CMB Anisotropy



Wayne Hu II Workshop Challenges Campos do Jordao, Decem<u>ber 2009</u>

Temperature and Polarization Spectra



Across the Horizon



Hu & White (2004); artist:B. Christie/SciAm; available at http://background.uchicago.edu

Acoustic Oscillations

WMAP Data



Gravitational Ringing

- Gravitational potential wells
- Fluid falls into wells, pressure resists: acoustic oscillations

Curvature in the Power SpectrumFeatures scale with angular diameter distance

• Angular location of the first peak

Baryons in the Power Spectrum

Dark Matter in the Power Spectrum

Planck First Light

• Planck satellite, first light survey



Power of Planck's Precision





Fixed Deceleration Epoch

- CMB determination of matter density controls all determinations in the deceleration (matter dominated) epoch
- WMAP5: $\Omega_m h^2 = 0.1326 \pm 0.0063 \rightarrow 5\%$
- Distance to recombination D_* determined to $\frac{1}{4}5\% \approx 1.25\%$
- Expansion rate during any redshift in the deceleration epoch determined to 5%
- Distance to any redshift in the deceleration epoch determined as

$$D(z) = D_* - \int_z^{z_*} \frac{dz}{H(z)}$$

- Volumes determined by a combination $dV = D_A^2 d\Omega dz / H(z)$
- Structure also determined by growth of fluctuations from z_*
- $\Omega_m h^2$ can be determined to ~ 1% from Planck.



- Peaks measure distance to recombination
- ISW effect constrains dynamics of acceleration





- Peaks measure distance to recombination
- ISW effect constrains dynamics of acceleration and early dark energy



Forecasts for $CMB+H_0$

• To complement CMB observations with $\Omega_m h^2$ to 1%, an H_0 of ~1% enables constant *w* measurement to ~2% in a flat universe



ISW Effect

- Gravitational blueshift on infall does not cancel redshift on climbing out
- Contraction of spatial metric doubles the effect: $\Delta T/T = 2\Delta \Phi$
- Effect from potential hills and wells cancel on small scales



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ISW-Galaxy Correlation

- Decaying potential: galaxy positions correlated with CMB
- Growing potential: galaxy positions anticorrelated with CMB
- Observations indicate correlation



Dark Energy Clustering

- ISW effect intrinsically sensitive to dark energy smoothness
- Large angle contributions reduced if clustered



Hu (1998); [plot: Hu & Scranton (2004)]

ISW-Galaxy Correlation

- ~4σ joint detection of ISW correlation with large scale structure (galaxies)
- $\sim 2\sigma$ high compared with Λ CDM



• Like light scattering off surface, scattering with electrons polarizes radiation



• Unlike sunlight, CMB radiation comes from and electrons respond in all directions



• If intensity differs at 90 degrees (quadrupole anisotropy), net linear polarization



Whence Quadrupoles?

- Temperature inhomogeneities in a medium
- Photons arrive from different regions producing an anisotropy



(Scalar) Temperature Inhomogeneity

Hu & White (1997)

Patterns on the Sky

 Projection of the quadrupole moment across sky determines polarization pattern



E-mode Polarization

 Polarization points along the direction of plane wave or amplitude variation



Acoustic Peaks in the Polarization

- Scalar quadrupole follows the velocity perturbation
- Acoustic velocity out of phase with acoustic temperature
- Correlation oscillates at twice the frequency



Hu & White (1997)

Power Spectrum Present



Polarized Landscape



B-mode Polarization

E-mode Polarization

 Polarization points along the direction of plane wave or amplitude variation



Electric & Magnetic Polarization

(a.k.a. gradient & curl)

• Alignment of principal vs polarization axes (curvature matrix vs polarization direction)



Kamionkowski, Kosowsky, Stebbins (1997) Zaldarriaga & Seljak (1997)

Gravitational Lensing

Example of CMB Lensing

- Toy example of lensing of the CMB primary anisotropies
- Shearing of the image
Polarization Lensing

• Since E and B denote the relationship between the polarization amplitude and direction, warping due to lensing creates B-modes



Zaldarriaga & Seljak (1998) [figure: Hu & Okamoto (2001)]

