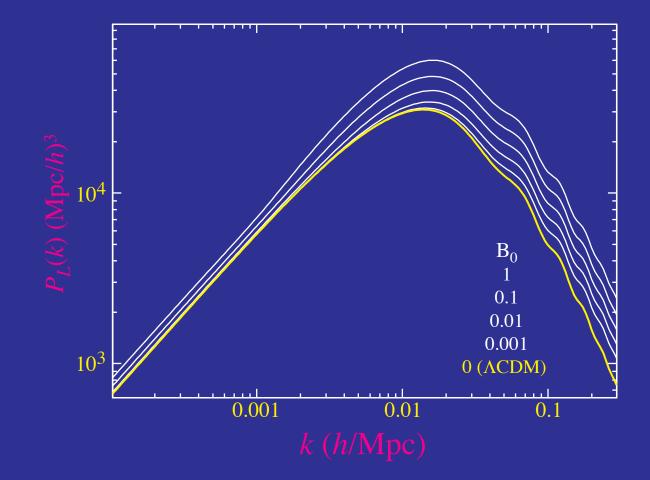
## Galaxy-ISW (Anti)Correlation

- Change in potential growth reduces galaxy-ISW correlation and for high  $B_0>1$  predicts anticorrelation
- Reported positive detections place upper limit of  $B_0 < 1$



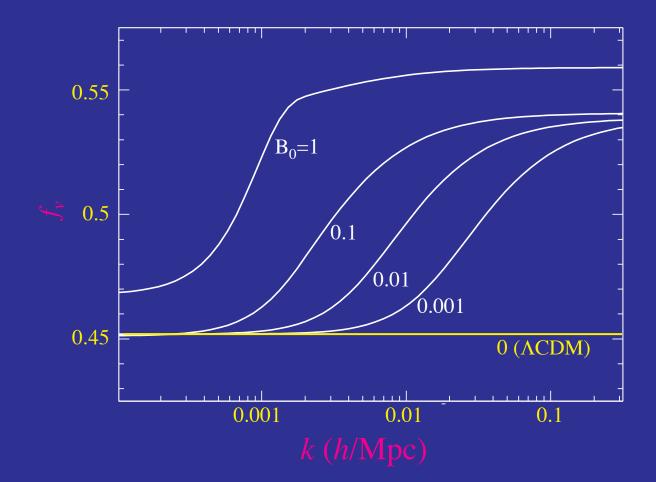
### Linear Power Spectrum

- Linear real space power spectrum enhanced on small scales
- Degeneracy with galaxy bias and lack of non-linear predictions leave constraints from shape of power spectrum



# **Redshift Space Distortion**

- Relationship between velocity and density field given by continuity with modified growth rate
- Redshift space power spectrum further distorted by Kaiser effect



### **PPF** Description

• On superhorizon scales, metric evolution given by conservation

$$\Phi'' - \Psi' - \frac{H''}{H'}\Phi' - \left(\frac{H'}{H} - \frac{H''}{H'}\right)\Psi = 0$$

requiring a closure relation between the metric fluctuations

$$\Psi = -f_1(a)\Phi$$

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requiring a closure relation between the metric fluctuations

 $\Psi = -f_1(a)\Phi$ 

• Below parameterized transition scale, modified Poisson equation

$$k^2\left(\frac{\Phi-\Psi}{2}\right) = \frac{1}{2}f_3(a)\mu^2 a^2\rho\Delta$$

with a potentially different closure relation

$$\Psi = -f_2(a)\Phi$$

and the usual quasistatic conservation laws

$$\Delta' = \left(\frac{k}{aH}\right)^2 H q, \quad H q' = \Psi,$$

# Summary

- Parameterized description of acceleration: background expansion history w(z) supplemented by
  - Transition scale where dark energy becomes smooth
  - Transition scale where modified gravity switches from Friedmann dynamics to quasistatic Newtonian dynamics (and a further non-linear transition to GR)
  - consistent with energy-momentum conservation and metric theory
- Test explanations of acceleration in absence of compelling models

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  - consistent with energy-momentum conservation and metric theory
- Test explanations of acceleration in absence of compelling models
- Expansion history alone tests specific models: e.g. DGP by  $H_0$
- PPF description of DGP shows disfavored enhanced ISW effect if Weyl anisotropy dominates during self-acceleration
- PPF description of f(R) shows previous models unstable but stable models do exist and are testable with linear phenomena