Polarized Landscape



Lensed Power Spectrum Observables

- Principal components show two observables in lensed power spectra
- Temperature and E-polarization: deflection power at *l*~100
 B-polarization: deflection power at *l*~500
- Normalized so that observables error = fractional lens power error



Redshift Sensitivity

• Lensing observables probe distance and structure at high redshift $\frac{\delta\Theta_i}{\Theta_i} = \left[\left(3 - \frac{d\ln\Delta_m^2}{d\ln k} \right) \frac{\delta D_A}{D_A} - \frac{\delta H}{H} + 2\frac{\delta G}{G} + 2\frac{\delta D_A(D_s - D)}{D_A(D_s - D)} \right]$



Smith, Hu & Kaplinghat (2006)

Lensing Observables

- Lensing observables provide a simple way of accounting for non-Gaussianity and parameter degeneracies
- Direct forecasts for Planck + 10% sky with noise $\Delta_P=1.4uK'$



Smith, Hu, Kaplinghat (2006) [see also: Kaplinghat et. al 2003, Acquaviva & Baccigalupi 2005, Smith et al 2005, Li et al 2006]

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Constraints on Lensing Observables

- Lensing observables in T,E are limited by CMB sample variance
- Lensing observables in **B** are limited by lens sample variance
- B-modes require 10x as much sky at high signal-to-noise or 3x as much sky at the optimal signal-to-noise with Δ_P =4.7uK'



Smith, Hu & Kaplinghat (2006)

High Signal-to-Noise B-modes

- Cosmic variance of CMB fields sets ultimate limit for *T*,*E*
- *B*-polarization allows mapping to finer scales and in principle is not limited by cosmic variance of *E* (Hirata & Seljak 2003)



mass

temp. reconstruction EB pol. reconstruction 100 sq. deg; 4' beam; 1µK-arcmin

Hu & Okamoto (2001)

Lensing-Galaxy Correlation

- ~3σ+ joint detection of WMAP lensing reconstruction with large scale structure (galaxies)
- Consistent with ACDM



Matter Power Spectrum

 Measuring projected matter power spectrum to cosmic variance limit across whole linear regime 0.002< k < 0.2 h/Mpc



Hu & Okamoto (2001)

Reionization

Across the Horizon



Hu & White (2004); artist:B. Christie/SciAm; available at http://background.uchicago.edu

Anisotropy Suppression

 A fraction τ~0.1 of photons rescattered during reionization out of line of sight and replaced statistically by photon with random temperature flucutuation - suppressing anisotropy as e^{-τ}



Reionization Suppression

• Rescattering suppresses primary temperature and polarization anisotropy according to optical depth, fraction of photons rescattered



Tilt-τ Degeneracy

Only anisotropy at reionization (high k), not isotropic temperature fluctuations (low k) - is suppressed leading to effective tilt for WMAP (not Planck)



Distance Predicts Growth

• With smooth dark energy, distance predicts scale-invariant growth to a few percent - a falsifiable prediction



Mortonson, Hu, Huterer (2008)

Temperature Inhomogeneity

- Temperature inhomogeneity reflects initial density perturbation on large scales
- Consider a single Fourier moment:



Locally Transparent

• Presently, the matter density is so low that a typical CMB photon will not scatter in a Hubble time (~age of universe)



Reversed Expansion

• Free electron density in an ionized medium increases as scale factor *a*-³; when the universe was a tenth of its current size CMB photons have a finite (~10%) chance to scatter



Polarization Anisotropy

• Electron sees the temperature anisotropy on its recombination surface and scatters it into a polarization



Temperature Correlation

• Pattern correlated with the temperature anisotropy that generates it; here an *m*=0 quadrupole



Instantaneous Reionization

- WMAP data constrains optical depth for instantaneous models of τ =0.087±0.017
- Upper limit on gravitational waves weaker than from temperature



Ionization History

• Two models with same optical depth τ but different ionization history



Kaplinghat et al. (2002); Hu & Holder (2003)

Distinguishable History

 Same optical depth, but different coherence - horizon scale during scattering epoch



Transfer Function

Linearized response to delta function ionization perturbation $T_{\ell i} \equiv rac{\partial \ln C_{\ell}^{EE}}{\partial x(z_i)}, \qquad \delta C_{\ell}^{EE} = C_{\ell}^{EE} \sum_i T_{\ell i} \delta x(z_i)$ $z_i = 25$ 0.4 0.2 $z_i = 8$ 0

10

100

Hu & Holder (2003)

Principal Components

• Eigenvectors of the Fisher Matrix



Hu & Holder (2003)

Representation in Modes

 Reproduces the power spectrum with sum over >3 modes more generally 5 modes suffices: e.g. total τ=0.1375 vs 0.1377



Hu & Holder (2003)

Large Scale Anomalies

Low Quadrupole

• Temperature quadrupole is low compared with best fit ACDM model



Quadrupole Origins

• Transfer function for the quadrupole



Gordon & Hu (2004)

Horizon-Scale Power

 Polarization is a robust indicator of horizon scale power and disfavors suppression as explanation of low quadrupole independently of ionization or acceleration model



Mortonson & Hu (2009)

Quadrupole Origins

• Transfer function for the quadrupole



Gordon & Hu (2004)

Model-Independent Reionization

- All possible ionization histories at *z*<30
- Detections at 20 < l < 30 required to further constrain general ionization which widens the τn_s degeneracy allowing $n_s = 1$
- Quadrupole & octopole predicted to better than cosmic variance test ACDM for anomalies



Mortonson & Hu (2008)

Large Angle Anomalies

- Low planar quadrupole aligned with planar octopole
- More power in south ecliptic hemisphere
- Non-Gaussian spot



Polarization Tests

Matching polarization anomalies if cosmological



Dvorkin, Peiris, Hu (2007)

Polarization Bumps

• If features in the temperature spectrum reflect features in the power spectrum (inflationary potential), reflected in polarization with little ambiguity from reionization


Inflation

Inflation Past

- Superhorizon correlations (acoustic coherence, polarization corr.)
- Spatially flat geometry (angular peak scale)
- Adiabatic fluctuations (peak morphology)
- Nearly scale invariant fluctuations (broadband power, small red tilt favored)
- Gaussian fluctuations
 (but *f*n1>few would rule out single field slow roll)

Inflationary Observables

• Curvature Power Spectrum:

$$\Delta_{\mathcal{R}}^{2} \approx \frac{8\pi G}{2} \frac{1}{\epsilon} \left(\frac{H}{2\pi}\right)^{2}, \quad \epsilon = \frac{1}{2} \frac{1}{8\pi G} \left(\frac{V'}{V}\right)^{2}$$

• Tilt

$$\frac{d\ln\Delta_{\mathcal{R}}^2}{d\ln k} = n_S - 1 = -4\epsilon - 2\delta$$

where

$$\delta = \epsilon - \frac{1}{8\pi G} \frac{V''}{V}$$

So for featureless potentials e.g. monomial ϕ^n , $\epsilon \sim |\delta|$

• Running $dn_S/d\ln k$ second order

Inflationary Observables

• Gravitational Wave (Tensor) Power Spectrum:

$$\Delta_{+,\times}^2 = 16\pi G \left(\frac{H}{2\pi}\right)^2$$

reflects energy scale of inflation $H^2 \propto V \equiv E_i^4$

$$\Delta B_{\mathrm{peak}} pprox 0.024 \left(rac{E_i}{10^{16} \mathrm{GeV}}
ight)^2 \mu \mathrm{K}$$

• Tensor-Scalar Ratio, Tilt:

$$r \equiv 4 \frac{\Delta_+^2}{\Delta_R^2} = 16\epsilon, \quad \frac{d\ln\Delta_+^2}{d\ln k} \equiv n_T = 2 \frac{d\ln H}{d\ln k} = -2\epsilon$$

• Consistency:

$$r = -8n_T$$

Quadrupoles from Gravitational Waves

- Transverse-traceless distortion provides temperature quadrupole
- Gravitational wave polarization picks out direction transverse to wavevector



Electric & Magnetic Polarization

(a.k.a. gradient & curl)

• Alignment of principal vs polarization axes (curvature matrix vs polarization direction)



Kamionkowski, Kosowsky, Stebbins (1997) Zaldarriaga & Seljak (1997)

Gravitational Wave Pattern

- Projection of the quadrupole anisotropy gives polarization pattern
- Transverse polarization of gravitational waves breaks azimuthal symmetry





density perturbation

gravitational wave

Energy Scale of Inflation

- Amplitude of B-mode peak scales as square of energy scale (Hubble parameter) during inflation, power as E_i^4
- Good: upper limits are at GUT scale. Bad: secondaries & foregrounds



Large Field, Small Field Models

• For detectable gravitational waves r>0.01, scalar field must roll by order $M_{\rm pl}=(8\pi G)^{-1/2}$

$$\frac{d\phi}{dN} = \frac{d\phi}{d\ln a} = \frac{d\phi}{dt}\frac{1}{H}$$

• The larger ϵ is the more the field rolls in an e-fold

$$\epsilon = \frac{r}{16} = \frac{3}{2V} \left(H \frac{d\phi}{dN} \right)^2 = \frac{8\pi G}{2} \left(\frac{d\phi}{dN} \right)^2$$

• Observable scales span $\Delta N \sim 5$ so

$$\Delta \phi > 5 \frac{d\phi}{dN} = 5 \left(\frac{r}{8}\right)^{1/2} M_{\rm pl} \approx 0.2 \left(\frac{r}{0.01}\right)^{1/2} M_{\rm pl}$$

- Does this make sense as an effective field theory? Lyth (1997)
- Small field models where ϕ near maximum more reasonable?
- Large field existence proof: monodromy Silverstein & Westphal (2008)
 ...theorists running around in circles...

Inflation Present

- Tilt indicates that one of the slow roll parameters finite (ignoring exotic high-z reionization)
- Upper limit on gravity waves put an upper limit on *V*/*V* and hence an upper limit on how far the inflaton rolls

- Constraints in the *r*-*n*_s plane test classes of models
- Given functional form of *V*, constraints on the flatness of potential when the horizon left the horizon predict too many (or few) efolds of further inflation

- Non-Gaussian fluctuations at *f*nl~50?
- Glitches and large scale anomalies

Inflationary Constraints

- Tilt mildly favored over tensors as explaining small scale suppression
- Specific models of inflation relate $r-n_s$ through V', V''
- Small tensors and $n_s \sim 1$ may make inflation continue for too many efolds



fnl

- Local second order non-Gaussianity: $\Phi_{nl}=\Phi+f_{nl}(\Phi^2-\langle\Phi^2\rangle)$
- WMAP3 Kp0+: 27<*f*_{nl}<147 (95% CL) (Yadav & Wandelt 2007)
- WMAP5 opt: $-4 < f_{nl} < 80 (95\% \text{ CL})$ (Smith, Senatore & Zaldarriaga 2009)



Local Non-Gaussianity

• Local non-Gaussianity couples long to short wavelength fluctuations

Temperatur $f_{NL}=0$



Ligouri et al (2007)

Local Non-Gaussianity

• Local non-Gaussianity couples long to short wavelength fluctuations

Temperature f_{NL} =3000



Ligouri et al (2007)

Inflation Future

- Planck can test Gaussianity down to fn1~few and make a high significance detection if fn1~50
- Planck will provide a high significance measurement of tilt (n_s-1)
- Planck will test constancy of tilt significant deviation would rule out all standard slow roll models
- Gravitational wave power proportional to energy scale to 4th power
- B-modes potentially observable for V^{1/4}>3 x 10¹⁵ GeV with removal of lensing B-modes and foregrounds
- Measuring both the reionization bump and recombination peak tests slow roll consistency relation by constraining tensor tilt

Consistency Relation & Reionization

- By assuming the wrong ionization history can falsely rule out consistency relation
- Principal components eliminate possible biases



Summary

- CMB acoustic peaks provide precision measurements of
 - baryon density $\Omega_b h^2$
 - matter density $\Omega_m h^2$
 - distance to recombination D_*
 - amplitude and tilt of spectrum A_S , n_S

and all should reach better than 1% precision with Planck

- CMB polarization provides windows on reionization, large scale anomalies, intervening matter and expansion history (through lensing)
- Inflationary origins will be tested further by precision spectrum, gravitational waves and non-Gaussianity
- Single-field slow roll inflation is highly falsifiable